

## ON THE PERIODIC WAVE SOLUTIONS OF THE (2+1)-D KONOPELCHENKO-DUBROVSKY EQUATION IN FLUID MECHANICS

by

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*The central orientation of this work is to plumb the (2+1)-D Konopelchenko-Dubrovsky equation that is utilized widely to describe certain non-linear phenomena in the field of the fluid mechanics. Two effective methods namely the variational method and the energy balance theory are employed to construct the periodic wave solutions. As predicted, the results extracted by these two approaches are almost identical, which is anticipated to offer some new viewpoints to the exploration of the periodic wave theory in physics.*

Key words: *periodic wave solution, variational principle, energy balance theory, variational method, Hamiltonian*

### Introduction

A good many kinds of complex problems in nature can be described and modeled by the non-linear partial differential equations (NPDE) [1-5]. However, because some basic properties of linear differential equations are no longer tenable in non-linear differential equations, it is difficult to use a unified method to solve the NPDE [6-8]. Therefore, for a long time, solving the exact solutions of NPDE has been a hot topic for researchers [9-12]. In this study, we will look into the (2+1)-D Konopelchenko-Dubrovsky equation (KDE) [13,14]:

$$\begin{aligned} \phi_t - \phi_{xxx} - 6b\phi\phi_x + \frac{3}{2}a\phi^2\phi_x - 3\phi_y + 3a\phi\phi_x &= 0 \\ \phi_y &= \varphi_x \end{aligned} \quad (1)$$

where  $\phi = \phi(x, y, t)$ ,  $\varphi = \varphi(x, y, t)$ , and  $a$  and  $b$  are the non-zero real parameters. Equation (1) was proposed by Konopelchenko and Dubrovsky [15] in 1984 and was used extensively to model complex phenomenon in fluid mechanics. Some different efficacious approaches have been adopted to deal with eq. (1) such as the extended F-expansion method [16], modified auxiliary equation approach [17], modified simplest equation method [18], tanh-sech method [19], improved tanh function approach [20], mapping approach [21] and others [22-25]. In this paper, we aim to construct the periodic wave solutions (PWS) of eq. (1) by taking advantage of the variational method (VM) and energy balance theory (EBT).

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### The variational principle and the Hamiltonian

For obtaining the periodic solutions, we first apply the traveling waves transformations for eq. (1):

$$\phi(x, y, t) = \phi(\eta), \quad \varphi(x, y, t) = \varphi(\eta), \quad \eta = x + y + cy + \eta_0 \quad (2)$$

Then, we obtain:

$$c\phi' - \phi''' - 6b\phi\phi' + \frac{3}{2}a\phi^2\phi' - 3\varphi' + 3a\varphi\phi' = 0 \quad (3)$$

$$\phi' = \varphi'$$

Taking a integration of the second equation in eq. (3) yields:

$$\phi = \varphi \quad (4)$$

Taking it into the first equation of eq. (3), then integrating the obtained equation one and taking the integration constant as zero, we get:

$$(c-3)\phi - \left(3b + \frac{3}{2}a\right)\phi^2 + \frac{a^2}{2}\phi^3 - \phi'' = 0 \quad (5)$$

Then the variational principle can be established by manipulating the semi-inverse approach [26-29]:

$$J(\phi) = \int \left\{ \frac{1}{2}(c-3)\phi^2 - \left(b + \frac{1}{2}a\right)\phi^3 + \frac{a^2}{8}\phi^4 + \frac{1}{2}(\phi')^2 \right\} d\eta \quad (6)$$

For simplicity, we re-write eq. (6) as:

$$J(\phi) = \int \left\{ m\phi^2 - n\phi^3 + k\phi^4 + \frac{1}{2}(\phi')^2 \right\} d\eta \quad (7)$$

where

$$m = \frac{1}{2}(c-3), \quad n = \left(b + \frac{1}{2}a\right), \quad k = \frac{a^2}{8}$$

eq. (7) can be written

$$J(\phi) = \int \left\{ m\phi^2 - n\phi^3 + k\phi^4 + \frac{1}{2}(\phi')^2 \right\} d\eta = \int (\Xi - \Theta) d\eta \quad (8)$$

where  $\Xi$  is the kinetic energy and  $\Theta$  – the stands for the potential energy:

$$\Omega = \frac{1}{2}(\phi')^2 \quad (9)$$

$$\Theta = -m\phi^2 + n\phi^3 - k\phi^4 \quad (10)$$

Thus we can get the system's Hamiltonian:

$$H = \Omega + \Theta = \frac{1}{2}(\phi')^2 - m\phi^2 + n\phi^3 - k\phi^4 \quad (11)$$

### The periodic wave solutions

#### The variational method

To seek the PWS of eq. (1) via the VM [30, 31], we postulate the periodic solution of eq. (5) is:

$$\phi(\eta) = \Xi \cos(\omega\eta) \quad (12)$$

Substituting it into eq. (7), we have:

$$J(\Xi) = \int_0^{T/4} \left\{ m\Xi^2 \cos^2(\omega\eta) - n\Xi^3 \cos^3(\omega\eta) + k\Xi^4 \cos^4(\omega\eta) + \frac{1}{2}\Xi^2 \omega^2 \sin^2(\omega\eta) \right\} d\eta \quad (13)$$

By the VM, there is:

$$\frac{dJ}{d\Xi} = 0 \quad (14)$$

That is:

$$\begin{aligned} \frac{dJ(\Xi)}{d\Xi} &= \frac{d}{d\Xi} \int_0^{T/4} \left\{ m\Xi^2 \cos^2(\omega\eta) - n\Xi^3 \cos^3(\omega\eta) + k\Xi^4 \cos^4(\omega\eta) + \frac{1}{2}\Xi^2 \omega^2 \sin^2(\omega\eta) \right\} d\eta = \\ &= \frac{1}{\omega} \frac{d}{d\Xi} \int_0^{\pi/2} \left\{ m\Xi^2 \cos^2(\psi) - n\Xi^3 \cos^3(\psi) + k\Xi^4 \cos^4(\psi) + \frac{1}{2}\Xi^2 \omega^2 \sin^2(\psi) \right\} d\psi = \\ &= \frac{1}{\omega} \int_0^{\pi/2} \left\{ 2m\Xi \cos^2(\psi) - 3n\Xi^2 \cos^3(\psi) + 4k\Xi^3 \cos^4(\psi) + \Xi \omega^2 \sin^2(\psi) \right\} d\psi = 0 \end{aligned} \quad (15)$$

The expression can be obtained:

$$\omega^2 = \frac{\int_0^{\pi/2} \left\{ -2m\Xi \cos^2(\psi) + 3n\Xi^2 \cos^3(\psi) - 4k\Xi^3 \cos^4(\psi) \right\} d\psi}{\int_0^{\pi/2} \Xi \sin^2(\psi) d\psi} \quad (16)$$

which leads to:

$$\omega = \sqrt{\frac{8n\Xi - 3k\pi\Xi^2 - 2m\pi}{\pi}} \quad (17)$$

Thus, the periodic solution of eq. (5) can be:

$$\phi(\eta) = \Xi \cos \left( \sqrt{\frac{8n\Xi - 3k\pi\Xi^2 - 2m\pi}{\pi}} \eta \right) \quad (18)$$

Finally, we can attain the PWS of eq. (1):

$$\phi(x, y, t) = \varphi(x, y, t) = \Xi \cos \left[ \sqrt{\frac{8n\Xi - 3k\pi\Xi^2 - 2m\pi}{\pi}} (x + y + ct + \eta_0) \right] \quad (19)$$

where

$$m = \frac{1}{2}(c-3), \quad n = \left( b + \frac{1}{2}a \right), \quad k = \frac{a^2}{8}$$

### The energy balance theory

To apply the EBT [32], the periodic solution of eq. (5) is assumed:

$$\phi(\eta) = \Xi \cos(\omega\eta), \quad \omega > 0 \quad (20)$$

Based on the Hamiltonian eq. (11), we can get the Hamiltonian constant :

$$H_0 = \Omega + \Theta = -m\Xi^2 + n\Xi^3 - k\Xi^4 \quad (21)$$

The EBT tells that the energy of the system keep unchanged with the  $\eta$  changes:

$$H = \Omega + \Theta = \frac{1}{2}[-\Xi\omega\sin(\omega\eta)]^2 - m[\Xi\cos(\omega\eta)]^2 + n[\Xi\cos(\omega\eta)]^3 - k[\Xi\cos(\omega\eta)]^4 = H_0 \quad (22)$$

which is:

$$\begin{aligned} & \frac{1}{2}[-\Xi\omega\sin(\omega\eta)]^2 - m[\Xi\cos(\omega\eta)]^2 + n[\Xi\cos(\omega\eta)]^3 - \\ & - k[\Xi\cos(\omega\eta)]^4 = -m\Xi^2 + n\Xi^3 - k\Xi^4 \end{aligned} \quad (23)$$

Without loss of generality, here we set:

$$\omega\eta = \frac{\pi}{4} \quad (24)$$

such that

$$\frac{1}{2}\left(-\frac{\sqrt{2}}{2}\Xi\omega\right)^2 - m\left(\frac{\sqrt{2}}{2}\Xi\right)^2 + n\left(\frac{\sqrt{2}}{2}\Xi\right)^3 - k\left(\frac{\sqrt{2}}{2}\Xi\right)^4 = -m\Xi^2 + n\Xi^3 - k\Xi^4 \quad (25)$$

which leads to:

$$\omega = \sqrt{(4-\sqrt{2})n\Xi - 3k\Xi^2 - 2m} \quad (26)$$

Obviously, we can discover that eq. (26) has a well agreement with eq. (18), this strongly confirms the two method are correct and reliable. By eq. (26), the PWS to eq. (1) is found as:

$$\phi(x, y, t) = \phi(x, y, t) = \Xi \cos \left[ \sqrt{(4-\sqrt{2})n\Xi - 3k\Xi^2 - 2m} (x + y + ct + \eta_0) \right] \quad (27)$$

where

$$m = \frac{1}{2}(c-3), \quad n = \left(b + \frac{1}{2}a\right), \quad k = \frac{a^2}{8}$$

## Conclusion

The (2+1)-D KDE is explored in this paper. By applying the VM and EBT, the periodic wave solutions are obtained. And it finds that the two wave solutions extracted by the two methods have a well agreement. The outcomes reveal that the proposed methods are correct and efficacious, and are anticipated to provide some new ideas to the research of the periodic wave theory.

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