

APPROXIMATE ANALYTICAL SOLUTION OF GENERALIZED FRACTAL EQUAL-WIDTH WAVE EQUATION

by

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In this paper, a generalized equal width wave equation involving space fractal derivatives and time Caputo fractional derivatives is studied and its approximate analytical solution is presented by the Adomian decomposition method. An example shows that the method is efficient to solve fractal non-linear partial differential equations.

Key word: *generalized fractal equal width wave equation, caputo derivative, space fractal derivative, adomian decomposition method*

Introduction

Fractional calculus has been received a skyrocketing attention from mathematics, physics to engineering, it can model effectively a seemingly stochastic diffusion [1, 2] or seemingly irregular Brownian motion [3], it can be also used for noise detection [4], and thermal displacement prediction [5]. The machine learning or deep learning [6-10] will become more powerful to deal with innumerable data for neural networks and imagine processing, and now a fractional model makes a complex and anomalous problem accessible.

A wave near the coast is affected by various random factors including the unsmooth boundary and its memory property, the traditional models can not model these factors, and a fractal-fractional model has to be considered. In this study, we consider the following generalized fractal-fractional equal width wave equation:

$$D_t^\alpha u(x,t) + au^k \frac{\partial u}{\partial x^\beta} - D_t^\alpha \left(b \frac{\partial u}{\partial x^{2\beta}} \right) = 0, \quad 0 < \alpha \leq 1, \quad 0 < \beta \leq 1 \quad (1)$$

subject to the initial condition:

$$u(x,0) = \phi \left[\frac{x^\beta}{\Gamma(1+\beta)} \right] \quad (2)$$

where D_t^α is the Caputo time fractional derivative of order α [11, 12], $\partial/\partial x^\beta$ – the He's space fractal derivative of order β [13-16], ϕ – the known function, and a , k , and b – the positive parameters.

When $\alpha = \beta = 1$, eq. (1) was first introduced by Morrison, *et al.* [17] as the model equation to describe non-linear dispersive waves, and it has a simple relation with the Benja-

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min-Bona-Mahony equation [18-20]. Other similar equations were widely appeared in literature, and much achievement has been obtained. In this paper, we consider a modification with the Caputo time-fractional derivative and He's space fractal derivative [13-16]. This adaptation considers that the Caputo fractional derivative has the memory property [21], and He's fractal derivative can effectively model phenomena in a porous medium [22-27].

Though eq. (1) can model exactly a memorial wave travelling along an unsmooth boundary, it is difficult to be solved analytically, some famous analytical methods, *e.g.*, the homotopy perturbation method [28-32], are not applicable to eq. (1) directly. This paper adopts the Adomian decomposition method (ADM) [33-36] for this purpose.

Basic definitions

In this section, we recall the basic definitions of fractional calculus and fractal calculus which shall be used in this paper. For more details see [37-41].

Definition 1. A real function $f(x)$, $x > 0$ is said to be in the space C_λ , $\lambda \in \mathbb{R}$ if there exists a real number $p > \lambda$, such that $f(x) = x^p f_1(x)$ where $f_1(x) \in C[0, \infty)$ and it is said to be in the space C_n if and only if $f^{(n)} \in C_\lambda$, $n \in \mathbb{N}$.

Definition 2. The Riemann-Liouville fractional integral operator of order $\alpha > 0$ of a function $f(x) \in C_\lambda$, $\lambda \geq -1$ is defined:

$$J^\alpha f(x) = \frac{1}{\Gamma(\alpha)} \int_0^x (x-s)^{\alpha-1} f(s) ds \quad (3)$$

while the definition of Riemann-Liouville fractional derivative is:

$${}^{RL}D_t^q u(x,t) = \frac{\partial^n}{\partial t^n} [J_t^{n-q} u(x,t)] \quad (4)$$

where $n-1 \leq q < n$ and n is an integer.

Properties of the operator J^α can be found in [37] and we mention only the following: For $\alpha, \beta \geq 0$, $x > 0$, and $\lambda > -1$:

$$J^\alpha J^\beta f(x) = J^{\alpha+\beta} f(x) \quad (5)$$

$$J^\alpha J^\beta f(x) = J^\beta J^\alpha f(x) \quad (6)$$

$$J^\alpha (x^\lambda) = \frac{\Gamma(\lambda+1)}{\Gamma(1+\alpha+\lambda)} x^{\lambda+\alpha} \quad (7)$$

Definition 3. The time fractional derivative of $u(x, t)$ in Caputo sense is defined:

$$D_t^\alpha u(x,t) = \frac{1}{\Gamma(m-\alpha)} \int_0^t (t-s)^{m-\alpha-1} u_s^{(m)}(x,s) ds \quad (8)$$

for $m-1 < \alpha \leq m$, $m \in \mathbb{N}^+$, $x > 0$, and $u(x,t) \in C_{-1}^m$.

Definition 4. The space fractal derivative of $u(x, t)$ in the He's sense is defined as [41]:

$$\frac{\partial u(x,t)}{\partial x^\beta} = \Gamma(1+\beta) \lim_{\substack{x_1-x \rightarrow \Delta x \\ \Delta x \neq 0}} \frac{u(x_1,t) - u(x,t)}{(x_1-x)^\beta}, \quad 0 < \beta \leq 1 \quad (9)$$

and

$$\frac{\partial u(x,t)}{\partial x^{2\beta}} = \frac{\partial}{\partial \beta} \left[\frac{\partial u(x,t)}{\partial x^\beta} \right] \quad (10)$$

The fractal derivative is also called as the two-scale fractal derivative, and it can be widely applied to various unsmooth boundary value problems and porous medium problems, [42-47].

Adomian decomposition method

The ADM [33-36] is a technique for solving nonlinear equations in the form:

$$u(x,t) = \tau + \Theta(u) \quad (11)$$

where $\Theta: H \rightarrow H$ is a non-linear mapping from a Banach space H into itself and $\tau \in H$ is known.

The Adomian decomposition method assumes that the solution u can be expanded as an infinite series:

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t) \quad (12)$$

and the non-linear term $\Theta(u)$ can be decomposed:

$$\Theta\left(\sum_{n=0}^{\infty} u_n\right) = \sum_{n=0}^{\infty} A_n(u) \quad (13)$$

where the polynomials $A_n(u)$ are given by:

$$A_n(u_0, u_1, \dots, u_n) = \frac{1}{n!} \frac{\partial^n}{\partial q^n} \left[\Theta\left(\sum_{k=0}^n q^k u_k\right) \right]_{q=0}, \quad n = 0, 1, 2, \dots \quad (14)$$

Substituting eqs. (12) and (13) into (11) gives:

$$\sum_{n=0}^{\infty} u_n = \tau + \sum_{n=0}^{\infty} A_n \quad (15)$$

which is satisfied formally if we set:

$$u_0(x,t) = \tau \quad (16)$$

$$u_1 = A_0 \quad (17)$$

⋮

$$u_{m+1} = A_m \quad (18)$$

Then k -term approximate solution of eq. (11) is given by $u = u_0 + u_1 + \dots + u_{k-1}$.

Approximate solutions

In this section, we derive the main algorithm for solving the problem (1)-(2). Now, eq. (1) can be written:

$$L_t u(x, t) = {}^{RL}D_t^{1-\alpha} \left[au^k \frac{\partial u}{\partial x^\beta} + D_t^\alpha \left(b \frac{\partial u}{\partial x^{2\beta}} \right) \right] \quad (19)$$

where

$$L_t \equiv \frac{\partial}{\partial t}$$

which is an easily invertible linear operator.

Taking the operator L_t^{-1} on both sides of eq. (19), we obtain:

$$u(x, t) = u(x, 0) + L_t^{-1} \left[{}^{RL}D_t^{1-\alpha} \left(au^k \frac{\partial u}{\partial x^\beta} \right) \right] + b \frac{\partial u}{\partial x^{2\beta}} \quad (20)$$

Suppose that the solutions take the form:

$$u(x, t) = \sum_{k=0}^{\infty} u_k(x, t) \quad (21)$$

and the non-linear term can be decomposed:

$$u^k \frac{\partial u}{\partial x^\beta} = \sum_{n=0}^{\infty} A_n \quad (22)$$

where the polynomials A_n can be expressed:

$$A_0 = u_0^k \frac{\partial}{\partial x^\beta} u_0 \quad (23)$$

$$A_1 = ku_0^{k-1} u_1 \frac{\partial}{\partial x^\beta} u_0 + u_0^k \frac{\partial}{\partial x^\beta} u_1 \quad (24)$$

$$A_2 = \frac{k(k-1)}{2} u_0^{k-2} u_1^2 \frac{\partial u_0}{\partial x^\beta} + ku_0 u_2 \frac{\partial}{\partial x^\beta} + \frac{k}{2} u_0^{k-1} u_1^2 + u_0^k \frac{\partial u}{\partial x^\beta} \quad (25)$$

and so on.

Therefore, by using the ADM, we have:

$$u_0(x, t) = u(x, 0) \quad (26)$$

$$u_1(x, t) = aL_t^{-1} \left({}^{RL}D_t^{1-\alpha} A_0 \right) + b \frac{\partial u_0}{\partial x^{2\beta}} \quad (27)$$

$$u_2(x, t) = aL_t^{-1} \left({}^{RL}D_t^{1-\alpha} A_1 \right) + b \frac{\partial u_1}{\partial x^{2\beta}} \quad (28)$$

$$u_3(x, t) = aL_t^{-1} \left({}^{RL}D_t^{1-\alpha} A_2 \right) + b \frac{\partial u_2}{\partial x^{2\beta}} \quad (29)$$

and so on.

Then k -term approximate solutions of eq. (19) are given by $u = u_0 + u_1 + \dots + u_{k-1}$. To illustrate the effectiveness of the previous algorithm, we consider a special case of eqs. (1) and (2):

$$D_t^\alpha u(x, t) = -u^2 \frac{\partial u}{\partial x^\beta} + D_t^\alpha \left(\frac{\partial u}{\partial x^{2\beta}} \right) \quad (30)$$

with the initial condition:

$$u(x, 0) = \frac{2\sqrt{6}E}{E^2 + 1} \quad (31)$$

where

$$E = \exp \left[\frac{-x^\beta}{\Gamma(1 + \beta)} \right]$$

By eqs. (26)-(29), and (23)-(25), we obtain:

$$\begin{aligned} u_0 &= \frac{2\sqrt{6}E}{E^2 + 1} \\ u_1(x, t) &= \frac{-2\sqrt{6}E(E^2 - 1)t^\alpha}{(E^2 + 1)^2 \Gamma(1 + \alpha)} \\ u_2(x, t) &= \frac{-2\sqrt{6}E(6E^2 - E^4 - 1)t^{2\alpha}}{(E^2 + 1)^3 \Gamma(1 + 2\alpha)} \\ u_3(x, t) &= \frac{2\sqrt{6}E(E^2 - 1)(22E^2 - E^4 - 1)t^{3\alpha}}{(E^2 + 1)^4 \Gamma(1 + 3\alpha)} \\ &\vdots \end{aligned} \quad (32)$$

Thus, the 4-term approximate solution of problem (30)-(31) is given by:

$$u(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + u_3(x, t) \quad (33)$$

where $u_i (i = 0 \sim 3)$ are given in eq. (32).

When $\alpha = \beta = 1$, eqs. (30) and (31) become, respectively:

$$\frac{\partial u}{\partial t} + u^2 \frac{\partial u}{\partial x} - \frac{\partial^3 u}{\partial x^2 \partial t} = 0 \quad (34)$$

$$u(x, 0) = \sqrt{6} \cosh^{-1}(x). \quad (35)$$

The approximate solution is:

$$u(x,t) = \frac{2\sqrt{6}e^{-x}}{1+e^{-2x}} + \frac{-2\sqrt{6}e^{-x}(e^{-2x}-1)t}{(1+e^{-2x})^2} - \frac{\sqrt{6}e^{-x}(6e^{-2x}-e^{-4x}-1)t^2}{(1+e^{-2x})^3} + \frac{\sqrt{6}e^{-x}(e^{-2x}-1)(22e^{-2x}-e^{-4x}-1)t^3}{3(1+e^{-2x})^4} \quad (36)$$

which converges to the exact solution if the iteration process continues.

Conclusion

In this paper, the ADM has been successfully applied to obtaining the approximate analytical solution of generalized fractal equal width wave equation. The example is given to illustrate the validity and accuracy of the method. The results show that the method is efficient to handle fractal wave equations with singularity [48] or He's fractional derivative [49, 50].

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