STUDY ON SINGLE-PHASE AND TWO-PHASE FLOW AND HEAT TRANSFER CHARACTERISTICS OF HFE-7100 IN MANIFOLD MICROCHANNEL HEAT SINK WITH CORRUGATED BOTTOM

Jianping Cheng¹, Hongsen Xu¹, Zhiguo Tang^{1*}, Pei Zhou²

¹School of Mechanical Engineering, Hefei University of Technology, Hefei, China ²School of Automotive and Transportation Engineering, Hefei University of Technology, Hefei, China

* Corresponding author; E-mail: 2017217349@mail.hfut.edu.cn

To solve the problem of high heat flux heat dissipation in microelectronic devices, a manifold microchannel heat sink with corrugated bottom (CB-MMC) is proposed on the basis of the manifold microchannel heat sink (MMC). The flow and heat transfer characteristics of HFE-7100 in MMC and CB-MMC are investigated numerically. The results show that CB-MMC reduces the pressure loss and enhances the heat transfer performance in single-phase flow. The orthogonal test method is used to obtain structural design solutions with optimal thermal performance. It is observed that the temperature reduction is always at the expense of the increase of the pressure drop. In addition, the optimization parameters combination obtained through comprehensive evaluation of temperature and pressure drop through weight matrix - optimized solution 19 (wavelength $A=800 \mu m$, amplitude $B=40~\mu\text{m}$, channel depth $C=180~\mu\text{m}$, outlet width $D=300~\mu\text{m}$, channel width $E=25 \mu m$). Its T_{ave} has decreased by 6.89°C, ΔP decreased by 10.27 kPa. Moreover, the subcooled boiling flow and heat transfer performance in MMC and CB-MMC are comparatively studied. The results demonstrate that the dynamic behavior of vapor bubbles causes large pressure fluctuations, which further leads to small temperature fluctuations, and thus reduces the stability of the flow and boiling heat transfer. Compared with MMC, CB-MMC exhibits more stable two-phase flow and better boiling heat transfer.

Key words: Manifold microchannel; Corrugated bottom; Orthogonal test; Flow boiling; Enhanced heat transfer

1. Introduction

With the development of the electronic technology, highly integrated and miniaturized electronic devices have been widely used, which led to the increasing heat flux of electronic devices. Some high-power electronic devices dissipate heat flux greater than 1000 W/cm² [1], and thus they work at higher temperature, which leads to thermal failure. Therefore, it is critical to select a suitable heat dissipation technology for high heat flux microelectronic devices. In 1981, Tuckerman and Pease [2] first proposed a microchannel liquid cooling method, which integrates microchannels on the back side of the electronic chip and uses water as coolant for forced convection heat transfer. Their

experiment proved the high efficiency heat transfer capability of the microchannel heat sink (TMC). Many studies on this technology were then conducted [3, 4].

The studies on TMC mainly focused on the design of sectional shapes [5-7], geometric parameters [8-10], and disturbance structures in the microchannel [11]. In the study of microchannel sectional shapes, it is considered that the rectangular section has smaller flow resistance [7] and higher heat transfer performance than the trapezoidal and triangular sections [6]. In the study of the geometrical parameters of rectangular microchannels, the ratio of depth to width [8] and the hydraulic diameter [10] of these microchannels are the main factors affecting the friction coefficient and Nusselt number. To obtain higher flow and heat transfer performance, many researchers designed different disturbance structures within the microchannels, including ribs [12], fins [13], and corrugated structures [14]. For instance, Wan et al. [14] introduced the semi-corrugated structure into the microchannel, which resulted in lower flow resistance and thermal resistance.

However, Kermani et al. [15] found two main disadvantages of the conventional TMC: the high pressure drop and large temperature difference due to the long channels. Therefore, Harpole and Eninger [16] proposed the manifold microchannel heat sink (MMC), which adds a distributor structure to the TMC. The new design was found to significantly reduce the pressure drop and temperature difference due to the shortened flow path of the coolant. Many studies on MMC were then conducted, while mainly focus on the design of inlet and outlet flow channels [17-19], geometric parameters of manifold microchannels [20, 21], and channel bottom structure [22]. In order to optimize the manifold inlet and outlet channels, Tang et al. [18] changed the intake manifold channel and the intake plenum chamber into a conical contraction structure based on the self-similar heat sink (SSHS), to alleviate the uneven flow distribution inside MMC. In the studies on the geometric parameters of MMC, there are mainly the inlet/outlet ratios of microchannel [20] and aspect ratios of channel section [21]. For instance, Pourfattah et al. [20] numerically found that the optimal heat transfer and flow performance of MMC can be obtained at the inlet/outlet ratio of 0.25, Reynolds number (Re) of 100, and nanoparticle volume fraction of 2%. Pan et al. [21] demonstrated that there is an optimum aspect ratio that allows the MMC to achieve optimal heat transfer. Furthermore, they determined the functional relationship between the optimal aspect ratio and the Prandtl number (Pr), Re number, and the ratio of the thermal conductivity between channel wall and coolant. In addition, many studies on adding enhanced heat transfer structures to the inner wall of MMC were conducted, including added ribs [23], fins [24, 25] and porous media [26]. For instance, Adio et al. [23] introduced four rib structures in MMC: forward triangular, rectangular, backward triangular, and semicircular ribs. They found that the overall heat transfer characteristics of all the sidewall ribbed MMC are higher than that of the smooth MMC, while the rectangular ribbed MMC obtained the largest performance enhancement. Pan et al. [25] introduced staggered pin-fin in MMC to obtain higher heat transfer performance and better temperature uniformity. Chen et al. [26] introduced porous fins in MMC, and showed that the MMC with 75% proportion of porous fins has the highest performance with a 19.8% reduction in thermal resistance and 46.2% increase in performance evaluation criterion (PEC).

The above studies mainly focus on the single-phase heat transfer in MMC. As we all know that the boiling heat transfer can effectively enhance the heat transfer performance [27]. Drummond et al. [28] manufactured and experimentally studied an embedded MMC that achieves heat flux dissipation of 1020 W/cm² under two-phase flow conditions using HFE-7100 as the coolant. The studies on boiling heat transfer in MMC mainly focus on the vapor bubbles dynamics [29], two-phase flow

distribution and evolution [30, 31], critical heat flux (CHF) [29], and boiling heat transfer enhancement [32]. In fact, the vapor bubble growth is the basic problem in the flow boiling process. Xie et al. [29] deduced that the *Re* number significantly affects the growth, distribution, and evolution of bubbles. Mukherjee et al. [30] deduced that the bubble growth rate decreases with increasing the *Re* number, and increases with increasing the superheat temperature of coolant. Lin et al. [31] studied the two-phase flow distribution in MMC, and explored U-shaped and HU-shaped manifold structures with low pressure drop and low thermal resistance. Xie et al. [29] also studied the mechanism of CHF in MMC. They deduced that CHF usually occurs in a channel having large aspect ratio and small flow rate. They also provided suggestions for the optimization of the flow field to delay CHF. In addition, the geometric parameters of MMC significantly affect the enhanced boiling heat transfer. Luo et al. [32] considered that when the ratio of inlet to outlet width is between 1 and 2, the heat sink shows the minimum pressure loss.

Therefore, based on the above literature, a novel manifold microchannel heat sink with corrugated bottom (CB-MMC) is proposed in this study in order to further reduce the thermal resistance and flow resistance by introducing a corrugated structure into the channel bottom of manifold microchannel heat sink (MMC). The HFE-7100, which has low boiling point and hydrophilic properties, is used as the coolant in this study, and the flow and heat transfer characteristics in MMC and CB-MMC are numerically simulated under single-phase and two-phase subcooled boiling conditions. The CB-MMC structure with the best flow and heat transfer characteristics is determined using the orthogonal test method. Finally, the two-phase flow and boiling heat transfer characteristics of MMC and CB-MMC are compared.

2. Physical models and numerical methods

2.1. Geometric model

Fig. 1(a) shows the structure diagram of MMC. Due to the advantages of the corrugated structure in terms of reduced flow resistance and enhanced heat transfer performance [33], a novel CB-MMC is proposed by applying corrugated structure to the channel bottom of the MMC, as shown in **Fig. 1**(b). The coolant enters the microchannel from the manifold inlet, and it is divided into two branches in the microchannel for fluid-solid coupled heat transfer. Finally, the coolant converges with the fluid entering from the adjacent manifold inlet and flows out. Due to the strict geometric symmetry, the single-cell model shown in **Figs. 1**(c) and (d) is used as the computational domain for simulation to reduce the computational cost. The geometric parameter values are shown in **Table 1**. The heat sink structures are all made of silicon, and HFE-7100 is used as coolant. The thermophysical properties of silicon and HFE-7100 are shown in **Table 2**.

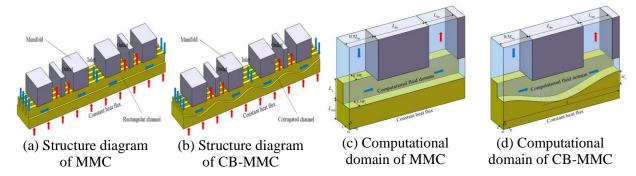


Fig. 1. Geometric model

Table 1. Geometric parameter values

Parameters	Variables	Values (µm)
Length of inlet	$L_{ m in}$	400
Length of outlet	$L_{ m out}$	200
Length of divider	$L_{ m div}$	350
Heightof substrate	$L_{ m sub}$	150
Height of channel	$L_{ m c}$	150
Width of fin	$W_{ m f}$	15
Width of channel	$W_{ m c}$	15
Wavelength	λ	800
Amplitude	$H_{ m a}$	100

Table 2. Thermophysical properties of silicon and HFE-7100

	Ciliaan	HFE-7100	HFE-7100	HFE-7100
	Silicon	(single-phase)	(liquid)	(vapor)
ho, kg·m ⁻³	2330	1420	1402	11.5
$C_{\rm p}$, J·kg ⁻¹ ·K ⁻¹	712	1500	1263	870
$k, \mathbf{W} \cdot \mathbf{m}^{-1} \cdot \mathbf{K}^{-1}$	148	0.059	0.058	0.01
μ , kg·m ⁻¹ ·s ⁻¹		4.34×10^{-4}	4.00×10^{-4}	1.32×10^{-5}
$i_{\rm lv},{\rm kJ}\cdot{\rm kg}^{-1}$			111.6	
σ , N·m ⁻¹			9.43×10 ⁻³	

^{*} The saturation temperature of liquid HFE-7100 $T_{\text{sat}} = T_{\text{ref}} = 339$ K.

2.2. Governing equations

To correctly describe the simulation process, the following assumptions are made: (1) the fluid is incompressible and Newtonian, (2) the fluid and solid have constant thermophysical properties, and (3) the viscous dissipation and contact thermal resistance are neglected.

When the single-phase heat transfer occurs in the channel, it can be treated as a fluid-solid coupled heat transfer problem, and the flow can be considered as a steady-state laminar flow [24]. Its continuity, momentum, and energy equations are given by [21]:

$$\nabla \cdot \boldsymbol{u} = 0 \tag{1}$$

$$(\boldsymbol{u}\cdot\nabla)\rho_{\scriptscriptstyle{f}}\boldsymbol{u} = -\nabla P + \mu\nabla^{2}\boldsymbol{u} \tag{2}$$

$$\rho_{\rm f} c_{\rm p,f} (\boldsymbol{u} \cdot \nabla T) = k_{\rm f} \nabla^2 T \tag{3}$$

where u denotes the velocity vector of the fluid.

Energy equation in the solid domain of heat sink:

$$k_{s}\nabla^{2}T = 0 \tag{4}$$

In subcooled flow boiling, evaporation processes of the coolant occur in the microchannel. The volume of fluid (VOF) [32] method is adopted to calculate the subcooled flow boiling in MMC. The governing equations are given by [24]:

Fluid volume fraction equation:

$$\frac{\partial \dot{\alpha}}{\partial t} + \nabla \cdot (\alpha \mathbf{u}) = \dot{\alpha}_{lv} \tag{5}$$

where α is the volume fraction of the liquid phase and $\dot{\alpha}_{\rm lv}$ is the volume fraction source term caused by phase change.

The continuity, momentum, and energy equations are as follows:

$$\nabla \cdot \boldsymbol{u} = \dot{\mathbf{v}}.\tag{6}$$

$$\frac{\partial(\rho \boldsymbol{u})}{\partial t} + \nabla \cdot (\rho \boldsymbol{u} \boldsymbol{u}) = \boldsymbol{f}_{\sigma} - \nabla p + \nabla \cdot [\mu(\nabla \boldsymbol{u} + \nabla \boldsymbol{u}^{T})] \qquad (7)$$

$$\frac{\partial(\rho i)}{\partial t} + \nabla \cdot (\rho \boldsymbol{u} i) = \nabla \cdot (k \nabla T) - \dot{q}_{l_{\text{to}}} \qquad (8)$$

$$\frac{\partial(\rho i)}{\partial t} + \nabla \cdot (\rho u i) = \nabla \cdot (k \nabla T) - \dot{q}_{_{lv}}$$
(8)

Energy equation of the heat sink solid domain:

$$\frac{\partial(\rho c_{p,s})}{\partial t} = k_s \nabla^2 T \tag{9}$$

where:

$$f_{\sigma} = \sigma \kappa \boldsymbol{n} \left| \nabla \alpha \right| \frac{2\rho}{\rho_{v} + \rho_{1}}$$

$$\kappa = -(\nabla \cdot \boldsymbol{n})$$
(10)

$$\kappa = -(\nabla \cdot \boldsymbol{n}) \tag{11}$$

$$\boldsymbol{n} = \frac{\nabla \alpha}{|\nabla \alpha|} \tag{12}$$

where f_{σ} denotes the surface tension, σ is the surface tension coefficient, κ is defined according to the dispersion of the unit normal, and n denotes the unit normal.

The enthalpy value i of intermediate state is calculated by mass-averaged:

$$i = \frac{(1-\alpha)\rho_{v}c_{p,v} + \alpha\rho_{l}c_{p,l}}{\rho}(T - T_{ref})$$
(13)

The density, viscosity, and thermal conductivity in the equations are calculated by volume fraction weighted average:

$$\rho = (1 - \alpha)\rho_{v} + \alpha\rho_{l} \tag{14}$$

$$\mu = (1 - \alpha)\mu_{v} + \alpha\mu_{1} \tag{15}$$

$$k = (1 - \alpha)k_{v} + \alpha k_{l} \tag{16}$$

The empirical rate parameter model is used in this study to describe the phase change process, which is based on the Lee model [34] proposed by Yang et al. [35]. This model does not require to artificially set nucleation points. The mass transfer process is driven by the difference between the local and saturation temperatures. Thus the mass transfer equation is as follows:

$$\dot{m}_{lv} = \begin{cases} \gamma_{l} \alpha \rho_{l} \frac{T - T_{sat}}{T_{sat}} & \text{if} \quad T \ge T_{sat} \\ \gamma_{v} (1 - \alpha) \rho_{v} \frac{T - T_{sat}}{T_{sat}} & \text{if} \quad T \le T_{sat} \end{cases}$$

$$(17)$$

The evaporation and condensation empirical rate coefficients (γ_1 and γ_y) highly depend on the simulated operating conditions, flowing geometry, and grid size. γ should have a compromise value to ensure that the vapor-liquid interface temperature is near the saturation temperature [36]. Therefore, based on the research results of Luo et al. [32], it is assumed that $\gamma_1 = \gamma_v = 100$.

The source term in the energy equation is determined by Eqs. (13) and (17):

$$\dot{q}_{\rm lv} = \dot{m}_{\rm lv} \, \dot{i}_{\rm lv} \tag{18}$$

Thus, the volume evaporation rate and the phase fraction source term can be expressed as follows:

$$\dot{v}_{lv} = \frac{\dot{q}_{lv}}{i_{lv}} (\frac{1}{\rho_{v}} - \frac{1}{\rho_{l}}) \tag{19}$$

$$\dot{\alpha}_{\rm lv} = -\frac{\dot{q}_{\rm lv}}{\rho \dot{l}_{\rm lv}} \tag{20}$$

2.3. Boundary conditions

In this study, the coolant enters the manifold at a constant velocity and temperature. An atmospheric pressure is specified at the manifold outlet. The bottom heat source surface is defined with a constant heat flux. Because it is far away from the heat source surface, the adiabatic boundary condition is applied to the upper solid surface. The non-slip boundary condition is applied to all fluid-solid coupled surfaces. Due to the geometrical and flow periodicity, the four sidewall surfaces of the calculation model are given symmetric boundary conditions, and the symmetry plane is set as zero gradient:

$$\frac{\partial u}{\partial n} = 0 \tag{21}$$

The coolant enters the heat sink at a constant velocity perpendicular to the inlet section, and its flow rate is computed as:

$$u = \frac{G_{\rm in}}{\rho_{\rm l}} \tag{22}$$

The fluid-solid coupled surface is continuously coupled by temperature and heat flux:

$$T_{\rm s} = T_{\rm f} \tag{23}$$

$$k_{\rm f} \frac{\partial T_{\rm f}}{\partial n} = k_{\rm s} \frac{\partial T_{\rm s}}{\partial n} \tag{24}$$

In the numerical calculation, the second-order upwind formats are used to discretize momentum and energy. The velocity-pressure coupling is solved using the SIMPLE algorithm, with a time step of 10^{-6} s. Commercial ANSYS FLUENT software is used to solve three-dimensional heat transfer and flow equations.

2.4. Verification of grid independence

Since the computational domain of CB-MMC is a rectangular structure, the structured mesh division of 3D model can improve the computational speed and save computational resources. **Fig. 2** shows the hexahedral mesh of the CB-MMC computational domain and the local mesh refinement.

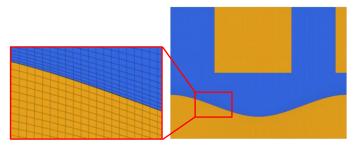


Fig. 2. CB-MMC computational domain mesh

The impact of the grid numbers is explored in this section. Five sets of grids with different grid numbers are considered. The number of grids ranging between 100,000 and 500,000. Simulations are performed for each of the five sets of grids with a microchannel flow rate of 1300 kg/(m²·s) and a heat flux at the bottom heat source surface of 200 W/cm². **Fig. 3** shows the impact of the grid numbers on the average temperature and the pressure drop of the CB-MMC. It is found that the trend of the average temperature and the pressure drop tends to stabilize as the grid numbers increase. The errors of the relative changes of the average temperature and the pressure drop in CB-MMC become less than 1% as the grid numbers increase from 300,000 to 400,000. Therefore, considering the efficiency of numerical calculation, a grid number of 300,000 is used for subsequent simulation studies of 3D model.

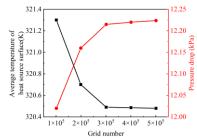


Fig. 3. The impact of the grid numbers on the calculation results

2.5. Validation of the simulation model

In this section, an experimental validation of the simulation model is performed. The structure of the calculation model is similar to that of the MMC experimentally studied in [28], as shown in **Fig.** 1(c). The geometric parameters of the calculation model are shown in **Table 1**. The HFE-7100 is used as the coolant. The inlet temperature is 332.15 K, and the channel mass flow rate is 1300 kg/($m^2 \cdot s$). The heat flux in the range of 0-400 W/cm² is set on the bottom of CB-MMC.

Fig. 4 shows the comparison between numerical simulation and experimental results. It is found that when the heat flux is less than 100 W/cm², almost no boiling occurs. At this time, single-phase heat transfer occurs in the channel, and the simulation results are consistent with the experimental data. When the heat flux increases from 150 W/cm² to 400 W/cm², flow boiling heat transfer occurs in the channel, the simulation results are slightly higher than the experimental data, and the difference between the average temperatures of the corresponding heat source surfaces is less than 3 K. Therefore, the validation results show that the simulation model is reliable.

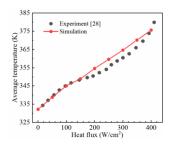


Fig. 4. Validation of the numerical model

3. Results and discussion

3.1. Comparison of single-phase flow and heat transfer characteristics of MMC and CB-MMC

In this section, MMC and CB-MMC, having the same inlet width, divider width, outlet width, and channel width are compared (**Figs. 1**(c) and (d)). The microchannel depth of MMC is consistent with the corrugated centerline depth of CB-MMC. Therefore, the two heat sinks have the same inlet velocity when the channel mass flows are equal. The specific structural dimensions are shown in **Table 1**. The bottom heat source surface of MMC and CB-MMC adopts a constant heat flux of 200 W/cm².

Fig. 5 shows the pressure nephogram, streamline, and temperature nephogram in the MMC and CB-MMC channels for a mass flow rate of 1500 kg/(m²·s) in the single-phase flow. It can be seen from Fig. 5(a) that in MMC, when the coolant flows into the microchannel through the manifold channel, the flow rate at the inlet increases and the pressure sharply decreases due to the sudden contraction of the section and the change of the flow direction. In CB-MMC, the corrugated bottom connects the inlet and outlet channels of the manifold into a similar U-shaped structure, which reduces the flow resistance when the coolant turns. It can be observed from Fig. 5(b) that the temperature distributions of MMC and CB-MMC are non-uniform, with a higher temperature of heat source surface near the manifold outlet. This is because the coolant collides with the incoming coolant from the adjacent inlet manifold at the outlet manifold where it accumulates, so that the heat cannot be discharged in time, which results in increasing the temperature in this area. In addition, the maximum temperature of the heat source surface of CB-MMC is lower than that of MMC, and its temperature uniformity is also better. This is mainly because the corrugated channel has a larger heat transfer area than the rectangular channel, and therefore the temperature of CB-MMC at the inlet is lower. In addition, the coolant impinges on the corrugated surface at the outlet, which destroys the thermal boundary layer, and thus enhances the convective heat transfer at the outlet.

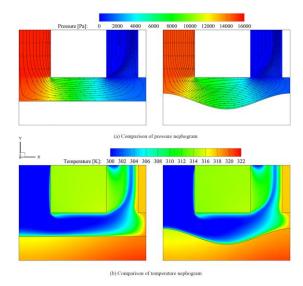


Fig. 5. Comparison of pressure and temperature nephograms between MMC and CB-MMC

Fig. 6 shows the variation of the average temperature of heat source surface (T_{ave}) and the pressure drop between inlet and outlet (ΔP) in MMC and CB-MMC. As can be seen from **Fig. 6**(a), all the T_{ave} values decrease as the channel mass flow rate increases. More precisely, when the flow rate increases from 1000 kg/(m²·s) to 1500 kg/(m²·s), the T_{ave} values of MMC and CB-MMC decrease by 4.3 °C and 3.5 °C, respectively. This is mainly due to the thinning of the thermal boundary layer caused by the coolant at high velocity. In addition, under the same conditions, CB-MMC increases the heat transfer area. Thus, it has a lower T_{ave} values than that of MMC. Fig. 6(b) shows that all the ΔP values increase as the channel flow rate increases, while the ΔP of CB-MMC is always lower than that of MMC. Under the same geometric conditions, the increase in flow velocity leads to greater disturbance to the bottom layer of the laminar flow [37], and thus enhances the heat transfer while also increasing the pressure loss. That is, the corrugated bottom plays a buffering role on the collision of the fluid, slows down the reduction of the flow velocity, and thus reduces the pressure loss, as shown in Fig. 5. Fig. 6(a) and (b) shows that at a channel mass flow rate of 1500 kg/(m²·s), the CB-MMC reduces the pressure drop by 4.1% at a temperature reduction of less than 1% compared to the MMC. It can be observed from Fig. 6(c) that a greater ΔP is required for MMC to reach the same T_{ave} obtained by CB-MMC. At both $T_{\rm ave}$ of 48 °C, the pressure drop of the MMC increases by 9.3% compared to the CB-MMC. A lower temperature and a lower pressure drop are both often required, and thus the corresponding structure near the lower left corner of the coordinate is superior. It is deduced that the comprehensive heat transfer performance of CB-MMC is higher than that of MMC.

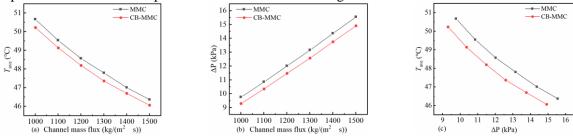


Fig. 6. Comparison of T_{ave} and ΔP between MMC and CB-MMC

3.2. The orthogonal test

In this section, a multi-factor analysis of the structural parameters affecting the flow and heat transfer of the CB-MMC is performed using the orthogonal test method, so that the influence of each factor level on the test index values can be determined. Thus, the structure design of CB-MMC with the best flow and heat transfer characteristics can be determined. The core of orthogonal test is the orthogonal array which consists of various factors and their corresponding levels [38]. Therefore, the factors and levels should be first determined. The corrugated structure wavelength (A), corrugated structure amplitude (B), channel depth (C), outlet width (D), and channel width (E) are considered as the influencing factors, and T_{ave} and ΔP are used as the test indexes to establish a five-factor four-level orthogonal test, as shown in **Table 3**.

Table 3. Factors and levels of the test

Levels	Factors				
	$A/\mu m$	$B/\mu m$	C/µm	$D/\mu m$	E/µm
1	400	40	90	150	10
2	800	80	120	200	15
3	1200	120	150	250	20
4	1600	160	180	300	25

The orthogonal design table for arranging various structural factors is established according to the L_{16} (4⁵) orthogonal table, as shown in **Table 4**. In the table, each row represents a set of CB-MMC structural parameter combinations. The computational geometry model corresponding to each scheme in **Table 4** is then developed, and CFD simulations are performed for all the models using ANSYS Fluent. The evaluation index values obtained from the simulation are shown in **Table 4**, where I and II are the evaluation indexes of T_{ave} and ΔP , respectively.

Table 4. L_{16} (4⁵) orthogonal designed table with the obtained results

T	Factors					Evaluation	on index
Test no.	$A/\mu m$	$B/\mu m$	C/µm	$D/\mu m$	$E/\mu m$	I (°C)	II (kPa)
1	400	40	90	150	10	58.93	30.22
2	400	80	120	200	15	50.71	16.08
3	400	120	150	250	20	49.42	12.28
4	400	160	180	300	25	51.80	10.91
5	800	40	120	250	25	54.53	5.85
6	800	80	90	300	20	59.77	6.18
7	800	120	180	150	15	42.43	18.08
8	800	160	150	200	10	46.03	27.37
9	1200	40	150	300	15	47.01	12.53
10	1200	80	180	250	10	43.72	28.58
11	1200	120	90	200	25	59.52	7.11
12	1200	160	120	150	20	49.60	10.76
13	1600	40	180	200	20	44.53	10.93
14	1600	80	150	150	25	50.33	8.88
15	1600	120	120	300	10	54.96	26.98

16 1600 160 90 250 15 59.71 22.33

The range of the evaluation indexes is obtained based on the evaluation index values in **Table 4**, as shown in **Tables 5** and **6**. In these tables, K_t and k_t (t=1, 2, 3, and 4) represent the sum and the average of the evaluation index values at different levels of a factor. Therefore, $k_t = K_t/4$. R_I and R_{II} respectively denote the ranges of the evaluation indexes I and II under a certain factor, and their values are calculated as follows: $R_{I/II} = \max(k_t) - \min(k_t)$, (t = 1, 2, 3, and 4), that is, the difference between the maximum value and the minimum value of the evaluation indexes under the corresponding factor. The value reflects the range of change of the evaluation index when each factor fluctuates at different levels. The larger the value, the greater the influence of the factor on the evaluation indexes. Therefore, the order of the effects of the factors on the evaluation indexes can be determined from the range value [39].

Table 5. Range analysis of the evaluation index I

Index	Factors				
	A/μm	B/µm	C/µm	D/µm	E/µm
K_1	210.86	205.00	237.93	201.29	203.64
K_2	202.76	204.53	209.80	200.79	199.86
K_3	199.85	206.33	192.79	207.38	203.32
K_4	209.53	207.14	182.48	213.54	216.18
k_1	52.72	51.25	59.48	50.32	50.91
k_2	50.69	<u>51.13</u>	52.45	50.20	<u>49.97</u>
k_3	<u>49.96</u>	51.58	48.20	51.85	50.83
k_4	52.38	51.79	<u>45.62</u>	53.39	54.05
$R_{\rm I}$	2.76	0.66	13.86	3.19	4.08
Optimize selections	A_3	B_2	C_4	D_2	E_2

Table 6. Range analysis of the evaluation index II

Index	Factors				
	$A/\mu m$	$B/\mu m$	C/µm	$D/\mu m$	E/µm
K_1	69.49	59.53	65.84	67.94	113.15
K_2	57.48	59.72	59.67	61.49	69.02
K_3	58.98	64.45	61.06	69.04	40.15
K_4	69.12	71.37	68.50	56.60	32.75
k_1	17.37	<u>14.88</u>	16.46	16.99	28.29
k_2	14.37	14.93	14.92	15.37	17.26
k_3	14.75	16.11	15.27	17.26	10.04
k_4	17.28	17.84	17.13	<u>14.15</u>	<u>8.19</u>
$R_{ { m II}}$	3.00	2.96	2.21	3.11	20.10
Optimize selections	A_2	B_1	C_2	D_4	E_4

It can be deduced from **Table 5** that the primary and secondary order of the impact of each factor on evaluation index I is C > E > D > A > B. This indicates that factor C has the most significant

effect on T_{ave} , while factor B has the least effect on it. Similarly, it can be seen from the range values in **Table 6** that the order of the influence of each factor on evaluation index II is E > D > A > B > C. Among the structural factors, factor E has the greatest effect on ΔP in heat sink, while factor C has the least effect on it. For evaluation indexes I and II, lower T_{ave} and ΔP are considered as more preferable selections. Therefore, the optimal solution for I obtained from **Table 5** is $A_3B_2C_4D_2E_2$ (Define it as Solution 17), and that for II obtained from **Table 6** is $A_2B_1C_2D_4E_4$ (Define it as Solution 18). The values of each structural parameter of the selected optimal solution are underlined in the table.

The above two sets of solutions are the superior results obtained from individual evaluation indexes. To meet the requirements of two evaluation indexes at the same time, a more optimal combination solution is obtained relative to I and II. In this study, the weight matrix analysis method is used to comprehensively measure the effect weight of each factor level on two evaluation indexes, so as to rapidly determine the optimal solution. The weight matrices of the two evaluation indexes are calculated as follows:

The evaluation index matrix is first established, where k_{ij} is the average of the evaluation indexes at the j-th level of the i-th factor. If the evaluation index is as small as possible, then K_{ij} is set to $1/k_{ij}$ and matrix M is constructed as:

$$M = \begin{bmatrix} K_{11} & 0 & 0 & \cdots & 0 \\ K_{12} & 0 & 0 & \cdots & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ K_{1j} & 0 & 0 & \cdots & 0 \\ 0 & K_{21} & 0 & \cdots & 0 \\ 0 & K_{22} & 0 & \cdots & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & K_{2j} & 0 & \cdots & 0 \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & 0 & 0 & \cdots & K_{1i} \\ 0 & 0 & 0 & 0 & \cdots & K_{1i} \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ 0 & 0 & 0 & 0 & \cdots & K_{1i} \\ 0 & 0 & 0 & \cdots & K_{1i} \end{bmatrix}$$

$$(25)$$

By setting $_{T_i=1/\sum K_{ij}}$, the factor matrix T is then constructed as:

$$T = \begin{bmatrix} T_1 & \cdots \\ T_2 & \cdots \\ \cdots & \cdots & \cdots \\ \cdots & \cdots & T_i \end{bmatrix}$$
 (26)

Afterwards, the level matrix is established and the range is defined as s_i . Finally, by setting $s_i = s_i / \sum_{i=1}^{i} s_i$, the S matrix is constructed as:

$$S = \begin{bmatrix} S_1 \\ S_2 \\ \dots \\ S_i \end{bmatrix}$$
 (27)

Therefore, the weight matrix of the evaluation indexes is calculated as:

$$\omega = MTS \tag{28}$$

There are two evaluation indexes in this test. Therefore, the total weight matrix of the orthogonal test is the average of the weight matrix of the two evaluation indexes, which is calculated as:

$$\omega = \frac{\omega_{l} + \omega_{ll}}{2} = \frac{1}{2} \begin{bmatrix} 0.027409 \\ 0.028507 \\ 0.027582 \\ 0.006744 \\ 0.006762 \\ 0.006765 \\ 0.1208188 \\ 0.006675 \\ 0.1208188 \\ 0.017002 \\ 0.006752 \\ 0.023206 \\ 0.017002 \\ 0.01832 \\ 0.018329 \\ 0.033182 \\ 0.033182 \\ 0.033182 \\ 0.0331278 \\ 0.0331278 \\ 0.032213 \\ 0.042073 \\ 0.042733 \\ 0.042733 \\ 0.042073 \\ 0.042073 \\ 0.042073 \\ 0.042073 \\ 0.042073 \\ 0.042073 \\ 0.039508 \end{bmatrix} \begin{bmatrix} 0.021769 \\ 0.02459 \\ 0.02748 \\ 0.02537 \\ 0.01593 \\ 0.01495 \\ 0.07184 \\ 0.02814 \\ 0.02940 \\ 0.02814 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02940 \\ 0.02951 \\ 0.06691 \\ 0.05691 \\ 0.05691 \\ 0.05691 \\ 0.05691 \\ 0.05691 \\ 0.05691 \\ 0.05691 \\ 0.14390 \end{bmatrix}$$

The comprehensive impact weight of each factor on the evaluation index is calculated using **Eq.** (29), as shown in **Table 7**. It can be deduced that the order of the effect of each factor on the evaluation index is E > C > D > A > B. Moreover, by comparing the weights of each level for each factor, it can be seen that at different levels of the five factors, the levels that have the greatest impact weight on the results are A_2 , B_1 , C_4 , D_4 , and E_4 . Therefore, the optimal solution obtained using the weight matrix method while considering evaluation indexes I and II is $A_2B_1C_4D_4E_4$ (Define it as Solution 19).

Table 7. Analysis results obtained from the weight matrix.

Factors	A	В	С	D	Е
Weights	0.10401	0.06060	0.31749	0.11453	0.40336
Weight order	E>C>D>.	A>B			
Superior levels	A_2	B_1	C_4	D_4	E_4
Optimal scheme	$A_2B_1C_4D_4$	E_4			

In order to verify the performance advantages of the optimal solution obtained using the weight matrix method, the evaluation index values obtained from the three optimal solutions are compared with those of the 16 groups of orthogonal test solutions, as shown in **Fig. 7**. It can be observed from **Fig. 7** that in the same case, the lowest T_{ave} is obtained in solution 7, while T_{ave} in solution 17 is only 0.32 °C higher than that in solution 7. However, its ΔP is 1.87 kPa lower than that in solution 7. Therefore, solution 17 is superior when only T_{ave} is considered. Similarly, when only considering ΔP , solution 18 has the lowest ΔP while its T_{ave} is relatively high. Solution 19 is obtained by comprehensively considering the temperature and pressure drop. Its T_{ave} is 6.89 °C lower than that in solution 18, and its ΔP is 10.27 kPa lower than that in solution 17. **Fig. 8** shows the comparison of temperature and pressure nephograms, and it can be seen that solution 19 has moderate T_{ave} and ΔP .

To further demonstrate the performance advantages of the optimized solutions, the performance evaluation criteria (PEC) for solutions 17, 18 and 19 are compared using the conventional MMC as a benchmark. The PEC is expressed as follows:

$$PEC = \frac{Nu}{Nu_0} / \left(\frac{f}{f_0}\right)^{1/3} \tag{30}$$

Where Nu_0 and f_0 represent the Nusselt number and friction factor of MMC, respectively. The comparison results are shown in **Table 8**. It can be seen that the PEC of the three optimized solutions is higher than that of MMC. It is worth noting that optimized solution 19 achieved the highest PEC, which increased by 67.5% compared to MMC. Therefore, The solution 19 obtained using the weight

matrix analysis method has the optimal comprehensive heat transfer performance.

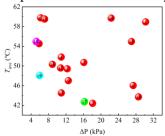


Fig. 7. Comparison of optimal solutions

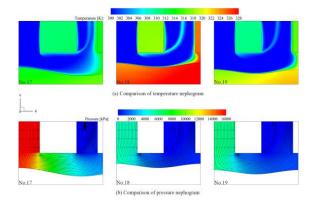


Fig. 8. Comparison of temperature and pressure nephograms of solutions 17, 18, and 19

Table 8. Comparison of performance evaluation criteria (PEC)

					•
solutions	MMC	17	18	19	
Nu	4.435	4.429	3.519	3.921	
f	0.927	0.664	0.266	0.145	
PEC	1	1.139	1.227	1.675	

3.3. Comparison of two-phase boiling flow and heat transfer characteristics of MMC and CB-MMC

The subcooled boiling flow and heat transfer characteristics of MMC and CB-MMC in twophase are numerically studied in this section. They both adopt the geometrical parameters of the optimal solution 19 ($A_2B_1C_4D_4E_4$). The boundary conditions are as follows: the channel mass flow rate is set at 1100 kg/(m²·s) and the boiling coolant enters the heat sink at 59 °C. The heat flux of 200 W/cm², 300 W/cm², and 400 W/cm² are set at the bottom of the heat sink. The thermophysical properties of the HFE-7100 during boiling are shown in **Table 2**.

Fig. 9 shows the variation of the T_{ave} and ΔP of MMC and optimized CB-MMC under subcooled flow boiling for three heat flux conditions. **Fig. 9**(a) and (b) shows that T_{ave} first rapidly increases with time, and then tends to be basically stable at almost 15 ms with a slight fluctuation. When the heat flux increases, the T_{ave} values in MMC and CB-MMC equally increase. In contrast to the single-phase flow and heat transfer characteristics presented in the previous study, the T_{ave} values of MMC and CB-MMC are basically the same. It can be observed from **Fig. 9**(c) and (d) that ΔP increases with time and reaches dynamic fluctuation equilibrium at almost 10 ms. It then oscillates at a specific value. This is basically consistent with the results of Luo et al. [40]. Moreover, it is shown that the average values of ΔP fluctuation increase with the heat flux increasing, and the fluctuation amplitudes become larger. However, CB-MMC shows a more stable fluctuation trend of ΔP compared

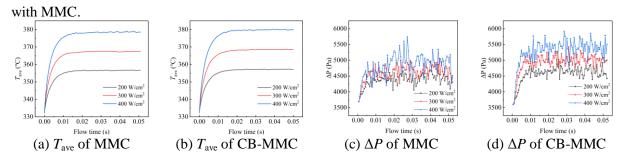


Fig. 9. The change of T_{ave} and ΔP for MMC and CB-MMC under subcooled flow boiling

Figs. 10 and 11 show the generation, growth, and movement of bubbles in MMC and CB-MMC. More precisely, Fig. 10 illustrates the variation of the vapor phase nephograms in MMC and CB-MMC with flow time for the heat flux of 400 W/cm². It can be seen that at 0.002 s, the bubbles are first generated at the bottom of the channel, and the isolated bubbles are mainly concentrated near the downstream outlet. This can be considered as the starting point of nuclear boiling [41]. Due to the hydrophilicity of HFE-7100, the bubbles quickly separate the channel wall after nucleation and flow out of the microchannel. The periodic generation and outflow of bubbles cause the pressure fluctuation in the channel, as shown in Fig. 9(c) and (d). At this time, the phase transition only locally occurs, and the heat taken away by small bubbles is limited. Thus, the temperature of MMC and CB-MMC continues to increase, as shown in Fig. 9(a) and (b). With the further occurrence of boiling, more nucleation sites are activated, the bubble generation area extends to the whole channel bottom, and more bubbles converge into large bubbles to form membrane boiling. From 0.004 s to 0.008 s, the vapor region changes from nuclear boiling to membrane boiling. From 0.008 s to 0.010 s, the volume fraction of vapor in the channel is basically unchanged and the boiling reaches a steady state. Afterwards, nuclear boiling and local membrane boiling alternately occur, and then exit the channel. This dynamic behavior of bubbles results in pressure fluctuations after boiling stabilization. In addition, the periodic removal of heat by bubbles causes T_{ave} to fluctuate in a similar way to the pressure drop. This is consistent with the results obtained by Pan et al. [24]. By comparing the vapor phase nephograms of MMC and CB-MMC, it can be deduced that when the boiling reaches a stable state, more bubbles are generated in the CB-MMC to ensure sufficient heat exchange, thus showing more stable temperature fluctuations.

Fig. 11 shows a comparison between the vapor phase nephograms of MMC and CB-MMC under different heat fluxes at a fixed time. It can be seen that when the heat flux increases, the proportions of vapor phase in the channel increase, and thus the ΔP values of MMC and CB-MMC both increase (Fig. 9(c) and (d)). Moreover, the addition of corrugated bottom increases the area of bubble nucleation at the bottom, which leads to more bubbles, thus taking more heat away. More nucleation points result in more bubbles growing, gathering, and discharging, which results in a higher pressure drop in the CB-MMC. As mentioned in Section 3.1, the corrugated bottom surface and the inlet and outlet channels in CB-MMC are connected into the U-shaped channel, which slows down the reduction of the flow velocity after the coolant enters the microchannel, and thus the bubbles in CB-MMC are more likely to regularly converge into large bubbles. Therefore, CB-MMC in Fig. 9(d) has a more uniform pressure fluctuation. In addition, it can be deduced that when the heat flux increases, a strip-shaped drying area starts to appear on the left side wall of the outlet manifold. This is because the fluid forms vortex at the outlet of the manifold, where bubbles gather and grow, while the low flow

velocity near the outlet cannot discharge the bubbles in time, which results in the generation of large strip-shaped bubbles [32].

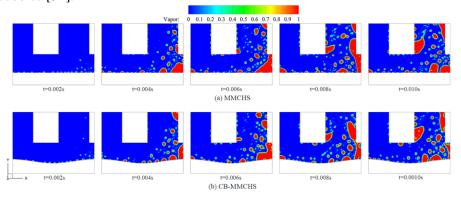


Fig. 10. The changes in vapor phase nephograms of (a) MMC and (b) CB-MMC at 400 W/cm²

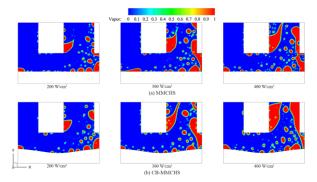


Fig. 11. Comparison between the vapor phase nephograms of (a) MMC and (b) CB-MMC at 0.025s

Conclusions

The flow and heat transfer characteristics of HFE-7100 in MMC and CB-MMC are numerically studied under the single-phase and two-phase flow conditions. The main conclusions are as follows:

- (1) In the single-phase flow of HFE-7100, in MMC and CB-MMC, the Tave values decrease with the flow rate increasing, while the ΔP values show an increasing trend. When the channel mass flow rate of 1500 kg/(m²·s), the CB-MMC reduces the pressure drop by 4.1% at a temperature reduction of less than 1% compared to the MMC. Under the same conditions, CB-MMC shows lower $T_{\rm ave}$ and ΔP . At both $T_{\rm ave}$ of 48 °C, the pressure drop of the MMC increases by 9.3% compared to the CB-MMC. Thus the CB-MMC has higher flow and convective heat transfer performance than MMC in single-phase flow.
- (2) A multi-factor analysis of the structural parameters in CB-MMC is performed using the orthogonal test method to obtain optimal solutions 17, 18 and 19 respectively. By comparison, it is found that solution 19, obtained using the weight matrix analysis method, obtained a more suitable comprehensive heat transfer performance with a $T_{\rm ave}$ 6.89 °C lower than that in solution 18 and a ΔP 10.27 kPa lower than that in solution 17. It is also found that the PEC of all three optimal solutions is higher than that of the MMC. The optimal solution 19 has the highest PEC with an increase of 67.5% compared to the MMC.
- (3) When subcooled boiling of HFE-7100 occurs in the microchannel, the speeds of bubbles growing, gathering, and discharging increase, and thus the magnitude of the fluctuations of ΔP increases with the increase of the heat flux, which also leads to small fluctuations of T_{ave} after boiling

equilibrium is reached. It can be deduced that the vapor bubbles in CB-MMC are more easily generated and discharged by comparing the boiling bubble characteristics in MMC and CB-MMC, and the fluctuations of ΔP and $T_{\rm ave}$ in the microchannel are more uniform. Therefore, CB-MMC also shows higher two-phase flow and boiling heat transfer performance.

Acknowledgment

This work was financially supported by the Major Science and Technology Project of Anhui Province (No. 202003a05020014), the Hefei Natural Science Foundation (No.2021045), and the Fundamental Research Funds for the Central Universities (No.JZ2021HGTA0150, No.JZ2021HGQA0239).

Nomenclature

$C_{ m p}$	specific heat capacity, [J·kg ⁻¹ ·K ⁻¹]	Greek let	ters
$G_{ m in}$	mass flow rate at inlet, [kg·m ⁻² ·s ⁻¹]	α	volume fraction, [-]
$H_{\rm a}$	amplitude of single-cell model, [µm]	$\dot{m{lpha}}_{ m lv}$	liquid fraction generation rate, [s ⁻¹]
$i_{ m lv}$	specific enthalpy, [kJ·kg ⁻¹]	λ	wavelength of single-cell model, [μm]
k	thermal conductivity, $[W \cdot m^{-1} \cdot K^{-1}]$	μ	dynamic viscosity, [kg·m ⁻¹ ·s ⁻¹]
L_{in}	length of inlet, [μm]	$\boldsymbol{\rho}$	density, [kg·m ⁻³]
$L_{ m out}$	length of outlet, [µm]	σ	surface tension coefficient, [N·m ⁻¹]
$L_{ m div}$	length of divider, [μm]	ω	the weight matrix, [-]
$L_{ m sub}$	height of substrate, [µm]	Subscript	ts
$L_{\rm c}$	height of channel, [µm]	1	liquid phase, [-]
	-, -		1
ΔP	pressure drop, [kPa]	lv	liquid-vapor phase change, [-]
ΔP $T_{ m ave}$	pressure drop, [kPa] average temperature, [K]	lv s	
			liquid-vapor phase change, [-]
T_{ave}	average temperature, [K]	s	liquid-vapor phase change, [-] solid, [-]
$T_{ m ave} \ \dot{m v}_{ m lv}$	average temperature, [K] volumetric evaporation rate, [s ⁻¹]	s sat	liquid-vapor phase change, [-] solid, [-] saturation state, [-]

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Submitted: 7.7.2023. Revised: 6.9.2023. Accepted: 11.9.2023.