

ANALYTIC ALGORITHM FOR LOCAL FRACTIONAL CAUDREY-DODD-GIBBON-KAEADA EQUATION BASED ON THE NEW ITERATIVE METHOD

by

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In this paper, the initial value problem is discussed for the local fractional Caudrey-Dodd-Gibbon-Kaeada equation. The fractional complex transform and the new iterative method are used to solve the problem, and the approximate analytical solutions are obtained.

Key words: Caudrey-Dodd-Gibbon-Kaeada equation, local fractional derivative, complex transform method, new iterative method

Introduction

The differential equations with local fractional derivatives have proved to be a suitable tool for modeling many non-differentiable phenomena. Many physical problems in fractal media lead to non-linear models involving local fractional derivatives [1, 2]. Recently, several authors investigated the non-linear local fractional heat equation for the anomalous diffusion on a fractal medium [3, 4], the fractal population dynamics [5], and the fractional generalized KdV system [6].

Our interest here is to solve the following non-linear local fractional partial differential equation:

$$\frac{\partial^\alpha u(x,t)}{\partial t^\alpha} = \frac{\partial^\beta}{\partial x^\beta} \left(au^3 + bu \frac{\partial^{2\beta} u}{\partial x^{2\beta}} + c \frac{\partial^{4\beta} u}{\partial x^{4\beta}} \right) \quad (1)$$

with the condition:

$$u(x,0) = \phi \left[\frac{x^\beta}{\Gamma(1+\beta)} \right] \quad (2)$$

where $\partial^\alpha u / \partial t^\alpha$ and $\partial^\beta u / \partial x^\beta$ are the local fractional derivatives ($0 < \alpha \leq 1, 0 < \beta \leq 1$), a, b , and c are constants, and $\phi(x)$ is given function.

When $\alpha = \beta = 1$, eq. (1) is called the Caudrey-Dodd-Gibbon-Kaeada equation [7], it has wide applications in the description of many non-linear phenomena in physics and chem-

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istry [8-10]. This problem is often challenging to be solved analytically. In recent years, some analytical methods for solving local fractional differential equations have been proposed, for examples, the variational iteration method [11-13], the homotopy perturbation method [14, 15], the local fractional series expansion [16], and the local fractional function decomposition [17]. The new iterative method (NIM) was proposed first by Daftardar-Gejji and Jafari in [18], it is an effective procedure to find approximate analytical solutions of a wide class of non-linear local fractional differential equations [19, 20].

The main purpose of this paper is to solve the problem of eqs. (1) and (2) by using the fractional complex transform [21-24] and NIM [18-20].

Preliminaries

Local fractional derivative

In this section, we give some basic definitions and properties of the local fractional calculus theory [1-3].

Definition 1. For arbitrary $\varepsilon > 0$, assume that the relation below exists:

$$|f(x) - f(x_0)| < \varepsilon^\alpha \quad (3)$$

with $|x - x_0| < \delta$. Then $f(x)$ is called local fractional continuous at x_0 which is denoted by $\lim_{x \rightarrow x_0} f(x) = f(x_0)$. If $f(x)$ is local fractional continuous on the interval (a, b) , it is denoted by $f(x) \in C_\alpha(a, b)$.

$$f(x) \in C_\alpha(a, b)$$

Definition 2. Let $f(x) \in C_\alpha(a, b)$. In fractal space, the local fractional derivative of $f(x)$ of order α at the point $x = x_0$ is given by:

$$D_x^\alpha f(x_0) = \left. \frac{d^\alpha}{dx^\alpha} f(x) \right|_{x=x_0} = f^{(\alpha)}(x_0) = \lim_{x \rightarrow x_0} \frac{\Delta^\alpha [f(x) - f(x_0)]}{(x - x_0)^\alpha} \quad (4)$$

where $\Delta[f(x) - f(x_0)] \cong \Gamma(\alpha + 1)[f(x) - f(x_0)]$.

Definition 3. The local fractional partial derivative of a high order is defined in the form:

$$\frac{\partial^{k\alpha} u(x, t)}{\partial x^{k\alpha}} = \overbrace{\frac{\partial^\alpha}{\partial x^\alpha} \frac{\partial^\alpha}{\partial x^\alpha} \cdots \frac{\partial^\alpha}{\partial x^\alpha}}^{k \text{ times}} u(x, t) \quad (5)$$

The following property holds true:

$$\frac{\partial^\alpha u[g(x, t)]}{\partial x^\alpha} = u'[g(x, t)] \frac{\partial^\alpha g(x, t)}{\partial x^\alpha} \quad (6)$$

where there exist $u'[g(x, t)]$ and $[\partial^\alpha g(x, t)]/\partial x^\alpha$.

The new iterative method

Below we illustrate the main points of NIM [18-20], by considering the following general function equation:

$$u = L(u) + \Phi(u) + \Pi \quad (7)$$

where L is a linear operator, Φ – a non-linear operator, and Π – a known function.

Assume that the solution of eq. (7) is of the form:

$$u = \sum_{i=0}^{\infty} u_i \quad (8)$$

The non-linear operator Φ can be decomposed as:

$$\Phi\left(\sum_{i=0}^{\infty} u_i\right) = \Phi(u_0) + \sum_{i=1}^{\infty} \left[\Phi\left(\sum_{j=0}^i u_j\right) - \Phi\left(\sum_{j=0}^{i-1} u_j\right) \right] \quad (9)$$

From eqs. (8) and (9), eq. (7) is equivalent to:

$$\sum_{i=0}^{\infty} u_i = \Pi + \sum_{i=0}^{\infty} T(u_i) + \Phi(u_0) + \sum_{i=1}^{\infty} \left[\Phi\left(\sum_{j=0}^i u_j\right) - \Phi\left(\sum_{j=0}^{i-1} u_j\right) \right] \quad (10)$$

Define the recurrence relation:

$$\begin{aligned} u_0 &= \Pi(x) \\ u_1 &= L(u_0) + M_0 \\ u_m &= L(u_m) + M_m, \quad m = 1, 2, \dots \end{aligned} \quad (11)$$

where

$$M_0 = \Phi(u_0) \quad (12)$$

$$M_m = \Phi\left(\sum_{i=0}^m u_i\right) - \Phi\left(\sum_{i=0}^{m-1} u_i\right), \quad m = 1, 2, \dots \quad (13)$$

Then the p -term approximate solution of (7) is given by:

$$u = u_0 + u_1 + \dots + u_{p-1} \quad (14)$$

Solution of the problem (1)-(2)

In this section, we consider the following initial value problem of non-linear local fractional Caudrey-Dodd-Gibbon-Kaeada equation:

$$\begin{aligned} \frac{\partial^\alpha u}{\partial t^\alpha} &= \frac{\partial^\beta}{\partial x^\beta} \left(au^3 + bu \frac{\partial^{2\beta} u}{\partial x^{2\beta}} + c \frac{\partial^{4\beta} u}{\partial x^{4\beta}} \right) \\ u(x, 0) &= \varphi \left[\frac{x^\beta}{\Gamma(1 + \beta)} \right] \end{aligned} \quad (15)$$

By using the fractional complex transform [21-24]:

$$X = \frac{x^\beta}{\Gamma(1+\beta)}, \quad T = \frac{t^\alpha}{\Gamma(1+\alpha)} \quad (16)$$

the problem (15) becomes:

$$\frac{\partial U}{\partial T} = \frac{\partial}{\partial X} \left(aU^3 + bU \frac{\partial^2 U}{\partial X^2} + \frac{\partial^4 U}{\partial X^4} \right) \quad (17)$$

$$U(X, 0) = \Phi(X)$$

The understanding of eq. (16) was discussed in detail by the two-scale fractal theory for various applications [25-34]. We rewrite the problem (17):

$$U(X, T) = U(X, 0) + \int_0^T \left(\sum_{k=1}^4 N_k + c \frac{\partial^4 U}{\partial X^4} \right) dT \quad (18)$$

where

$$N_1 = aU^3 \frac{\partial U}{\partial X}, \quad N_2 = bU^2 \frac{\partial^3 U}{\partial X^3}, \quad N_3 = bU \frac{\partial U}{\partial X} \frac{\partial^2 U}{\partial X^2}, \quad N_4 = b \frac{\partial U}{\partial X} \frac{\partial^2 U}{\partial X^2} \frac{\partial^3 U}{\partial X^3}$$

Suppose that the solution of eq. (18) takes the form:

$$U(X, T) = \sum_{k=0}^{\infty} U_k(X, T) \quad (19)$$

and the non-linear term in eq. (18) is decomposed:

$$N_1 = aU_0^2 \frac{\partial U_0}{\partial X} + a \sum_{m=1}^{\infty} N_{1m} \quad (20)$$

$$N_2 = bU_0 \frac{\partial^3 U_0}{\partial X^3} + b \sum_{m=1}^{\infty} N_{2m} \quad (21)$$

$$N_3 = b \frac{\partial U_0}{\partial X} \frac{\partial^2 U_0}{\partial X^2} + b \sum_{m=1}^{\infty} N_{3m} \quad (22)$$

where

$$N_{1m} = \left(\sum_{k=0}^m U_k \right)^2 \frac{\partial}{\partial X} \left(\sum_{k=0}^m U_k \right) - \left(\sum_{k=0}^{m-1} U_k \right)^2 \frac{\partial}{\partial X} \left(\sum_{k=0}^{m-1} U_k \right)$$

$$N_{2m} = \left(\sum_{k=0}^m U_k \right) \frac{\partial^3}{\partial X^3} \left(\sum_{k=0}^m U_k \right) - \left(\sum_{k=0}^{m-1} U_k \right) \frac{\partial^3}{\partial X^3} \left(\sum_{k=0}^{m-1} U_k \right)$$

$$N_{3m} = \frac{\partial}{\partial X} \left(\sum_{k=0}^m U_k \right) \frac{\partial^2}{\partial X^2} \left(\sum_{k=0}^m U_k \right) - \frac{\partial}{\partial X} \left(\sum_{k=0}^{m-1} U_k \right) \frac{\partial^2}{\partial X^2} \left(\sum_{k=0}^{m-1} U_k \right)$$

Now, we let:

$$N_0 = aU_0^2 \frac{\partial U_0}{\partial X} + bU_0 \frac{\partial^3 U_0}{\partial X^3} + b \frac{\partial U_0}{\partial X} \frac{\partial^2 U_0}{\partial X^2}$$

then according to NIM, we obtain:

$$U_0(X, T) = \Phi(X)$$

$$U_1(X, T) = \int_0^T \left(N_0 + c \frac{\partial^5 U_0}{\partial X^5} \right) dT$$

$$U_2(X, T) = \int_0^T \left(N_{11} + N_{21} + N_{31} + c \frac{\partial^5 U_1}{\partial X^5} \right) dT$$

⋮

$$U_m(X, T) = \int_0^T \left(N_{1,m-1} + N_{2,m-1} + N_{3,m-1} + c \frac{\partial^5 U_{m-1}}{\partial X^5} \right) dT \quad (23)$$

Thus the p -term approximate solution of eq. (18) is given by:

$$U(X, T) = U_0(X, T) + U_1(X, T) + U_2(X, T) + \cdots + U_p(X, T)$$

From eq. (14), we can get the solution of the problem (15):

$$u(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + \cdots u_n(x, t) + \cdots$$

Application

Consider the problem (1)-(2) in the form:

$$\begin{aligned} \frac{\partial^\alpha u}{\partial t^\alpha} &= \frac{\partial^\beta}{\partial x^\beta} \left(-60u^3 - 30bu \frac{\partial^{2\beta} u}{\partial x^{2\beta}} - \frac{\partial^{4\beta} u}{\partial x^{4\beta}} \right) \\ u(x, 0) &= \frac{1}{4} \sec h^2 \left[\frac{x^\beta}{2\Gamma(1+\beta)} \right] \end{aligned} \quad (24)$$

By the relations (16), we obtain:

$$\begin{aligned} \frac{\partial U}{\partial T} &= \frac{\partial}{\partial X} \left(-60U^3 - 30U \frac{\partial^{2\beta} U}{\partial X^{2\beta}} - \frac{\partial^{4\beta} U}{\partial X^{4\beta}} \right) \\ U(X, 0) &= \frac{1}{4} \sec h^2 \left(\frac{X}{2} \right) \end{aligned} \quad (25)$$

$$\text{Let } \exp(-2X) = E, \quad \exp\left[\frac{-2x^\beta}{\Gamma(1+\beta)}\right] = F.$$

Then by eq. (23), we have:

$$\begin{aligned} U_0(X, T) &= \frac{4E}{(1+E)^2} \\ U_1(X, T) &= \frac{128(E-E^2)}{(1+E)^2} T \\ U_2(X, T) &= \frac{2048(E^3-4E^2+E)}{(1+E)^4} T^2 \\ U_3(X, T) &= \frac{65536(11E^3-11E^2-E^4+E)}{3(1+E)^5} T^2 \\ &\vdots \end{aligned}$$

and so on.

Hence, from eq. (17), we obtain:

$$\begin{aligned} u_0(x, t) &= \frac{4F}{(1+F)^2} \\ u_1(x, t) &= \frac{128(F-F^2)}{(1+F)^3 \Gamma(1+\alpha)} t^\alpha \\ u_2(x, t) &= \frac{2048(F^3-4F^2+F)}{(1+F)^4 \Gamma^2(1+\alpha)} t^{2\alpha} \\ u_3(x, t) &= \frac{65536(11F^3-11F^2-F^4+F)}{3(1+F)^5 \Gamma^3(1+\alpha)} t^{3\alpha} \\ &\vdots \end{aligned}$$

Finally, the solution of eq. (23) is given by:

$$u(x, t) = u_0(x, t) + u_1(x, t) + u_2(x, t) + u_3(x, t) + u_4(x, t) + \dots$$

When $\alpha = \beta = 1$, if let $\exp(-2x) = G$, then we get:

$$u(x, t) = \frac{4G}{(1+G)^2} + \frac{128(G-G^2)}{(1+G)^3} t + \frac{2048(G^3-4G^2+G)}{(1+G)^4} t^2 + \dots$$

which is close to the exact solution [10]:

$$u(x, t) = \text{sech}^2(x - 16t)$$

Conclusion

We have presented an analytic algorithm for non-linear local fractional Caudrey-Dodd-Gibbon-Kawada equation. The fractional complex transform method and NIM have been successfully applied to find the approximate analytical solutions of the equation. The results show that NIM is a powerful and efficient technique in finding approximate analytical solutions for non-linear local fractional differential equations.

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