

## FRACTIONAL RESIDUAL METHOD COUPLED WITH ADOMIAN DECOMPOSITION METHOD FOR SOLVING LOCAL FRACTIONAL DIFFERENTIAL EQUATIONS

by

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*In order to solve the local fractional differential equations, we couple the fractional residual method with the Adomian decomposition method via the local fractional calculus operator. Several examples are given to illustrate the solution process and the reliability of the method.*

Key words: *local fractional derivative, fractional residual method, Adomian decomposition method*

### Introduction

It is an obvious fact that fractional differential equations can describe non-differentiable problems much better than classic differential equations [1-8]. Now the fractional calculus has been developing rapidly, and it has been widely applied to physics, mathematics, engineering, and many other fields [1-5]. However, there are still many problems which are needed to be solved urgently. The main barrier is the definition of the fractional derivative. There are too many definitions, and new definitions always appear in literature, however, most fractional derivatives are almost incompatible, and the physical understanding of the fractional models is unclear. The aforementioned issues are worthy of further study in the future.

The local fractional derivative [9, 10] is defined in Cantor fractal space and can model many practical problems. Fractal theory is also a powerful tool to the analysis of biologic and material phenomena [11, 12]. How to solve such problems has become a hot topic in mathematics, and some effective methods have appeared in literature, *e.g.*, the variational iteration method [13-15], the fractional residual method [16], the homotopy perturbation method [17-20], the local fractional Fourier series method [21-23], The Yang Laplace transform-DJ iteration method [24], He-Laplace method [25, 26], the fractional complex transform (two-scale transform) [27, 28], the coupled method of the variational iteration and reduced differential transform method [29], the differential transform approach [30], the asymptotic perturbation method [31], the coupled method of the Sumudu transform and the variational iteration method [32], the direct algebraic method [33, 34], the exp-function method [35], the variational approach [36-39], the Fourier spectral method [40] and the reproducing kernel method [41].

In this paper, we attempt to solve the local differential equation by coupling the fractional residual method [16] with the Adomian decomposition method [42] *via* the local differential equations.

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### Basic definitions of local fractional calculus

In this section, we introduce some mathematical preliminaries of the local fractional calculus theory in fractal space for our subsequent development [6].

*Definition 1.* In fractal space, let  $u(t) \in C_\alpha(a, b)$ , the local fractional derivative of  $u(t)$  of order  $\alpha$  at  $t = t_0$  is given by [6, 9, 10]:

$$D_t^{(\alpha)} u(t_0) = u^{(\alpha)}(t_0) = \frac{d^\alpha u(t)}{dt^\alpha} \Big|_{t=t_0} = \lim_{t \rightarrow t_0} \frac{\Delta^\alpha [u(t) - u(t_0)]}{(t - t_0)^\alpha} \quad (1)$$

where

$$\Delta^\alpha [u(t) - u(t_0)] \cong \Gamma(1 + \alpha) \Delta [u(t) - u(t_0)]$$

*Definition 2.* [6] let  $u(t) \in C_\alpha(a, b)$ , the local fractional integral of  $(t)$  of order  $\alpha$  in the interval  $[a, b]$  is defined by [6, 9, 10]:

$${}_a I_b^{(\alpha)} u(t) = \frac{1}{\Gamma(1 + \alpha)} \int_a^b u(t) (dt)^\alpha = \frac{1}{\Gamma(1 + \alpha)} \lim_{\Delta t \rightarrow 0} \sum_{j=0}^{j=N-1} u(t_j) (\Delta t_j)^\alpha \quad (2)$$

where  $\Delta t_j = t_{j+1} - t_j$ ,  $\Delta t = \max\{\Delta t_1, \Delta t_2, \Delta t_j, \dots\}$  and  $(t_j, t_{j+1})$ ,  $j = 0, \dots, N-1$ ,  $t_0 = a$ ,  $t_N = b$ , is a partition of the interval  $[a, b]$ .

### The fractional residual method coupled with the Adomian decomposition method

In this section, we shall present the iterative process of the fractional residual method [16] coupled with the Adomian decomposition method [42] to search for the exact solution of some local fractional differential equations.

Firstly, we discuss the following local fractional differential equation on the fractal set, and its form is given:

$$Lu(x, t) - Ru(x, t) - Pu(x, t) = g(x, t) \quad (3)$$

where  $L, P$  are all linear operators,  $R$  is a non-linear operator and  $g(x, t)$  is an inhomogeneous term.

In order to solve eq. (3), we apply the local fractional reverse operator  $L^{-1}(\bullet)$  on both sides of eq. (3), then we obtain:

$$u(x, t) = \tilde{u}(x, t) + L^{-1}[Ru(x, t) + Pu(x, t) + g(x, t)] \quad (4)$$

where  $\tilde{u}(x, t)$  is derived from the initial condition.

Supposing  $u_0(x, t)$  is a function to be determined, we can rewrite eq. (4):

$$u(x, t) = \tilde{u}(x, t) + L^{-1}u_0 + L^{-1}[-u_0 + Ru(x, t) + Pu(x, t) + g(x, t)] \quad (5)$$

If we let:

$$L^{-1}[-u_0 + Ru(x, t) + Pu(x, t) + g(x, t)] = 0 \quad (6)$$

Solving eq. (6), we can determine  $u_0(x, t)$ , and then substituting  $u_0(x, t)$  into eq. (5), we can get the exact solution of eq. (4):

$$u(x, t) = \tilde{u}(x, t) + L^{-1}u_0 \quad (7)$$

According to eq. (6), we can construct the following equation:

$$-u_0 + R[\tilde{u}(x, t) + L^{-1}u_0] + P[\tilde{u}(x, t) + L^{-1}u_0] + g(x, t) = 0 \quad (8)$$

By analyzing the process of this method, we can deduce that it is critical to get the exact solution of eq. (3) to choose  $u_0(x, t) = 0$ .

Let the solution  $u_0(x, t)$  of eq. (8) has the following series form:

$$u_0(x, t) = \sum_{n=0}^{\infty} v_n(x, t), \quad n = 1, 2, \dots \quad (9)$$

Substituting eq. (9) into eq. (8), we can get:

$$\sum_{n=0}^{\infty} v_n(x, t) = R \left\{ \tilde{u}(x, t) + L^{-1} \left[ \sum_{n=0}^{\infty} v_n(x, t) \right] \right\} + P \left\{ \tilde{u}(x, t) + L^{-1} \left[ \sum_{n=0}^{\infty} v_n(x, t) \right] \right\} + g(x, t) \quad (10)$$

We let  $A_n(u)$  be the Adomian polynomials, which are:

$$A_n(u) = \frac{1}{\Gamma(1+n\alpha)} \frac{d^{n\alpha}}{d\lambda^{n\alpha}} \left( R \left\{ \tilde{u}(x, t) + L^{-1} \left[ \sum_{i=0}^{\infty} \lambda^{i\alpha} v_i(x, t) \right] \right\} \right) \Big|_{\lambda=0} \quad (11)$$

And we suppose  $g(x, t) + P[\tilde{u}(x, t)]$  can be decomposed as the following series:

$$g(x, t) + P[\tilde{u}(x, t)] = \sum_{n=0}^{\infty} g_n \quad (12)$$

According to eqs. (11) and (12) and the Adomian decomposition method, we can construct the following new recursion scheme:

$$\begin{aligned} v_0 &= g_0 \\ v_1 &= g_1 + P(v_0) + A_0 \\ v_{m+1} &= g_m + P(v_m) + A_m \end{aligned} \quad (13)$$

Obviously, similarly to the classic Adomian decomposition method, we can easily verify that:

$$u_0 = \sum_{n=0}^{\infty} v_n \quad (14)$$

converges.

Then, by virtue of eqs. (7) and (14), the exact solution of eq. (4) can be easily derived.

### Illustrative examples

To demonstrate the effectiveness of the method, several local fractional PDE are presented.

*Example 1.* Consider the following local fractional Schrodinger conduction equation:

$$i^\alpha \frac{\partial^\alpha u}{\partial x^\alpha}(x, t) + \frac{\partial^{2\alpha} u}{\partial t^{2\alpha}}(x, t) + |u(x, t)|^2 [u(x, t)] + E_\alpha [2i(t-x)^\alpha] = 0 \quad (15)$$

subject to the initial condition:

$$u(0, t) = E_\alpha (2i^\alpha t^\alpha) \quad (16)$$

According to eq. (5), eq. (15) can be transformed into the following equation:

$$u_0(x, t) = i^\alpha E_\alpha [2i(t-x)^\alpha] + i^\alpha \left( \frac{\partial^{2\alpha} [\tilde{u}(x, t) + L^{-1}u_0]}{\partial t^{2\alpha}}(x, t) + \left| L^{-1}[\tilde{u}(x, t) + L^{-1}u_0](x, t) \right|^2 \{L^{-1}[\tilde{u}(x, t) + L^{-1}u_0](x, t)\} \right) \quad (17)$$

Letting  $u_0(x, t) = \sum_{n=0}^{\infty} v_n(x, t)$  and substituting it into eq. (17), we can derive:

$$\begin{aligned} \sum_{n=0}^{\infty} v_n(x, t) &= \frac{-3x^\alpha i^\alpha}{\Gamma(1+\alpha)} E_\alpha (2i^\alpha t^\alpha) + i^\alpha E_\alpha [2i(t-x)^\alpha] + i^\alpha \left[ \sum_{n=0}^{\infty} \frac{\partial^{2\alpha} L^{-1}v_n}{\partial t^{2\alpha}}(x, t) + \sum_{n=0}^{\infty} L^{-1}A_n(u) \right] = \\ &= \frac{-3x^\alpha i^\alpha}{\Gamma(1+\alpha)} E_\alpha (2i^\alpha t^\alpha) + i^\alpha E_\alpha [2i(t-x)^\alpha] + i^\alpha \left[ \sum_{n=0}^{\infty} \frac{\partial^{2\alpha} L^{-1}v_n}{\partial t^{2\alpha}}(x, t) + \sum_{n=0}^{\infty} L^{-1}A_n(u) \right] \quad (18) \end{aligned}$$

where  $A_n(u)$  is the Adomian polynomials for the non-linear term  $|u(t, x)|^2 u(t, x)$  and  $A_n(u)$  is given:

$$A_n(u) = \frac{1}{\Gamma(1+n\alpha)} \frac{d^{n\alpha}}{d\lambda^{n\alpha}} \left\{ \left[ \sum_{i=0}^{\infty} \lambda^{i\alpha} L^{-1}v_i(x, t) \right]^2 \left[ \sum_{i=0}^{\infty} \lambda^{i\alpha} L^{-1}v_i(x, t) \right] \right\}_{\lambda=0} \quad (19)$$

Obviously, according to eq. (13), the first few Adomian polynomials are respectively given by:

$$\begin{aligned}
 A_0(x, t) &= v_0 \bar{v}_0 v_0 = E_\alpha(2it)^\alpha \\
 A_1(x, t) &= v_0 \bar{v}_0 v_1 + v_0 v_0 \bar{v}_1 + v_0 \bar{v}_0 v_1 = \frac{(-2ix)^\alpha}{\Gamma(1+\alpha)} E_\alpha(2it)^\alpha \\
 A_2(x, t) &= v_0 v_0 \bar{v}_2 + v_0 \bar{v}_0 v_2 + v_0 \bar{v}_1 v_1 + v_0 v_0 \bar{v}_2 + v_1 v_1 \bar{v}_0 + v_0 \bar{v}_1 v_1 = \frac{(-2ix)^{2\alpha}}{\Gamma(1+2\alpha)} E_\alpha(2it)^\alpha \quad (20) \\
 &\dots\dots \\
 A_n(x, t) &= \dots = \frac{(-2ix)^{n\alpha}}{\Gamma(1+n\alpha)} E_\alpha(2it)^\alpha
 \end{aligned}$$

In the light of eq. (12), we let:

$$G = \sum_{n=0}^{\infty} G_n = \frac{-3x^\alpha t^\alpha}{\Gamma(1+\alpha)} E_\alpha(2i^\alpha t^\alpha) + i^\alpha E_\alpha[2i(t-x)^\alpha] \quad (21)$$

where

$$\begin{aligned}
 G_0 &= i^\alpha E_\alpha(2it)^\alpha, \quad G_1 = \frac{-5(ix)^\alpha}{\Gamma(1+\alpha)} E_\alpha(2it)^\alpha \\
 G_2 &= 2^2 \frac{(ix)^{2\alpha}}{\Gamma(1+\alpha)} E_\alpha(2it)^\alpha, \dots, \quad G_n = (-2)^m \frac{(ix)^{m\alpha}}{\Gamma(1+m\alpha)} E_\alpha(2it)^\alpha \quad (22)
 \end{aligned}$$

Using eqs. (12), (13), (20), and (22), gives:

$$\begin{aligned}
 v_0 &= G_0 = i^\alpha E_\alpha(2it)^\alpha \\
 v_1 &= \frac{-5(ix)^\alpha}{\Gamma(1+\alpha)} E_\alpha(2it)^\alpha - i^\alpha \left[ \frac{\partial^{2\alpha} L^{-1} v_0}{\partial t^{2\alpha}}(x, t) + L^{-1} A_0 \right] d\tau = -(2i)^\alpha \frac{x^\alpha}{\Gamma(1+\alpha)} E_\alpha(2it)^\alpha \\
 v_2 &= -2^2 \frac{(ix)^{2\alpha}}{\Gamma(1+2\alpha)} E_\alpha(2it)^\alpha - i^\alpha \left[ \frac{\partial^{2\alpha} L^{-1} v_1}{\partial t^{2\alpha}}(x, t) + L^{-1} A_1 \right] d\tau = (2i)^{2\alpha} \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} E_\alpha(2it)^\alpha \\
 v_3 &= 8 \frac{(ix)^{3\alpha}}{\Gamma(1+3\alpha)} E_\alpha(2it)^\alpha - i^\alpha \left[ \frac{\partial^{2\alpha} L^{-1} v_2}{\partial t^{2\alpha}}(x, t) + L^{-1} A_2 \right] d\tau = -(2i)^{3\alpha} \frac{x^{3\alpha}}{\Gamma(1+3\alpha)} E_\alpha(2it)^\alpha \\
 v_m &= (-2)^m \frac{(ix)^{m\alpha}}{\Gamma(1+m\alpha)} E_\alpha(2it)^\alpha + i^\alpha \left[ \frac{\partial^{2\alpha} v_{m-1}}{\partial t^{2\alpha}}(x, t) + A_m \right] d\tau = (-2i)^{m\alpha} \frac{x^{m\alpha}}{\Gamma(1+m\alpha)} E_\alpha(2it)^\alpha \quad (23)
 \end{aligned}$$

Proceeding in this manner, the rest of the components  $u_n(x, t)$  can be completely determined and then, making using of eq. (7), the exact solution of eq. (15) is:

$$u(x, t) = \sum_{m=0}^{\infty} (-2i)^{m\alpha} \frac{x^{m\alpha}}{\Gamma(1+m\alpha)} E_\alpha(2it)^\alpha = E_\alpha[2i(t-x)^\alpha] \quad (24)$$

*Example 2.* Consider the following non-linear gas dynamic equation:

$$\frac{\partial^{2\alpha} u(x,t)}{\partial x^{2\alpha}} + u(x,t) \frac{\partial^{2\alpha} u(x,t)}{\partial t^{2\alpha}} - u(x,t)[1 - u(x,t)] = 0 \quad (25)$$

subject to the boundary and initial conditions:

$$\begin{aligned} u(0,t) &= \sin_{\alpha}(t^{\alpha}) \\ \frac{\partial^{\alpha} u(0,t)}{\partial x^{\alpha}} &= \sin_{\alpha}(t^{\alpha}) \end{aligned} \quad (26)$$

Applying the inverse operator  $L^{-1}(\bullet) = {}_0I_x^{(2\alpha)}(\bullet)$  on both sides of eq. (25), and making use of eq. (5), we obtain:

$$\begin{aligned} u(x,t) &= \sin_{\alpha}(t^{\alpha}) \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] + {}_0I_x^{(2\alpha)} u_0 - \\ &- {}_0I_x^{(2\alpha)} \left\{ u_0 + u(x,t) \frac{\partial^{2\alpha} u(x,t)}{\partial t^{2\alpha}} - u(x,t)[1 - u(x,t)] \right\} \end{aligned} \quad (27)$$

According to eqs. (6) and (7), we let:

$$u(x,t) = \sin_{\alpha}(t^{\alpha}) \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] + {}_0I_x^{(2\alpha)} u_0 \quad (28)$$

and

$${}_0I_t^{(2\alpha)} \left\{ u_0 + u(x,t) \frac{\partial^{2\alpha} u(x,t)}{\partial t^{2\alpha}} - u(x,t)[1 - u(x,t)] \right\} = 0 \quad (29)$$

Then, we suppose:

$$u_0 + u(x,t) \frac{\partial^{2\alpha} u(x,t)}{\partial t^{2\alpha}} - u(x,t)[1 - u(x,t)] = 0 \quad (30)$$

and

$$u_0(x,t) = \sum_{n=0}^{\infty} v_n(x,t) \quad , n=1,2,\dots \quad (31)$$

where  $u_n(t)$  are all constants to be determined.

Substituting the eqs. (28) and (31) into the eq. (30), we can obtain:

$$\begin{aligned} &\sum_{n=0}^{\infty} v_n(x,t) - \sin_{\alpha}(t^{\alpha}) \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] - \sum_{n=1}^{\infty} {}_0I_t^{(2\alpha)} v_n(x,t) + \\ &+ \left\{ \sin_{\alpha}(t^{\alpha}) \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] + \sum_{n=1}^{\infty} {}_0I_x^{(2\alpha)} v_n(x,t) \right\} \sin_{\alpha} t^{\alpha} \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] = 0 \end{aligned} \quad (32)$$

In the light of eq. (12), we let:

$$G = \sum_{n=0}^1 G_n = \sin_{\alpha}(t^{\alpha}) \left[ 1 + \frac{x^{\alpha}}{\Gamma(1+\alpha)} \right] \quad (33)$$

where

$$G_0 = \sin_{\alpha} t^{\alpha}, \quad G_1 = \sin_{\alpha}(t^{\alpha}) \frac{x^{\alpha}}{\Gamma(1+\alpha)} \quad (34)$$

Using eqs. (11), (13), (32), and (34), gives:

$$\begin{aligned} v_0 &= G_0 = \sin_{\alpha}(t^{\alpha}) \\ v_1 &= \sin_{\alpha}(t^{\alpha}) \frac{x^{\alpha}}{\Gamma(1+\alpha)} + {}_0I_x^{(2\alpha)} v_0(x, t) = \sin_{\alpha}(t^{\alpha}) \left[ \frac{x^{\alpha}}{\Gamma(1+\alpha)} + \frac{x^{2\alpha}}{\Gamma(1+2\alpha)} \right] \\ v_2 &= {}_0I_x^{(2\alpha)} v_1(x, t) = \sin_{\alpha}(t^{\alpha}) \left[ \frac{x^{3\alpha}}{\Gamma(1+\alpha)} + \frac{x^{4\alpha}}{\Gamma(1+2\alpha)} \right] \\ v_m &= {}_0I_x^{(2\alpha)} v_{m-1}(x, t) = \sin_{\alpha}(t^{\alpha}) \left\{ \frac{x^{(m+1)\alpha}}{\Gamma[1+(m+1)\alpha]} + \frac{x^{(m+2)\alpha}}{\Gamma[1+(m+2)\alpha]} \right\} \\ &\dots \end{aligned} \quad (35)$$

Substituting the eq. (35) into eq. (31), and then substituting the result into eq. (30), we can obtain the following exact solution of the eq. (25):

$$u(x, t) = E_{\alpha}(x^{\alpha}) \sin_{\alpha}(t^{\alpha}) \quad (36)$$

## Conclusions

In this article, we have suggested the fractional residual method coupled with the Adomian decomposition method for solving the local fractional equation. The test examples show that the method is simple and feasible, and it is effective for solving some local differential equation problems.

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