

PERIODIC SOLUTION OF FRACTAL PHI-4 EQUATION

by

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This paper focuses on a fractal Phi-4 equation with time-space fractal derivatives, though its solitary solutions have been deeply studied, its periodic solution was rarely revealed due to its strong non-linearity. Now the condition is completely changed, He's frequency formulation provides with a universal tool to having a deep insight into the periodic property of the fractal Phi-4 equation. The two-scale transform is used to convert approximately the fractal Phi-4 equation a differential model, and a criterion is suggested for the existence of a periodic solution of the equation, the effect of fractal orders on the periodic property is also elucidated.

Key words: fractal calculus, periodic solution, solitary wave, Duffing oscillator, two-scale mathematics

Introduction

This paper studies the following fractal modification of the Phi-4 equation [1-4]:

$$D_t^{2\alpha} u - D_x^{2\beta} u + m^2 u + w u^3 = 0, \quad 0 < \alpha \leq 1, \quad 0 < \beta \leq 1 \quad (1)$$

where m and w are constants, D_t^α and D_x^β are fractal derivatives with respect to t and x , respectively. Their definitions are given [5, 6]:

$$D_t^\alpha u(t_0, x) = \Gamma(1 + \alpha) \lim_{\substack{t \rightarrow t_0 + \Delta t \\ \Delta t \neq 0}} \frac{u(t, x) - u(t_0, x)}{(t - t_0)^\alpha} \quad (2)$$

$$D_x^\beta u(t, x_0) = \Gamma(1 + \beta) \lim_{\substack{x \rightarrow x_0 + \Delta x \\ \Delta x \neq 0}} \frac{u(t, x) - u(t, x_0)}{(x - x_0)^\beta} \quad (3)$$

The chain rule works for the fractal derivatives:

$$D_t^{2\alpha} = D_t^\alpha (D_t^\alpha) \quad (4)$$

$$D_x^{2\beta} = D_x^\beta (D_x^\beta) \quad (5)$$

When $\alpha = \beta = 1$, the classic Phi-4 equation is obtained:

$$u_{tt} - u_{xx} + m^2 u + w u^3 = 0 \quad (6)$$

The Phi-4 equation can model many non-linear phenomena arising in optics, thermal science, nanofluid, and non-linear vibration [1-4]. Much literature focused on its solitary solutions, while its periodic solution was rarely studied. Furthermore, the traditional Phi-4 equation cannot figure out the effect of porous structure on the solution property, and a fractal modifica-

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tion is much needed. Now the fractal differential models can reveal many interesting properties which the traditional models cannot reveal. Now fractal calculus has witnessed various applications in various fields [7-20].

Two-scale transform

The two-scale transform [21-23] is to convert eq. (1) on a small scale to an differential equation on a large-scale:

$$T = t^\alpha \quad (7)$$

$$X = x^\beta \quad (8)$$

We can convert eq. (1) into the following:

$$\frac{\partial^2 u}{\partial T^2} - \frac{\partial^2 u}{\partial X^2} + m^2 u + wu^3 = 0 \quad (9)$$

Introducing a complex variable ζ :

$$u(T, X) = U(\zeta), \quad \zeta = c(X - vT) \quad (10)$$

Equation (9) becomes:

$$c^2 v^2 U'' - c^2 U'' + m^2 U + wU^3 = 0 \quad (11)$$

where c and v are constants.

We assume the initial conditions:

$$U(0) = A, U'(0) = 0 \quad (12)$$

Periodic solution

This section applies He's frequency formulation [24-26] to figure out the periodic property of the fractal Phi-4 equation. We re-write eq. (11):

$$U'' + F(U) = 0 \quad (13)$$

where the prime implies the derivative with respect to ζ , $F(U)$ is given:

$$F(U) = \frac{m^2 U + wU^3}{c^2 v^2 - c^2} \quad (14)$$

According to He's frequency formulation [19, 20]:

$$\Omega^2 = F'(U) \Big|_{U=\frac{A}{2}} \quad (15)$$

where Ω is the frequency and A is the amplitude. He's frequency formulation has been caught much attention due to its simplest solution process and relatively high accuracy [27-37].

It is easy to calculate the derivative of $F(U)$ with respect to U :

$$F'(U) = \frac{m^2 + 3wU^2}{c^2 v^2 - c^2} \quad (16)$$

Using He's frequency formulation:

$$\Omega = \sqrt{F'(U) \Big|_{U=\frac{A}{2}}} = \sqrt{\frac{m^2 + 3w \frac{A^2}{4}}{c^2 v^2 - c^2}} \quad (17)$$

The criterion for the existence of a periodic solution:

$$c^2 v^2 - c^2 > 0 \tag{18}$$

or

$$v > 1 \text{ or } v < -1 \tag{19}$$

In case:

$$\frac{m^2}{c^2 v^2 - c^2} = 1 \tag{20}$$

and

$$\frac{w}{c^2 v^2 - c^2} = \varepsilon \tag{21}$$

Equation (11) becomes the standard Duffing oscillator:

$$U'' + U + \varepsilon U^3 = 0, \quad U(0) = A, \quad U'(0) = 0 \tag{22}$$

By the homotopy perturbation [38, 39] method or the variational iteration method, its frequency:

$$\Omega = \sqrt{1 + \frac{3}{4} \varepsilon A^2} \tag{23}$$

The approximate periodic solution of eq. (11):

$$U(\xi) = A \cos(\Omega \xi + \varphi) \tag{24}$$

where φ is a constant.

In view of eq. (10):

$$U(\xi) = A \cos(\Omega \xi + \varphi) \tag{25}$$

By eqs. (7) and (8), we finally obtain the periodic solution of eq. (1):

$$U(t, x) = A \cos[\Omega c(x^\beta - vt^\alpha) + \varphi] \tag{26}$$

or

$$U(t, x) = A \cos \left[\sqrt{\frac{m^2 + \frac{3w}{4} A^2}{v^2 - 1}} (x^\beta - vt^\alpha) + \varphi \right] \tag{27}$$

Solution morphology

From eq. (26), we have:

$$\frac{\partial}{\partial t} U(t, x) = \Omega c v A \alpha t^{\alpha-1} \sin[\Omega c(x^\beta - vt^\alpha) + \varphi] \tag{28}$$

and

$$\frac{\partial}{\partial x} U(t, x) = -\Omega c A \beta x^{\beta-1} \sin[\Omega c(x^\beta - vt^\alpha) + \varphi] \tag{29}$$

In case of $\alpha = \beta = 1$, we have:

$$\frac{\partial}{\partial t} U(t, x) = \Omega c v A \sin[\Omega c(x - vt) + \varphi] \tag{30}$$

$$\frac{\partial}{\partial x} U(t, x) = -\Omega c A \sin[\Omega c(x - vt) + \varphi] \tag{31}$$

This case leads to the standard periodic property.

In case of $\alpha < 1$ and $\beta < 1$, we have:

$$\frac{\partial}{\partial t} U(0, x) \rightarrow \infty \tag{32}$$

$$\frac{\partial}{\partial x} U(t, 0) \rightarrow \infty \tag{33}$$

Equation (32) sees an extremely large change of U at the initial time, while eq. (33) predicts an extremely large slope at $x = 0$.

From aforementioned analysis, we can see that the fractal orders will greatly affect the solution morphology, see fig. 1 for the case $A = m = w = 1$, $v = 2$, and $\varphi = 0$.

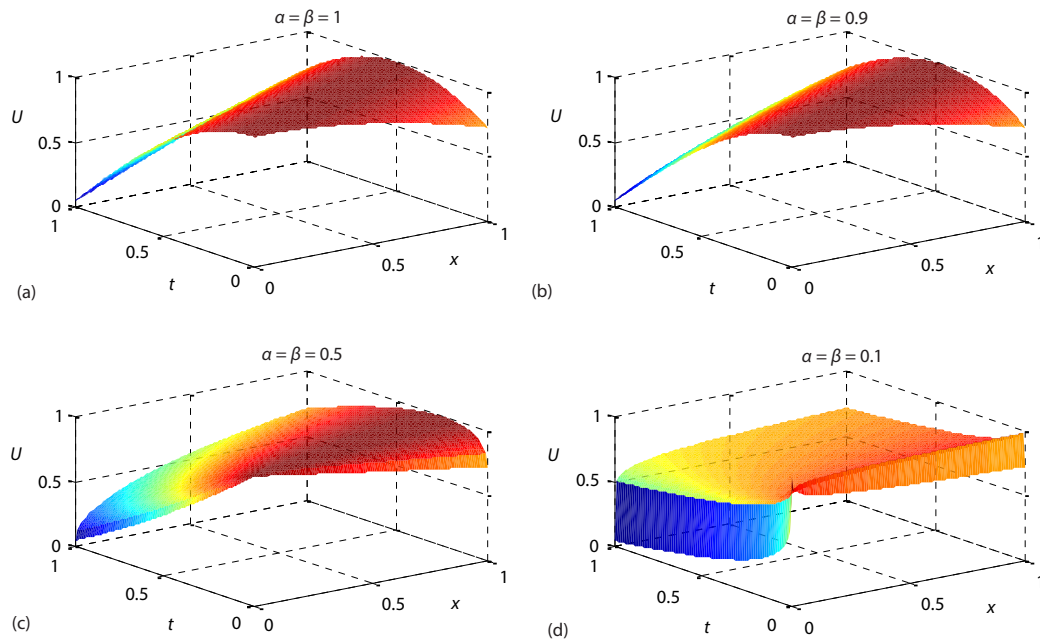


Figure 1. Solution morphology for different values of α and β

Conclusions

In this paper, we use the Hira Tariq method to transform a FPDE to a FODE and it was very effective, however, FODE remains difficult to solve. Fortunately, there is He’s frequency effective method to solve it. As Dr. Ji-Huan He has been emphasized, the simpler is the better. In engineering applications, a fast and effective estimation of a nonlinear vibration problem is very needed, and He’s frequency formation becomes a universal tool for this purpose.

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