

## NEW OPTICAL EXPLICIT PLETHORA OF THE RESONANT SCHRODINGER'S EQUATION VIA TWO RECENT COMPUTATIONAL SCHEMES

by

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*This research paper investigates the computational solutions of the resonant schrödinger's equation. The modified Khater method and Adomian decomposition method are applied for construct new analytical traveling and semi-analytical wave solutions. This model describes the pulse phenomena and studied in non-linear optics. For further illustration of our obtained solutions, some distinct types of sketches are given.*

**Key words:** *(3+1) resonant Schrodinger's equation, modified Khater method, Adomian decomposition method, exact and approximate solutions*

### Introduction

Marvellous applications in the field of science and engineering have been formulating in non-linear PDE with an-integer order or fractional order [1-15]. Thus, many researchers in different fields have been focusing on studying the exact travelling and solitary wave solutions. These solutions are used to discover more physical, dynamical behaviour, and the nature of non-linear problems and explain the different scientific non-linear phenomena. For achieving to this fundamental goal, many computational and numerical schemes have been being formulated such as the improved and novel expansion method [1-3], extended simplest equation method [4-7], first integral method [8, 9], Khater method which is derived by Khater in 2017 [10-13], the exp -expansion method [14, 15], the Kudryashov method [16], the Jacobi elliptic functions [17, 18], the modified Khater method, [19-22], the generalized Kudryashov method [23, 24], the new extended direct algebraic method [25], the functional variable method [26], Adomian decomposition method [27-29], and the sub equation method [30].

In this paper, with the availability of symbolic computation packages, the resonant non-linear Schrodinger's equation (RNLS) is investigated by employing the modified Khater

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method and Adomian decomposition method to construct analytical and semi-analytical solutions [31-41]. The (3+1)-D RNLS model is given:

$$iQ_t + r_1 \nabla^2 Q + r_2 |Q|^2 S + r_2 \left( \frac{\nabla^2 |Q|}{|Q|} \right) S = 0 \quad (1)$$

where

$$Q = Q(x, y, z, t), \quad \nabla = \frac{\partial}{\partial x} + \frac{\partial}{\partial y} + \frac{\partial}{\partial z}, \quad x, y, z \in \mathcal{R}$$

Employing the following traveling wave transformation:

$$Q(x, y, z, t) = \mathcal{G}(\mathcal{h}) e^{i\varphi}$$

where

$$\mathcal{h} = \cos(s_1)x + \cos(s_2)y + \cos(s_3)z + s_4 t, \quad \varphi = \left\{ s_5 \left[ \cos(s_1)x + \cos(s_2)y + \cos(s_3)z \right] + s_6 t \right\}$$

yields

$$\mathcal{L}_1 \mathcal{G}'' - \mathcal{L}_2 \mathcal{G} + r_2 \mathcal{G}^3 = 0 \quad (2)$$

where

$$\mathcal{L}_1 = \wp(r_3 + r_1), \quad \mathcal{L}_2 = (\wp r_1 s_5 + s_6), \quad s_4 = -2\wp r_1 s_5, \quad \wp = \cos(s_1)^2 + \cos(s_2)^2 + \cos(s_3)^2$$

Balancing the highest order derivative term and non-linear term in eq. (2), leads:

$$\mathcal{G}'', \mathcal{G}^3 \Rightarrow N + 2 = 3N \Rightarrow N = 1$$

### Application

In this part, we apply the modified Khater method and Adomian decomposition method [42-44] to the (3+1) RNLS equation.

#### Modified Khater method

According to the general solutions suggested by the method, we get the general solution of eq. (2) in the form:

$$\mathcal{G}(\mathcal{h}) = a_0 + \sum_{i=1}^N a_i \mathcal{K}^{i f(\mathcal{h})} + \sum_{i=1}^N b_i k^{-i f(\mathcal{h})} = a_0 + k^{f(\mathcal{h})} a_1 + k^{-f(\mathcal{h})} b_1 \quad (3)$$

where  $a_0, a_1, b_1$ , and  $k$  are arbitrary constants while  $f(\mathcal{h})$  satisfies the auxiliary equation

$$f'(\mathcal{h}) = \frac{1}{\text{Ln}(\mathcal{K})} \left( \varrho \mathcal{K}^{-f(\mathcal{h})} + \chi + \delta \mathcal{K}^{f(\mathcal{h})} \right), \quad \text{where } \varrho, \chi, \delta \text{ are arbitrary constants}$$

Substituting eq. (3) and its derivatives into eq. (2). Collecting all terms of the same power of  $\mathcal{K}^{f(\mathcal{h})}$ . Solving the obtained algebraic system by any computer software program, leads:

Family 1.

$$a_1 \rightarrow \frac{2\delta a_0}{\chi}, \quad b_1 \rightarrow 0, \quad \mathcal{L}_1 \rightarrow -\frac{2a_0^2 r_2}{\chi^2}, \quad \mathcal{L}_2 \rightarrow \frac{-4\delta \varrho a_0^2 r_2 + \chi^2 a_0^2 r_2}{\chi^2}$$

Family 2.

$$a_1 \rightarrow 0, b_1 \rightarrow \frac{2\varrho a_0}{\chi}, \mathcal{L}_1 \rightarrow -\frac{2a_0^2 r_2}{\chi^2}, \mathcal{L}_2 \rightarrow \frac{-4\delta\varrho a_0^2 r_2 + \chi^2 a_0^2 r_2}{\chi^2}$$

According to the value of parameters in Family 1 and 2 the explicit wave solutions of eq. (1) are given:

– for  $\beta^2 - 4\alpha\sigma < 0, \sigma \neq 0$

$$Q_{1,1}(x, y, x, t) = e^{i\varphi} \frac{\sqrt{4\delta\varrho - \chi^2} a_0 \tan\left[\frac{1}{2} \hbar \sqrt{4\delta\varrho - \chi^2}\right]}{\chi} \quad (4)$$

$$Q_{1,2}(x, y, x, t) = e^{i\varphi} \frac{\sqrt{4\delta\varrho - \chi^2} \cot\left[\frac{1}{2} \hbar \sqrt{4\delta\varrho - \chi^2}\right] a_0}{\chi} \quad (5)$$

– for  $\beta^2 - 4\alpha\sigma > 0, \sigma \neq 0$

$$Q_{1,3}(x, y, x, t) = -e^{i\varphi} \frac{\sqrt{-4\delta\varrho + \chi^2} a_0 \tanh\left[\frac{1}{2} \hbar \sqrt{-4\delta\varrho + \chi^2}\right]}{\chi} \quad (6)$$

$$Q_{1,4}(x, y, x, t) = -e^{i\varphi} \frac{\sqrt{-4\delta\varrho + \chi^2} \coth\left[\frac{1}{2} \hbar \sqrt{-4\delta\varrho + \chi^2}\right] a_0}{\chi} \quad (7)$$

– for  $\chi = \delta = \kappa$  and  $\varrho = 0$

$$Q_{1,5}(x, y, x, t) = -a_0 e^{i\varphi} \coth\left[\frac{\kappa}{2} \hbar\right] \quad (8)$$

– for  $\varrho = 0, \chi \neq 0$ , and  $\delta \neq 0$

$$Q_{1,6}(x, y, x, t) = -a_0 e^{i\varphi} \left(1 + \frac{4}{-2 + e^{\hbar z} \delta}\right) \quad (9)$$

– for  $\chi^2 - 4\varrho\delta = 0$

$$Q_{1,7}(x, y, x, t) = -a_0 e^{i\varphi} \left(-1 + \hbar \frac{4\delta\varrho(2 + \hbar\chi)}{\chi^3}\right) \quad (10)$$

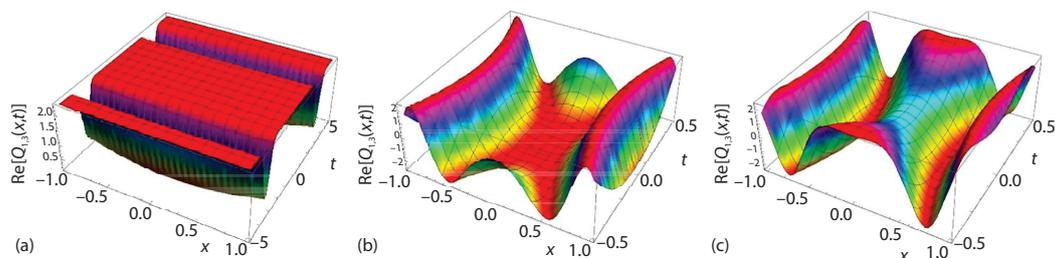
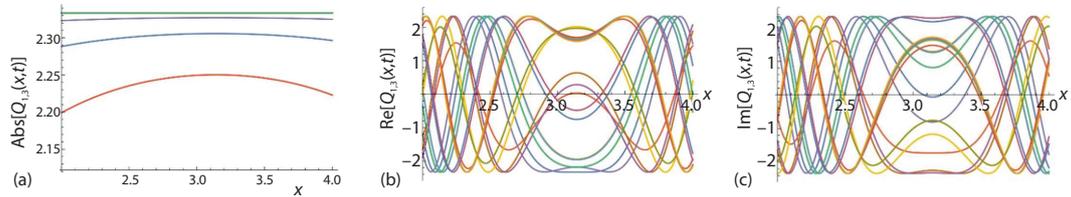
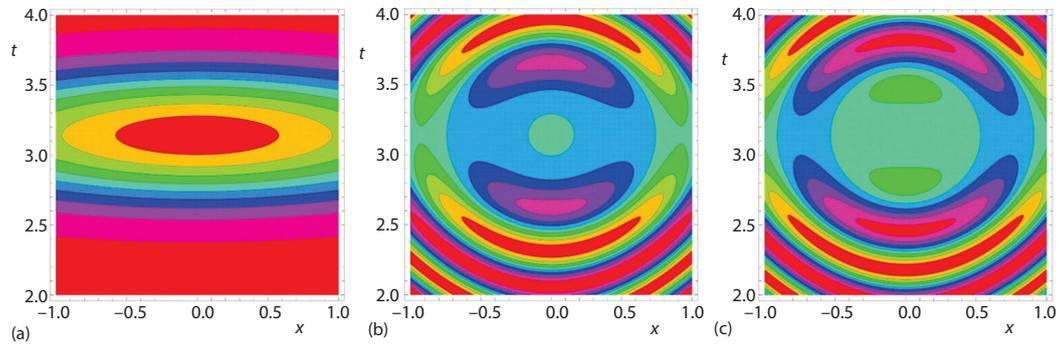


Figure 1. Solitary wave solutions of absolute, real, and imaginary value of eq. (6) in 3-D for  $s_1 = 10, s_2 = 12, s_3 = 13, s_4 = 14, s_5 = 15, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$



**Figure 2. Solitary wave solutions of absolute, real, and imaginary value of eq. (6) in 2-D for  $s_1 = 10, s_2 = 12, s_3 = 13, s_4 = 14, s_5 = 15, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$**



**Figure 3. Solitary wave solutions of absolute, real, and imaginary value of eq. (6) in contour plot for  $s_1 = 10, s_2 = 12, s_3 = 13, s_4 = 14, s_5 = 15, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$**

According to the value of parameters in *Family 2*, we get the solitary wave solutions of eq. (1):

– for  $\beta^2 - 4\alpha\sigma < 0, \sigma \neq 0$

$$Q_{2,1}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 - \frac{4\delta\varrho}{\chi^2 - \chi\sqrt{4\delta\varrho - \chi^2} \tan\left[\frac{1}{2}h\sqrt{4\delta\varrho - \chi^2}\right]} \right) \tag{11}$$

$$Q_{2,2}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 - \frac{4\delta\varrho}{\chi^2 - \chi\sqrt{4\delta\varrho - \chi^2} \cot\left[\frac{1}{2}h\sqrt{4\delta\varrho - \chi^2}\right]} \right) \tag{12}$$

– for  $\beta^2 - 4\alpha\sigma > 0, \sigma \neq 0$

$$Q_{2,3}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 - \frac{4\delta\varrho}{\chi^2 + \chi\sqrt{-4\delta\varrho + \chi^2} \tanh\left[\frac{1}{2}h\sqrt{-4\delta\varrho + \chi^2}\right]} \right) \tag{13}$$

$$Q_{2,4}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 - \frac{4\delta\varrho}{\chi^2 + \chi\sqrt{-4\delta\varrho + \chi^2} \coth\left[\frac{1}{2}h\sqrt{-4\delta\varrho + \chi^2}\right]} \right) \tag{14}$$

– for  $\chi = \varrho/2 = \kappa$  and  $\delta = 0$

$$Q_{2,5}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 + \frac{4}{-2 + e^{\kappa h}} \right) \quad (15)$$

– for  $\delta = 0, \chi \neq 0$ , and  $\varrho \neq 0$

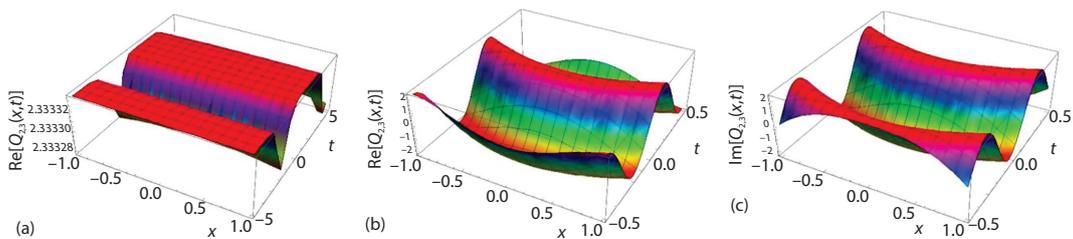
$$Q_{2,6}(x, y, x, t) = a_0 e^{i\varphi} \left( 1 - \frac{2\varrho}{\varrho - e^{h\chi}} \right) \quad (16)$$

– for  $\chi^2 - 4\varrho\delta = 0$

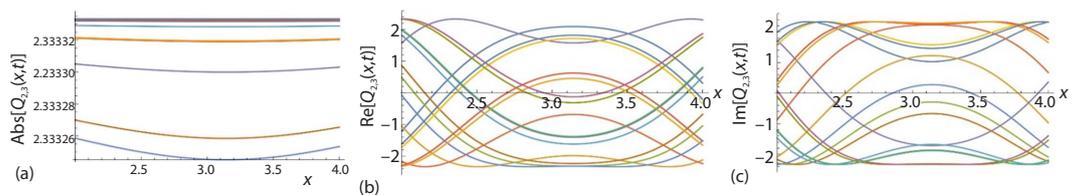
$$Q_{2,7}(x, y, x, t) = \frac{2a_0 e^{i\varphi}}{2 + h\chi} \quad (17)$$

where

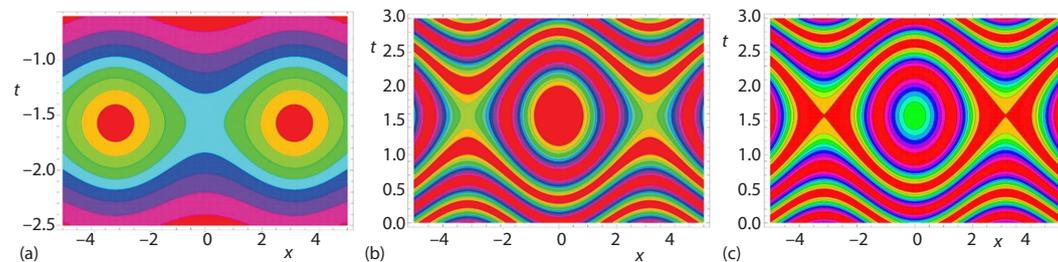
$$h = \cos(s_1)x + \cos(s_2)y + \cos(s_3)z + s_4 t, \quad \varphi = \{s_5 [\cos(s_1)x + \cos(s_2)y + \cos(s_3)z] + s_6 t\}$$



**Figure 4.** Solitary wave solutions of absolute, real, and imaginary value of eq. (13) in 3-D for  $s_1 = 5, s_2 = 6, s_3 = 7, s_4 = 4, s_5 = 10, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$



**Figure 5.** Solitary wave solutions of absolute, real, and imaginary value of eq. (13) in 2-D for  $s_1 = 5, s_2 = 6, s_3 = 7, s_4 = 4, s_5 = 10, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$



**Figure 6.** Solitary wave solutions of absolute, real, and imaginary value of eq. (13) in contour plot for  $s_1 = 5, s_2 = 6, s_3 = 7, s_4 = 4, s_5 = 10, z = 9, y = 8, a_0 = 7, \delta = 1, \varrho = 2, \chi = 3$

*Adomian decomposition method*

Applying the Adomian decomposition method on eq. (2) enables rewriting it to be in the form:

$$L\mathcal{G}(\mathcal{h}) + R\mathcal{G}(\mathcal{h}) + N\mathcal{G}(\mathcal{h}) = 0 \quad (18)$$

where ( $L$ ,  $R$ , and  $N$ ) represent a differential operator, a linear operator and non-linear term, respectively.

Using the inverse operator ( $L^{-1}$ ) on (18), we get:

$$\sum_{i=0}^{\infty} \mathcal{G}_i(\mathcal{h}) = \mathcal{G}(0) + \mathcal{G}'(0)\mathcal{h} + \frac{\mathcal{L}_2}{\mathcal{L}_1} L^{-1} \left[ \sum_{i=0}^{\infty} \mathcal{G}_i(\mathcal{h}) \right] - \frac{\mathcal{r}_1}{\mathcal{L}_1} L^{-1} \left( \sum_{i=0}^{\infty} A_i \right) \quad (19)$$

Under the following condition [ $\delta = 1$ ,  $\varrho = 2$ ,  $\chi = 3$ ,  $a_0 = 5$ ,  $r_2 = 4$ ] on eq. (6), we get:

$$\mathcal{G}_0 = -\frac{5\mathcal{h}}{6} \quad (20)$$

$$\mathcal{G}_1 = \frac{5\mathcal{h}^3}{72} - \frac{\mathcal{h}^5}{192} \quad (21)$$

$$\mathcal{G}_2 = -\frac{\mathcal{h}^5}{576} + \frac{\mathcal{h}^7}{16128} - \frac{25\mathcal{h}^8}{64512} + \frac{\mathcal{h}^{10}}{55296} \quad (22)$$

$$\mathcal{G}_3 = \frac{31\mathcal{h}^7}{48384} - \frac{67\mathcal{h}^9}{1161216} + \frac{5\mathcal{h}^{10}}{2322432} + \frac{\mathcal{h}^{11}}{337920} - \frac{\mathcal{h}^{12}}{14598144} - \frac{\mathcal{h}^{13}}{12779520} \quad (23)$$

According to eqs. (20)-(23), we get an approximate solution of eq. (2):

$$\begin{aligned} \mathcal{G}(\mathcal{h}) = & -\frac{5\mathcal{h}}{6} + \frac{5\mathcal{h}^3}{72} - \frac{\mathcal{h}^5}{144} + \frac{17\mathcal{h}^7}{24192} - \frac{25\mathcal{h}^8}{64512} - \frac{67\mathcal{h}^9}{1161216} + \\ & + \frac{47\mathcal{h}^{10}}{2322432} + \frac{\mathcal{h}^{11}}{337920} - \frac{\mathcal{h}^{12}}{14598144} - \frac{\mathcal{h}^{13}}{12779520} + \dots \end{aligned} \quad (24)$$

In tab. 1, we discuss the exact and approximate solutions of the (3+1) RNLS equation show the value of the absolute error between them.

**Table 1. Shows for increasing the value  $\mathcal{h}$ , the absolute error increases gradually; that means the Adomian decomposition method gives more accurate solutions for the values near to zero**

Value of $\mathcal{h}$	Exact solution	Approximate solution	Absolute error
0.01	0.00833326	0.00833326	$1.73472 \cdot 10^{-18}$
0.02	0.0166661	0.0166661	$6.93889 \cdot 10^{-18}$
0.03	0.0249981	0.0249981	$2.56739 \cdot 10^{-16}$
0.04	0.0333289	0.0333289	$2.5327 \cdot 10^{-15}$
0.05	0.041658	0.041658	$1.51129 \cdot 10^{-14}$
0.06	0.049985	0.049985	$6.4948 \cdot 10^{-14}$
0.07	0.0583095	0.0583095	$2.22801 \cdot 10^{-13}$
0.08	0.0666311	0.0666311	$6.4812 \cdot 10^{-13}$
0.09	0.0749494	0.0749494	$1.66225 \cdot 10^{-12}$
0.1	0.083264	0.083264	$3.85979 \cdot 10^{-12}$

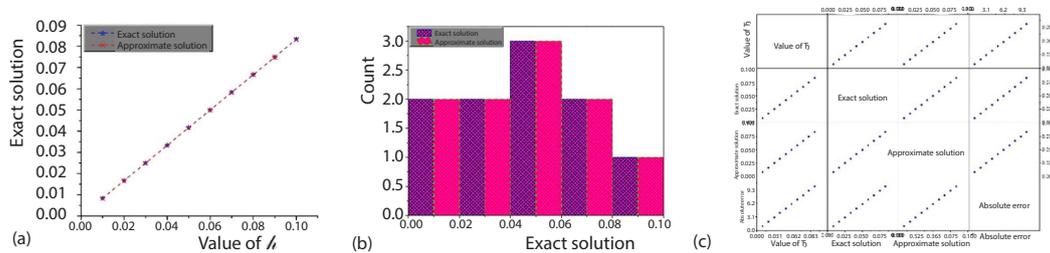


Figure 7. Exact and approximate wave solutions according to the shown values in tab. 1, show the accuracy of the our obtained solutions.

### Conclusion

In this paper, the (3+1) RNLS equation have been investigated via the modified Khater method. Some new distinct types of computational solutions have been obtained. These solutions have been used to evaluate the initial and boundary conditions of the model. Furthermore, the Adomian decomposition method have been applied for construct the semi-analytical wave solutions based on these conditions, Some solitary and approximate solutions are sketched to investigate more of the physical properties of this model figs. 1-7. The performance of both methods shows useful and powerful in studying many of non-linear partial differential equations.

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