# AN ANALYTICAL SOLUTION FOR SOLVING A NEW GENERAL FRACTIONAL-ORDER MODEL FOR WAVE IN MINING ROCK

by

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In this paper, we consider the general fractional-order derivatives of the Liouville-Sonine-Caputo and Liouville-Sonine type containing the Lorenzo-Hartley kernel. A general fractional-order model for the wave equation with the analytical solution is discussed in detail. The general fractional-order formula is accurate and efficient for description of the complex, power-law and memory behaviors for the mining rock.

Key words: general fractional-order derivative, Lorenzo-Hartley kernel, wave equation, analytical solution, mining rock

### Introduction

The linear and non-linear models for the wave motion have been successfully observed an investigated in the mining rock [1-3]. For example, the linear Guyer-McCall-Boitnott equation of the motion describe 1-D wave propagation in the mining rock was reported [4]. The linear McCall equation of the motion describe 1-D wave propagation in the mining rock was considered [5]. The linear McCall-Guyer wave equation of the motion in the mining rock was presented [6].

Recently, the general fractional-order derivatives of the Liouville-Sonine and Liouville-Sonine-Caputo types containing the Lorenzo-Hartley kernel was proposed [7]. The main goal of the paper is to the linear McCall equation of the motion describe 1-D wave propagation in the mining rock in the sense of the general fractional-order derivative of the Liouville-Sonine-Caputo type.

## General fractional-order derivatives containing the Lorenzo-Hartley kernel

In this section, we investigate the general fractional-order calculus, inclusive of the general fractional-order derivatives and integrals [7].

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General fractional-order integral operators containing the Lorenzo-Hartley kernel

The left-sided general fractional-order integral operator containing the Lorenzo-Hartley kernel are expressed [7]:

$${}_{H}I_{a+}^{\alpha,\gamma}k(t) = \int_{a}^{t} \aleph_{\alpha} \left[ -\gamma (t-\tau)^{\alpha} \right] k(\tau) d\tau$$
 (1)

and the right-sided general fractional-order integral operator containing the Lorenzo-Hartley kernel:

$$_{\mathrm{H}}I_{b-}^{\alpha,\gamma}k\left(t\right) = \int_{t}^{b} \aleph_{\alpha}\left[-\gamma\left(\tau-t\right)^{\alpha}\right]k\left(\tau\right)\mathrm{d}\tau\tag{2}$$

respectively, where the Lorenzo-Hartley function is defined [8]:

$$\aleph_{\alpha} \left( \gamma t^{\alpha} \right) = \sum_{\kappa=0}^{\infty} \frac{\gamma^{\kappa} t^{(\kappa+1)\alpha-1}}{\Gamma \left[ (\kappa+1)\alpha \right]} \tag{3}$$

If the Laplace transform operator is defined [7]:

$$\wp\left[\iota(t)\right] = \iota(s) = \int_{0}^{\infty} e^{-st} \iota(t) dt$$
 (4)

then the Laplace transform of eq. (3) can be given [7]:

$$\wp\left\{\aleph_{\alpha}\left(\gamma t^{\alpha}\right)\right\} = s^{-\alpha}\left(1 - \gamma s^{-\alpha}\right)^{-1} \left(\left|\nu s^{-\alpha}\right| < 1\right) \tag{5}$$

When a = 0, the general fractional-oder integral operator containing the Lorenzo-Hartley kernel are expressed [7]:

$${}_{\mathrm{H}}I_{0+}^{\alpha,\gamma}k(t) = \int_{\alpha}^{t} \aleph_{\alpha} \left[ -\gamma(t-\tau)^{\alpha} \right] k(\tau) d\tau \tag{6}$$

General fractional-order derivatives of the Liouville-Sonine type

The left-sided general fractional-order derivative of the Liouville-Sonine type containing the Lorenzo-Hartley kernel is defined [7]:

$${}_{\mathrm{H}}^{\mathrm{LS}} \mathbf{D}_{a+}^{\alpha,\gamma} k(t) = {}_{\mathrm{H}} I_{a+}^{\alpha,\gamma} \left[ k^{(1)}(t) \right] = \int_{a}^{t} \aleph_{\alpha} \left[ -\gamma (t-\tau)^{\alpha} \right] k^{(1)}(\tau) \, \mathrm{d}\tau \tag{7}$$

and the right-sided general fractional-order derivative of the Liouville-Sonine type containing the Lorenzo-Hartley kernel [7]:

$${}_{\mathrm{H}}^{\mathrm{LS}} \mathbf{D}_{b^{-}}^{\alpha,\lambda} f\left(t\right) = {}_{\mathrm{H}} I_{b^{-}}^{\alpha,\gamma} \left[-f^{(1)}\left(t\right)\right] = -\int_{0}^{b} \aleph_{\alpha} \left[-\gamma \left(\tau - t\right)^{\alpha}\right] f^{(1)}\left(\tau\right) \mathrm{d}\tau \tag{8}$$

When a = 0, eq. (5) can be re-written [7]:

$${}_{\mathrm{H}}^{\mathrm{LS}} \mathbf{D}_{0+}^{\alpha,\gamma} k(t) = {}_{\mathrm{H}} I_{0+}^{\alpha,\gamma} \left[ k^{(1)}(t) \right] = \int_{0}^{t} \mathbf{\aleph}_{\alpha} \left[ -\gamma (t-\tau)^{\alpha} \right] k^{(1)}(\tau) \, \mathrm{d}\tau \tag{9}$$

and the Laplace transform of eq. (9) [7]:

$$\mathcal{O}\left\{ {}_{H}^{LS} \mathbf{D}_{0+}^{\alpha,\gamma} k\left(t\right) \right\} = s^{-\alpha} \left(1 - \gamma s^{-\alpha}\right)^{-1} \left[ sk\left(s\right) - k\left(0\right) \right]$$
(10)

General fractional-order derivatives of the Liouville-Sonine-Caputo type

The left-sided general fractional-order derivative of the Liouville-Sonine-Caputo type containing the Lorenzo-Hartley kernel is defined [7]:

$${}^{\mathrm{LSC}}_{\mathrm{H}} \mathbf{D}_{a+}^{\alpha,\kappa,\gamma} k\left(t\right) = {}_{\mathrm{H}} I_{a+}^{\alpha,\gamma} \left[ k^{(\kappa)}\left(t\right) \right] = \int_{0}^{t} \aleph_{\alpha} \left[ -\gamma \left(t-\tau\right)^{\alpha} \right] k^{(\kappa)}\left(\tau\right) \mathrm{d}\tau$$

$$\tag{11}$$

and the right-sided general fractional-order derivative of the Liouville-Sonine-Caputo type containing the Lorenzo-Hartley kernel [7]:

$${}^{\mathrm{LSC}}_{\mathrm{H}} \mathbf{D}_{b^{-}}^{\alpha,\lambda} k\left(t\right) = {}_{\mathrm{H}} I_{b^{-}}^{\alpha,\gamma} \left[ \left(-1\right)^{\kappa} k^{(\kappa)}\left(t\right) \right] = \left(-1\right)^{\kappa} \int_{t}^{b} \mathbf{N}_{\alpha} \left[ -\gamma \left(\tau - t\right)^{\alpha} \right] k^{(\kappa)} \left(\tau\right) \mathrm{d}\tau \tag{12}$$

When a = 0, eq. (8) becomes [7]:

$${}^{\mathrm{LSC}}_{\mathrm{H}} \mathbf{D}_{0+}^{\alpha,\kappa,\gamma} k\left(t\right) = {}_{\mathrm{H}} I_{0+}^{\alpha,\gamma} \left[ k^{(\kappa)}\left(t\right) \right] = \int_{0}^{t} \aleph_{\alpha} \left[ -\gamma \left(t - \tau\right)^{\alpha} \right] k^{(\kappa)}\left(\tau\right) \mathrm{d}\tau \tag{13}$$

and the Laplace transform of eq. (13) [7]:

$$\mathcal{O}\left\{ {}^{\mathrm{LSC}}_{\mathrm{H}} \mathbf{D}^{\alpha,\kappa,\gamma}_{0+} k\left(t\right) \right\} = s^{-\alpha} \left(1 - \gamma s^{-\alpha}\right)^{-1} \left[ s^{\kappa} k\left(s\right) - \sum_{j=1}^{\kappa} s^{\kappa-j} k^{(j-1)}\left(0\right) \right]$$

$$\tag{14}$$

General fractional-order integrals

The left-sided general fractional-order integral is defined [7]:

$${}_{\rm H}^{\rm R}I_{a+}^{\alpha,\kappa,\gamma}k\left(t\right) = \int_{-\infty}^{\infty} \left(t-\tau\right)^{\kappa-\alpha-1}E_{\alpha,\kappa-\alpha}^{-1}\left[-\gamma\left(t-\tau\right)^{\alpha}\right]k\left(\tau\right)\mathrm{d}\tau\tag{15}$$

and the right-sided general fractional-order integral [7]:

$${}_{\rm H}^{\rm R}I_{b-}^{\alpha,\kappa,\gamma}k(t) = \left(-1\right)^{\kappa}\int_{0}^{t} \left(\tau - t\right)^{\kappa-\alpha-1}E_{\alpha,\kappa-\alpha}^{-1}\left[-\gamma(\tau - t)^{\alpha}\right]k(\tau)d\tau \tag{16}$$

When a = 0, eq. (15) can be given [7]:

$${}_{\mathrm{H}}^{\mathrm{R}}I_{0+}^{\alpha,\kappa,\gamma}k\left(t\right) = \int_{0}^{t} \left(t-\tau\right)^{\kappa-\alpha-1}E_{\alpha,\kappa-\alpha}^{-1}\left[-\gamma\left(t-\tau\right)^{\alpha}\right]k\left(\tau\right)d\tau\tag{17}$$

with the Laplace transform [7]:

$$\mathscr{O}\left\{ {}_{H}^{R}I_{0+}^{\alpha,\kappa,\gamma}k(t)\right\} = s^{\alpha-\kappa}\left(1+\gamma s^{-\alpha}\right)k(s) \tag{18}$$

For more relations between the general fractional-order derivatives and integrals, [7].

# A general fractional-order wave model with the analytical solution

We now consider the general fractional-order wave model within the general fractional-order derivative of the Liouville-Sonine-Caputo type containing the Lorenzo-Hartley kernel:

$${}^{LSC}_{H} \partial_{0+}^{\alpha,2,\gamma} \Pi(x,t) = \frac{\partial^{2} \Pi(x,t)}{\partial x^{2}} + \sin(\pi x) \quad (0 < x < 1, t > 0)$$

$$\tag{19}$$

with the initial and boundary conditions:

$$\Pi(0,t) = 0$$
,  $\Pi(1,t) = 0$ ,  $\Pi(x,0) = 0$ , and  $\frac{\partial \Pi(x,0)}{\partial x} = 0$  (20)

where the general fractional-order partial derivative is defined:

$${}^{LSC}_{H}\partial^{\alpha,2,\gamma}_{0+}\Pi(x,t) = {}_{H}I^{\alpha,\gamma}_{0+}\left[\Pi^{(2)}(x,t)\right] = \int_{0}^{t} \aleph_{\alpha}\left[-\gamma(t-\tau)^{\alpha}\right]\Pi^{(2)}(x,\tau)d\tau \tag{21}$$

Making use of the Laplace transform of eq. (19), we obtain:

$$\frac{\partial^2 \Pi(x,s)}{\partial x^2} - s^{2-\alpha} \left( 1 + \gamma s^{-\alpha} \right)^{-1} \Pi(x,s) = \frac{\sin(\pi x)}{s}$$
 (22)

Suppose that

$$\Pi(x,s) = \mu_1 e^{x\sqrt{s^{2-\alpha}(1+\gamma s^{-\alpha})^{-1}}} + \mu_2 e^{-x\sqrt{s^{2-\alpha}(1+\gamma s^{-\alpha})^{-1}}} + \mu_3 \sin(\pi x) + \mu_4 \cos(\pi x)$$
(23)

where  $\mu_1$ ,  $\mu_2$ ,  $\mu_3$ , and  $\mu_4$  are the constants, then:

$$\left[-\pi^2 - s^{2-\alpha} \left(1 + \gamma s^{-\alpha}\right)^{-1}\right] \left[\mu_3 \sin\left(\pi x\right) + \mu_4 \cos\left(\pi x\right)\right] = \frac{\sin\left(\pi x\right)}{s}$$
 (24)

which leads:

$$\mu_3 = \frac{1}{s} \frac{1}{-\pi^2 - s^{2-\alpha} \left(1 + \gamma s^{-\alpha}\right)^{-1}}, \ \mu_4 = 0$$
 (25)

Thus, we obtain:

$$\Pi(x,s) = \mu_1 e^{x\sqrt{s^{2-\alpha}(1+\gamma s^{-\alpha})^{-1}}} + \mu_2 e^{-x\sqrt{s^{2-\alpha}(1+\gamma s^{-\alpha})^{-1}}} + \frac{1}{s} \frac{1}{-\pi^2 - s^{2-\alpha}(1+\gamma s^{-\alpha})^{-1}} \sin(\pi x)$$
 (26)

which from eq. (20) implies:

$$\mu_1 = 0, \ \mu_2 = 0$$

and

$$\Pi(x,s) = \frac{1}{s} \frac{1}{-\pi^2 - s^{2-\alpha} \left(1 + \gamma s^{-\alpha}\right)^{-1}} \sin(\pi x)$$
 (27)

With the use:

$$\frac{1}{1 + \frac{1}{\pi^2} s^{2-\gamma} \left(1 + \nu s^{-\alpha}\right)^{-1}} =$$

$$= \sum_{i=0}^{n} \left( -\frac{1}{\pi^2} \right)^i s^{(2-\gamma)i} \left( 1 + \nu s^{-\gamma} \right)^{-i}, \ \left| \frac{1}{\pi^2} s^{2-\gamma} \left( 1 + \nu s^{-\gamma} \right)^{-1} \right| < 1$$
 (28)

eq. (27) can be re-written:

$$\Pi(x,s) = -\frac{\sin(\pi x)}{\pi^2} \sum_{i=0}^{n} \left(-\frac{1}{\pi^2}\right)^{i} s^{(2-\alpha)i-1} \left(1 + \gamma s^{-\alpha}\right)^{-i}$$
(29)

Therefore, the analytical solution of eq. (19) takes the form:

$$\Pi(x,t) = -\frac{\sin(\pi x)}{\pi^2} \sum_{i=0}^{n} \left(-\frac{1}{\pi^2}\right)^{j} t^{(2-\alpha)j} \left(1 + \gamma s^{-\alpha}\right)^{-j} E_{\alpha,\alpha i-2i+1}^{i} \left(-\gamma t^{\alpha}\right)$$
(30)

### Conclusion

We have investigated the general fractional-order calculus containing the Loren-zo-Hartley kernel. A general fractional-order model for the wave equation containing the gener-

al fractional-order derivative of the Liouville-Sonine-Caputo type containing the Lorenzo-Hart-ley kernel with the analytical solution is presented with the use of the Laplace transform. The result is accurate and efficient for designing of the complex, power-law and memory behaviors for the mining rock.

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### Nomenclature

t – time co-ordinate, [m] x – space co-ordinate, [m]  $\alpha$  – fractional order, [–]

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