

## A DIRECT ALGORITHM OF EXP-FUNCTION METHOD FOR NON-LINEAR EVOLUTION EQUATIONS IN FLUIDS

by

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*In this paper, a direct algorithm of the exp-function method is proposed for exactly solving non-linear evolution equations. To illustrate the validity and advantages of the algorithm, the Korteweg-de Vries and Jimbo-Miwa equations are considered. As a result, exact solutions are obtained. It is shown that the exp-function method with the direct algorithm provides a simpler but effective mathematical tool for constructing exact solutions of non-linear evolution equations in fluids.*

Key words: *exp-function method, exact solution, Korteweg-de Vries equation, Jimbo-Miwa equation*

### Introduction

Searching for exact solutions of non-linear evolution equations (NLEE) plays an important role in the study of non-linear physical phenomena in fluid dynamics. In 2006, He and Wu [1] proposed the so-called exp-function method to solve non-linear wave equations. It is shown that the exp-function method and its improvements are available for many NLEE such as those in [2-10]. This is due to the method possesses a more general ansatz:

$$u = \frac{\sum_{n=0}^g a_n e^{n\xi}}{\sum_{m=0}^p b_m e^{m\xi}}, \quad \xi = \sum_{i=1}^s k_i x_i - \omega t \quad (1)$$

supposed for a given NLEE with independent variables  $t, x_1, x_2, \dots, x_s$  and dependent variable  $u$ :

$$F(u, u_t, u_{x_1}, u_{x_2}, \dots, u_{x_m}, u_{x_1 t}, u_{x_2 t}, \dots, u_{x_s t}, u_{tt}, u_{x_1 x_1}, u_{x_2 x_2}, \dots, u_{x_s x_s}, \dots) = 0 \quad (2)$$

Here  $a_n$  and  $b_m$  are undetermined constants,  $f, p, g$ , and  $q$  can be determined by balancing the highest order linear term with the highest order non-linear term in eq. (2). In general, the final solution does not strongly depend on the choices of values of  $f, p, g, q$ , and  $f=p=g=q=1$  is the simplest choice. In this paper, we would like to introduce the following general approach:

$$u = \sum_{j=1}^n u_j(t, x_1, x_2, \dots, x_s) \varphi^{j-n}, \quad \varphi = 1 + e^\xi \quad (3)$$

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where non-negative integer  $n$  is determined by balancing the highest order linear term with the highest order non-linear term in eq. (2),  $u_j = u_j(t, x_1, x_2, \dots, x_s), j = 1, 2, \dots, s$ , are coefficient functions to be determined by the system of partial differential equations resulted from substituting eq. (3) into eq. (2). In the next section, the Korteweg-de Vries equation [11]:

$$u_t - 6uu_x - u_{xxx} = 0 \quad (4)$$

and the Jimbo-Miwa equation [12]:

$$u_{xxy} - 3u_y u_{xx} - 3u_x u_{xy} - 2u_{yt} - 3u_{xz} = 0 \quad (5)$$

are taken as two examples to illustrate the steps of our algorithm.

### Example 1

For the Korteweg-de Vries eq. (4), we balance the highest order linear term  $u_{xxx}$  with the highest order non-linear term  $6uu_x$  and then have  $n + 3 = n + n + 1$ , i. e.,  $n = 2$ . In this case, eq. (3) becomes:

$$u = u_0(x, t)\varphi^2 + u_1(x, t)\varphi^1 + u_2(x, t), \quad \varphi = 1 + e^\xi, \quad \xi = kx - ct \quad (6)$$

Substituting eq. (6) into eq. (4) and then equating each coefficient of the same order power of  $\varphi^i (i = 0, -1, -2, \dots, -5)$  to zero yields a set of partial differential equations:

$$\varphi^5: 24k^3 e^{3\xi} u_0 - 12k e^\xi u_0^2 = 0 \quad (7)$$

$$\varphi^4: 18k^3 e^{2\xi} u_0 - 6k^3 e^{3\xi} u_1 - 18k e^\xi u_0 u_1 - 18k^2 e^{2\xi} u_{0x} - 6u_0 u_{0x} = 0 \quad (8)$$

$$\varphi^3: 2c e^\xi u_0 - 2k^3 e^\xi u_0 - 6k^3 e^{2\xi} u_1 - 6k e^\xi u_1^2 - 12k e^\xi u_0 u_2 - 6k^2 e^\xi u_{0x} - 6u_0 u_{1x} - 6k e^\xi u_{0xx} = 0 \quad (9)$$

$$\varphi^2: c e^\xi u_1 - k^3 e^\xi u_1 - 6k e^\xi u_1 u_2 - u_{0x} - 6u_2 u_{0x} - 3k^2 e^\xi u_{1x} - 6u_1 u_{1x} - 6u_0 u_{2x} - 3k e^\xi u_{1xx} - u_{0xxx} = 0 \quad (10)$$

$$\varphi^1: u_{1t} - 6u_2 u_{1x} - 6u_1 u_{2x} - u_{1xxx} = 0 \quad (11)$$

$$\varphi^0: u_{2t} - 6u_2 u_{2x} - u_{2xxx} = 0 \quad (12)$$

Solving eqs. (7)-(12), we have:

$$u_0 = 2k^2 e^{2\xi}, \quad u_1 = 2k^2 e^\xi, \quad u_2 = \frac{k^3 - c}{6k} \quad (13)$$

and hence obtain an exact solution of the Korteweg-de Vries eq. (4):

$$u = \frac{2k^2 e^{2\xi}}{(1 - e^\xi)^2} - \frac{2k^2 e^\xi}{1 - e^\xi} + \frac{k^3 - c}{6k} - \frac{1}{2} k^2 \operatorname{sech}^2 \frac{1}{2} \xi + \frac{k^3 - c}{6k}, \quad \xi = kx - ct \quad (14)$$

### Example 2

Let us consider the Jimbo-Miwa eq. (5). To begin with, we balance the highest order linear term  $u_{xxy}$  with the highest order non-linear term  $3u_y u_{xx}$  and then have  $n + 4 = 2n + 3$ , i. e.,  $n = 1$ . In this case, eq. (3) becomes:

$$u = u_0(x, y, z, t)\varphi^{-1} + u_1(x, y, z, t), \quad \varphi = 1 + e^\xi, \quad \xi = kx + w(y, z, t) \quad (15)$$

Substituting eq. (15) into eq. (5) and then equating each coefficient of the same order power of  $\varphi^i (i = 0, -1, -2, \dots, -5)$  to zero yields a set of partial differential equations:

$$\varphi^{-5}: 24k^3 e^{4\xi} w_{1y} u_0 - 12k^2 e^{3\xi} w_{1y} u_0 = 0 \quad (16)$$

$$\varphi^{-4}: 36k^3 e^{3\xi} w_y u_0 - 6k^2 e^{2\xi} w_y u_0^2 - 6k^3 e^{3\xi} u_{0y} - 9k^2 e^{2\xi} u_{0y} u_{0y} - 18k^2 e^{3\xi} w_y u_{0y} - 15k e^{2\xi} w_y u_{0x} u_{0x} = 0 \quad (17)$$

$$\varphi^{-3}: 6k e^{2\xi} w_2 u_0 - 14k^3 e^{2\xi} w_3 u_0 - 4e^{2\xi} w_y w_z u_0 - 6k^3 e^{2\xi} u_{0y} - 3k^2 e^\xi u_{0y} u_{0y} - 6k^2 e^{2\xi} u_{0y} u_{1y} - 18k^2 e^{2\xi} w_y u_{0x} - 3k e^\xi w_y u_{0x} u_{0x} - 9k e^\xi u_{0x} u_{0y} - 3e^\xi w_y u_{0x}^2 - 6k e^{2\xi} w_y u_{0x} u_{1x} - 6k^2 e^{2\xi} u_{0xy} - 3k e^\xi u_{0x} u_{0xy} - 6k e^{2\xi} w_y u_{0xx} - 3e^\xi w_y u_{0xx} = 0 \quad (18)$$

$$\varphi^{-2}: 3k e^\xi w_z u_0 - k^3 e^\xi w_y u_0 - 2e^\xi w_y w_t u_0 - 2e^\xi w_{yt} u_0 - 2e^\xi w_y u_{0t} - 3k e^\xi u_{0z} - k^3 e^\xi u_{0y} - 2e^\xi w_t u_{0y} - 3k^2 e^\xi u_{0y} u_{1y} - 3e^\xi w_z u_{0x} - 3k^2 e^\xi w_y u_{0x} - 6k e^\xi u_{1y} u_{0x} - 3k e^\xi w_y u_{0x} u_{1x} - 3k e^\xi u_{0y} u_{1x} - 3e^\xi w_y u_{0x} u_{1x} - 3k^2 e^\xi u_{0xy} - 3u_{0x} u_{0xy} - 3k e^\xi u_{0x} u_{1xy} - 3k e^\xi w_y u_{0xx} - 3u_{0y} u_{0xx} - 3e^\xi w_y u_{0x} u_{1xx} - 3k e^\xi u_{0xxy} - e^\xi w_y u_{0xxx} = 0 \quad (19)$$

$$\varphi^{-1}: 2u_{0yt} - 3u_{0xz} - 3u_{1x} u_{0xy} - 3u_{0x} u_{1xy} - 3u_{1y} u_{0xx} - 3u_{0y} u_{1xx} - u_{0xxx} = 0 \quad (20)$$

$$\varphi^0: 2u_{1yt} - 3u_{1xz} - 3u_{1x} u_{1xy} - 3u_{1x} u_{1xx} - u_{1xxx} = 0 \quad (21)$$

From eqs. (16)-(21), we have:

$$u_0 = 2k e^\xi, \quad u_1 = \frac{1}{3k^2} \{3k[f_3(z) - f_1(y, z)] - [k^3 - 2p]f_{1y}(y, z)\} dy \quad (22)$$

$$w = f_1(y, z) - f_2(z) - pt \quad (23)$$

where  $f_1(y, z)$  and  $f_2(z)$  are arbitrary functions of the indicated variables,  $f_3(z) = df_3(z)/dz$  and  $p$  is an arbitrary constant.

We therefore obtain an exact solution of the Jimbo-Miwa equation (5):

$$u = \frac{2k e^\xi}{1 - e^\xi} - \frac{1}{3k^2} \{3k[f_3(z) - f_1(y, z)] - [k^3 - 2p]f_{1y}(y, z)\} dy \quad (24)$$

where  $\xi = kx + f_1(y, z) + f_2(z) + pt$ .

It is shown from computer running that the general approach (3) has an advantage over the one in eq. (1) to deal with the so-called middle expression expansion problem appeared in the process of using the exp-function method to transform a given non-linear evolution equation into the over determined system of algebraic or differential equations.

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### References

- [1] He, J.-H., Wu, X. H., Exp-Function Method for Non-Linear Wave Equations, *Chaos, Solitons & Fractals*, 30 (2006), 3, pp. 700-708

- [2] He, J.-H., Abdou, M. A., New Periodic Solutions for Non-Linear Evolution Equations Using Exp-Function Method, *Chaos, Solitons & Fractals*, 34 (2006), 5, pp. 1421-1429
- [3] Ebaid, A., Exact Solitary Wave Solutions for Some Non-Linear Evolution Equations Via Exp-Function Method, *Physics Letters A*, 365 (2007), 3, pp. 213-219
- [4] Zhang, S., Application of Exp-Function Method to a KdV Equation with Variable Coefficients, *Physics Letters A*, 365 (2007), 5-6, 448-453
- [5] Zhang, S., Application of Exp-Function Method to High-Dimensional Non-Linear Evolution Equation, *Chaos, Solitons & Fractals*, 38 (2008), 1, pp. 270-276
- [6] Mohyud-Din, S. T., et al., A. Exp-Function Method for Solitary and Periodic Solutions of Fitzhugh-Nagumo Equation, *International Journal of Numerical Methods for Heat and Fluid Flow*, 22 (2012), 3-4, pp. 335-341
- [7] Zhang, S., Zhang, H. Q., An Exp-Function Method for New N-Soliton Solutions with Arbitrary Functions of a (2+1)-Dimensional vcBK System, *Computers and Mathematics with Applications*, 61 (2011), 8, pp. 1923-1930
- [8] Zhang, S., et al., Multi-Wave Solutions for a Non-Isospectral KdV-Type Equation with Variable Coefficients, *Thermal Science*, 16 (2012), 5, pp. 1576-1579
- [9] He, J.-H., Exp-Function Method for Fractional Differential Equations, *Journal of Non-Linear Sciences and Numerical Simulation*, 14 (2013), 6, pp. 363-366
- [10] Cui, L. X., et al., New Exact Solutions of Fractional Hirota-Satsuma Coupled Korteweg-de Vries Equations, *Thermal Science*, 19 (2015), 4, pp. 1173-1176
- [11] Ablowitz, M. J., Clarkson P. A., *Solitons, Non-Linear Evolution Equations and Inverse Scattering*, Cambridge University Press, Cambridge, New York, USA, 1991
- [12] Zhang, S., A Generalization of the G'/G-Expansion Method and its Application to Jimbo-Miwa Equation, *Bulletin of the Malaysian Mathematical Sciences Society*, 36 (2013), 3, pp. 699-708