

Open forum

ON THE SEMI-INVERSE METHOD AND VARIATIONAL PRINCIPLE

by

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Short paper
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In this Open Forum, Liu et al. proved the equivalence between He-Lee 2009 variational principle and that by Tao and Chen (Tao, Z. L., Chen, G. H., Thermal Science, 17(2013), pp. 951-952) for one dimensional heat conduction. We confirm the correction of Liu et al.'s proof, and give a short remark on the history of the semi-inverse method for establishment of a generalized variational principle.

Key words: variational principle, heat conduction, semi-inverse method,
Lagrange multiplier, parameterized variational principle

Introduction

Liu et al. proved that the following variational principle for 1-D heat conduction [1]:

$$\tilde{J}_{He-Lee} = \int_0^{t_0} \int_a^b \left[\alpha \left(\frac{\partial^2 T}{\partial x^2} + k^2 \frac{\partial T}{\partial t} + \lambda k^2 T \right) + \beta \right] dx dt \quad (1)$$

is equivalent to He-Lee 2009 variational principle [2]:

$$J_{He-Lee} = \int_0^{t_0} \int_a^b \left[\frac{\partial^2 T}{\partial x^2} + k^2 \frac{\partial T}{\partial t} + \lambda k^2 T \right] dx dt \quad (2)$$

and Tao-Chen 2013 variational principle [3]:

$$J_{Tao-Chen} = \int_0^{t_0} \int_a^b \left[\frac{\partial T}{\partial t} + \lambda T + \frac{1}{k^2} \frac{\partial^2 T}{\partial x^2} \right] dx dt \quad (3)$$

for all non-zero constants α and β . Liu et al.'s proof is straightforward and it is very easy for understanding. To show this, we consider a simple function:

$$y(x) = x^2 - x \quad (4)$$

Its extreme value is same with the following one:

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$$y(x) = \alpha(x^2 - x) + \beta \quad (5)$$

for all non-zero constants α and β .

Tao and Chen applied the semi-inverse method proposed in 2007 [4], hereby we give a tutorial introduction to the method for beginners.

The semi-inverse method

To elucidate basic property of the semi-inverse method [4], we consider a 2-D incompressible and potential flow, its governing equations are:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (6)$$

$$\frac{\partial \Phi}{\partial x} = u, \quad \frac{\partial \Phi}{\partial y} = v \quad (7)$$

where u and v are velocities in x - and y -directions, respectively, Φ is the potential. There is a known variational principle for the problem, which is:

$$J(\Phi) = \iint \frac{1}{2}(u^2 + v^2) dx dy \quad (8)$$

which is subject to the constraints, eq. (7).

The general approach to establishment of a generalized variational principle is the Lagrange multiplier method [5]:

$$J(\Phi, u, v, \lambda_1, \lambda_2) = \iint \left[\frac{1}{2}(u^2 + v^2) + \lambda_1 \left(\frac{\partial \Phi}{\partial x} - u \right) + \lambda_2 \left(\frac{\partial \Phi}{\partial y} - v \right) \right] dx dy \quad (9)$$

where λ_1 and λ_2 are Lagrange multipliers.

Considering the fact that the Lagrange multipliers involved in eq. (9) are unknown, the semi-inverse method [4] is to replace the terms involving the Lagrange multipliers by an unknown function, F , in the form:

$$J(\Phi, u, v) = \iint \left[\frac{1}{2}(u^2 + v^2) + F \right] dx dy \quad (10)$$

where F is an unknown function of the variables u , v , Φ and/or their derivatives $F = F(u, v, \Phi, \Phi_x, \Phi_y, \dots)$

We call eq. (10) a trial-functional. Making eq. (10) stationary with respect to u , v , and Φ :

$$\delta J(\Phi, u, v) = \iint \left[(u \delta u + v \delta v) + \frac{\delta F}{\delta u} \delta u + \frac{\delta F}{\delta v} \delta v + \frac{\delta F}{\delta \Phi} \delta \Phi \right] dx dy = \quad (11)$$

$$= \iint \left[\left(u + \frac{\delta F}{\delta u} \right) \delta u + \left(v + \frac{\delta F}{\delta v} \right) \delta v + \frac{\delta F}{\delta \Phi} \delta \Phi \right] dx dy = 0 \quad (11)$$

we have the following Euler-Lagrange equations:

$$u + \frac{\delta F}{\delta u} = 0 \quad (12)$$

$$v + \frac{\delta F}{\delta v} = 0 \quad (13)$$

$$\frac{\delta F}{\delta \Phi} = 0 \quad (14)$$

where $\delta F/\delta \Phi$ is the variational derivative with respect to Φ , defined as:

$$\frac{\delta F}{\delta \Phi} = \frac{\partial F}{\partial \Phi} - \frac{\partial}{\partial x} \left(\frac{\partial F}{\partial \Phi_x} \right) - \frac{\partial}{\partial y} \left(\frac{\partial F}{\partial \Phi_y} \right) + \frac{\partial^2}{\partial x^2} \left(\frac{\partial F}{\partial \Phi_{xx}} \right) + \dots \quad (15)$$

The above equations should be equivalent to eqs. (1) and (2). To this end, we set:

$$u + \frac{\delta F}{\delta u} = a \left(u - \frac{\partial \Phi}{\partial x} \right) \quad (16)$$

$$v + \frac{\delta F}{\delta v} = b \left(v - \frac{\partial \Phi}{\partial y} \right) \quad (17)$$

$$\frac{\delta F}{\delta \Phi} = c \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \quad (18)$$

where a , b , and c are non-zero constants. From eqs. (16) and (17), we have:

$$F = \frac{1}{2}(a-1)u^2 + \frac{1}{2}(b-1)v^2 - au \frac{\partial \Phi}{\partial x} - bv \frac{\partial \Phi}{\partial y} + F_1 \quad (19)$$

where F_1 is an unknown function of Φ and/or its derivatives. Equations (18) and (19) imply that:

$$a \frac{\partial u}{\partial x} + b \frac{\partial v}{\partial y} + \frac{\delta F_1}{\delta \Phi} = c \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \quad (20)$$

from which we can obtain the following relations:

$$F_1 = 0 \text{ and } a = b = k \quad (21)$$

Finally we obtain a parameterized variational principle, which reads:

$$\begin{aligned} J(\Phi, u, v) &= \iint \left[\frac{1}{2}(u^2 + v^2) + \frac{1}{2}(k-1)u^2 + \frac{1}{2}(k-1)v^2 - ku \frac{\partial \Phi}{\partial x} - kv \frac{\partial \Phi}{\partial y} \right] dx dy = \\ &= - \iint k \left[\frac{1}{2}(u^2 + v^2) + u \frac{\partial \Phi}{\partial x} - v \frac{\partial \Phi}{\partial y} \right] dx dy \end{aligned} \quad (21)$$

where k is a non-zero constant.

Conclusions

Equation (1) is a parameterized variational principle, and we further confirm hereby that the variational principle by Tao and Chen is not a new one, but it is equivalent to He Lee 2009 variational principle. The development of the semi-inverse method was summarized in [6-9].

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