

## NUMERICAL SOLUTION OF A COUPLED REACTION-DIFFUSION MODEL USING BARYCENTRIC INTERPOLATION COLLOCATION METHOD

by

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*In thermal science, chemical and mechanics, the non-linear reaction-diffusion model is very important, and an approximate solution with high precision is always needed. In this article, the barycentric interpolation collocation method is proposed for this purpose. Numerical experiments show that the proposed approach is highly reliable.*

**Key words:** *non-linear reaction-diffusion model, numerical experiments, barycentric interpolation collocation method,*

### Introduction

The paper is devoted to the numerical solution of a class of non-linear reaction-diffusion models:

$$\begin{cases} \frac{\partial u}{\partial t} = d_1 \frac{\partial^2 u}{\partial x^2} + h_1(u, v) \\ \frac{\partial v}{\partial t} = d_2 \frac{\partial^2 v}{\partial x^2} + h_2(u, v) \end{cases} \quad a \leq x \leq b, \quad t \geq 0 \quad (1)$$

with the following initial conditions and boundary conditions:

$$\begin{aligned} u(x, 0) &= f_0(x), v(x, 0) = g_0(x), \quad a \leq x \leq b \\ u(a, t) &= f_1(t), u(b, t) = f_2(t), \quad t \geq 0 \\ v(a, t) &= g_1(t), v(b, t) = g_2(t), \quad t \geq 0 \end{aligned} \quad (2)$$

where  $d_1 > 0$  and  $d_2 > 0$  are diffusion coefficients and may also include the parametrizations of turbulence.

The reaction-diffusion model has wide applications in thermal science, chemical and mechanics [1-5]. There are some analytical and numerical methods [6-10] for solving the advection-reaction-diffusion system, among which the barycentric interpolation collocation method [11-17] is a high precision method. In this paper, we mainly employ the Lagrange barycentric interpolation collocation method to solve the reaction-diffusion model (1).

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The barycentric interpolation formula of the eq. (1)

We give two initial hypothesis functions  $u_0, v_0$  and construct the following linear iterative format:

$$\begin{cases} \frac{\partial u_n}{\partial t} = d_1 \frac{\partial^2 u_n}{\partial x^2} + h_1(u_{n-1}, v_{n-1}) \\ \frac{\partial v_n}{\partial t} = d_2 \frac{\partial^2 v_n}{\partial x^2} + h_2(u_{n-1}, v_{n-1}) \end{cases} \quad n = 1, 2, 3, \dots \quad (3)$$

Next, we give a meshless barycentric interpolation collocation method of eq. (3).

Let  $a \leq x_1 < x_2 < \dots < x_M \leq b$ ,  $0 \leq t_1 < t_2 < \dots < t_N \leq T$ . The barycentric interpolation form of function  $u(x, t)$  can be expressed:

$$u(x, t) = \sum_{i=1}^M \sum_{j=1}^N \xi_i(x) \eta_j(t) u(x_i, t_j) \quad (4)$$

where  $\xi_i(x)$  is the barycentric interpolation basis function of  $M$  nodes on interval  $[a, b]$ ,  $\eta_j(t)$  is the barycentric interpolation basis function of  $N$  nodes on interval  $[0, T]$ .

$$\xi_i(x) = \frac{\prod_{k=1, k \neq i}^M (x - x_k)}{\prod_{k=1, k \neq i}^M (x_i - x_k)}, \quad i = 1, 2, \dots, M, \quad \eta_j(t) = \frac{\prod_{k=1, k \neq j}^N (t - t_k)}{\prod_{k=1, k \neq j}^N (t_j - t_k)}, \quad j = 1, 2, \dots, N$$

Using the method of [11-12, 15], eq. (3) can be written in following partitioned matrix form:

$$\begin{bmatrix} D^{(0,1)} - d_1 D^{(2,0)} & 0 \\ 0 & D^{(0,1)} - d_2 D^{(2,0)} \end{bmatrix} \begin{bmatrix} u_n \\ v_n \end{bmatrix} = \begin{bmatrix} \text{diag}[h_1(u_{n-1}, v_{n-1})] \\ \text{diag}[h_2(u_{n-1}, v_{n-1})] \end{bmatrix} \quad (5)$$

where  $D^{(0,1)}$ ,  $D^{(0,2)}$ , and  $\text{diag}[h_1(u_{n-1}, v_{n-1})]$  are explained in [11-12, 15]. In this paper, we use displacement method to impose the initial boundary conditions. The detailed procedure can be found in [15].

## Numerical experiments

*Experiment 1.* We consider the following reaction-diffusion model [3]:

$$\begin{cases} \frac{\partial u}{\partial t} = a_1 \frac{\partial^2 u}{\partial x^2} - uv^2 + a_2(1 - u) \\ \frac{\partial v}{\partial t} = b_1 \frac{\partial^2 v}{\partial x^2} + uv^2 - (a_2 + b_2)v \end{cases} \quad (6)$$

with the following initial conditions:

$$u(x, 0) = 1 - 0.5 \sin^{100} \frac{\pi(x-50)}{100}, \quad v(x, 0) = 0.25 \sin^{100} \frac{\pi(x-50)}{100}, \quad x \in [-50, 50] \quad (7)$$

and the boundary conditions:

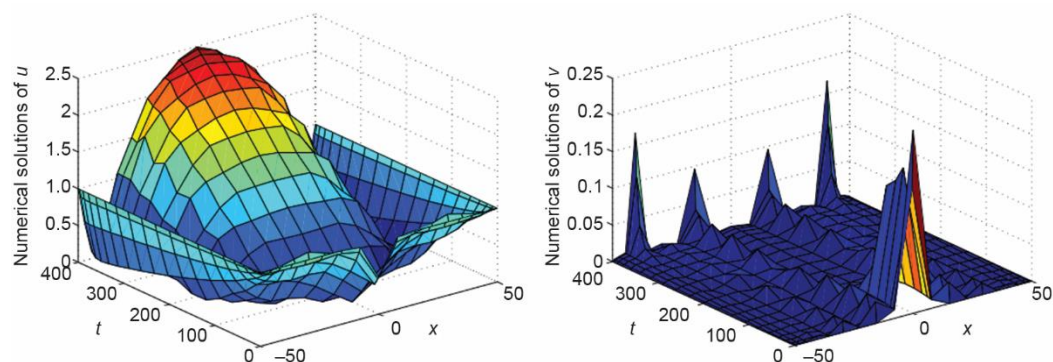
$$u(-50,t) = u(50,t) = 1, \quad v(-50,t) = v(50,t) = 0, \quad t \in [0,400] \quad (8)$$

In tab. 1, we analyze the influence of  $b_1, b_2$  on  $u$  and  $v$  at  $t = 100, M = N = 20$ , with the given parameters  $a_1 = 1.0, a_2 = 0.064$ .

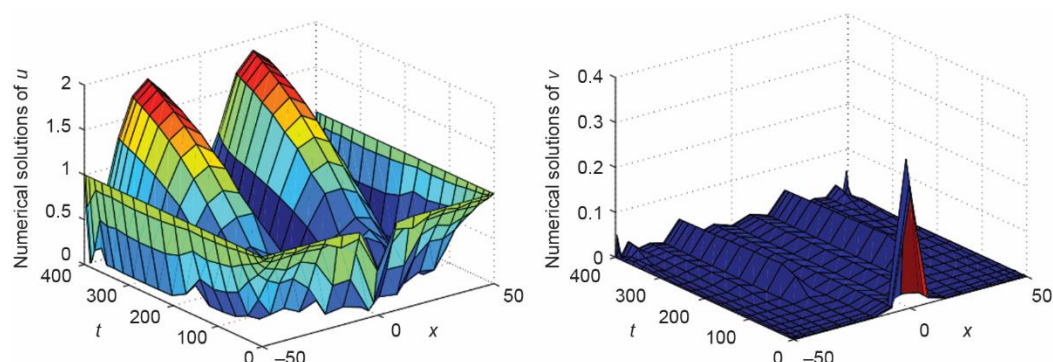
Figures 1-3 show the numerical solutions of  $u$  and  $v$  obtained by the present method, where  $b_1 = 0.01, b_2 = 0.062, M = 20, N = 20$ .

**Table 1. Comparison of numerical solutions at  $t = 100$  for Experiment 1 ( $M = N = 20$ )**

$x$	$u(x, t)$	$v(x, t)$	$u(x, t)$	$v(x, t)$	$u(x, t)$	$v(x, t)$
	$b_1 = 0.01$	$b_2 = 0.062$	$b_1 = 0.5$	$b_2 = 0.062$	$b_1 = 0.01$	$b_2 = -0.3$
-50	1.0000	0.0000	1.0000	0.0000	1.0000	0.0000
-25	0.5493	-0.0012	0.5156	-0.0090	0.5153	-0.0024
0	0.0006	0.1575	0.2454	0.3662	0.2805	0.2998
25	0.5493	-0.0012	0.5162	-0.0095	0.5153	-0.0024
50	1.0000	0.0000	1.0000	0.0000	1.0000	0.0000



**Figure 1. The numerical solution for Experiment 1 at  $a_1 = 1.0, a_2 = 0.064$**



**Figure 2. The numerical solution for Experiment 1 at  $a_1 = 1.0, a_2 = 0.5$**

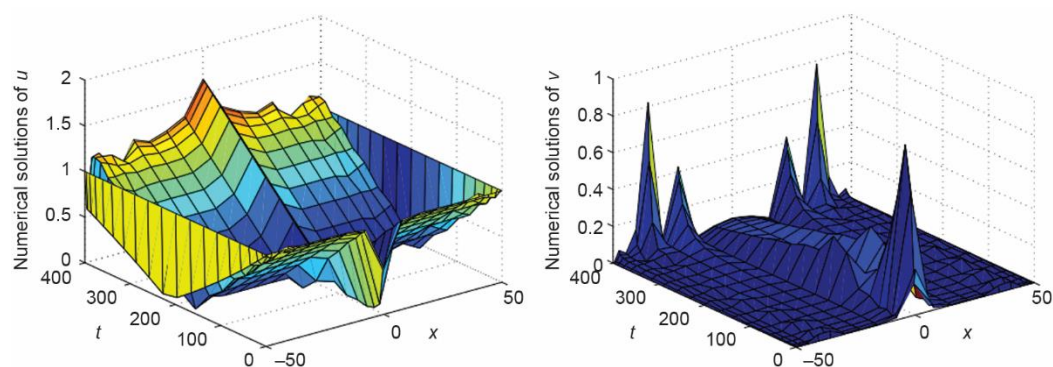


Figure 3. The numerical solution for *Experiment 1* at  $a_1 = 0.001$ ,  $a_2 = 0.064$

*Experiment 2.* Consider the following Brusselator model [4, 5]:

$$\begin{cases} \frac{\partial u}{\partial t} = 10^{-4} \frac{\partial^2 u}{\partial x^2} - 4.4u + u^2v + 1 \\ \frac{\partial v}{\partial t} = 10^{-4} \frac{\partial^2 v}{\partial x^2} + 3.4u - u^2v \end{cases} \quad (9)$$

with the following initial conditions:

$$u(x, 0) = 0.5, v(x, 0) = 1 + 5x, \quad x \in [0, 1] \quad (10)$$

and the boundary conditions:

$$u_x(0, t) = u_x(1, t) = 0, \quad v_x(0, t) = v_x(1, t) = 0, \quad t \in [0, 1] \quad (11)$$

The results of *Experiment 2* are given in fig. 4.

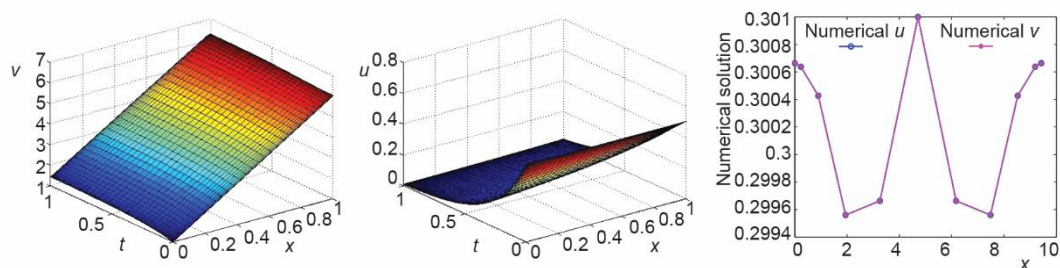


Figure 4. Numerical solutions obtained by the present method for *Experiment 2*

*Experiment 3.* Consider the Schnakenberg model [1]:

$$\begin{cases} \frac{\partial u}{\partial t} = d_1 \frac{\partial^2 u}{\partial x^2} + a - u + u^2v \\ \frac{\partial v}{\partial t} = d_2 \frac{\partial^2 v}{\partial x^2} + b - u^2v \end{cases} \quad (12)$$

- (1) We select the parameter  $d_1 = 0.2, d_2 = 0.1, a = 0.14, b = 0.66$ , the initial condition  $u(x, 0) = 0.8 + 0.1 \cos(x), v(x, 0) = 1.03 + 0.1 \cos(x)$  and the boundary condition  $u_x(0, t) = 0, u_x(3\pi, t) = 0, v_x(0, t) = 0, v_x(3\pi, t) = 0$ .
- (2) We select the parameter  $d_1 = 0.01, d_2 = 1.0, a = 0.14, b = 0.16$  and the initial condition  $u(x, 0) = 0.3 + 0.001 \sin(3x), v(x, 0) = 1.778 + 0.001 \cos(2x)$ . The results of *Experiment 3* are given in figs. 5 and 6.

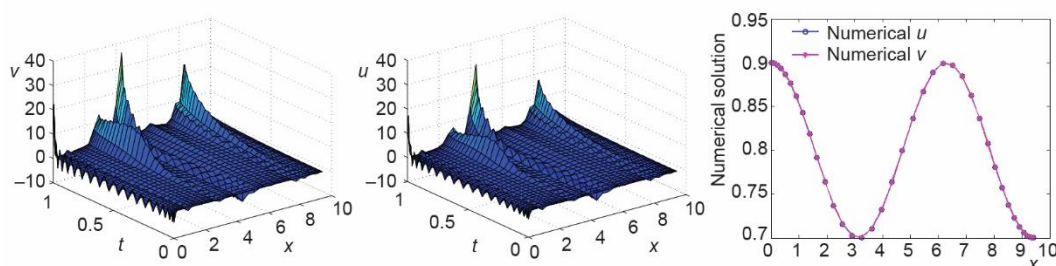


Figure 5. Numerical solutions obtained by the present method for *Experiment 3* (1)

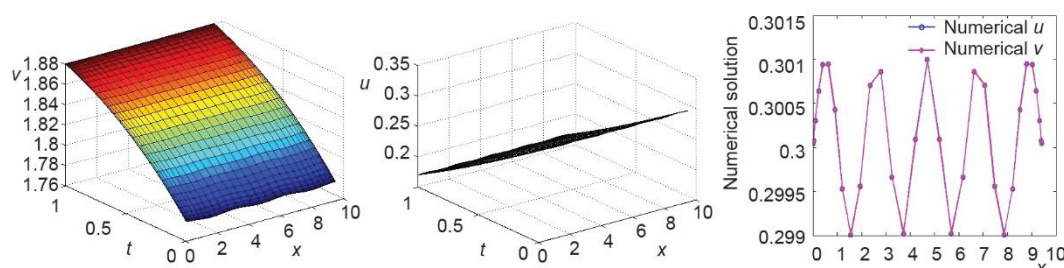


Figure 6. Numerical solutions obtained by the present method for *Experiment 3* (2)

*Experiment 4.* Consider the isothermal chemical model [2]:

$$\begin{cases} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} - uv \\ \frac{\partial v}{\partial t} = \frac{\partial^2 v}{\partial x^2} - kv + uv \end{cases} \quad (13)$$

$$\begin{aligned} u(x, 0) &= 1, & v(x, 0) &= \exp(-x^2), & 0 \leq x &\leq 200 \\ u_x(0, t) &= 0, & u(200, t) &= 1, & v_x(0, t) &= 0, & u(200, t) &= 0 \end{aligned}$$

We select the parameter  $k = 0.1, 0.3, 0.5, 0.7, 0.9$ . The results of *Experiment 4* are given in fig. 7 and tab. 2.

Table 2. Comparison of numerical solution with the  $k$  for *Experiment 4*

$k$	Numerical solution of $u$	Numerical solution of $v$	$k$	Numerical solution of $u$	Numerical solution of $v$
0.1	0.0038529	0.0053738	0.7	0.0046887	0.0049609
0.3	0.0037315	0.0050254	0.9	0.0050133	0.0051748
0.5	0.0042162	0.0043605			

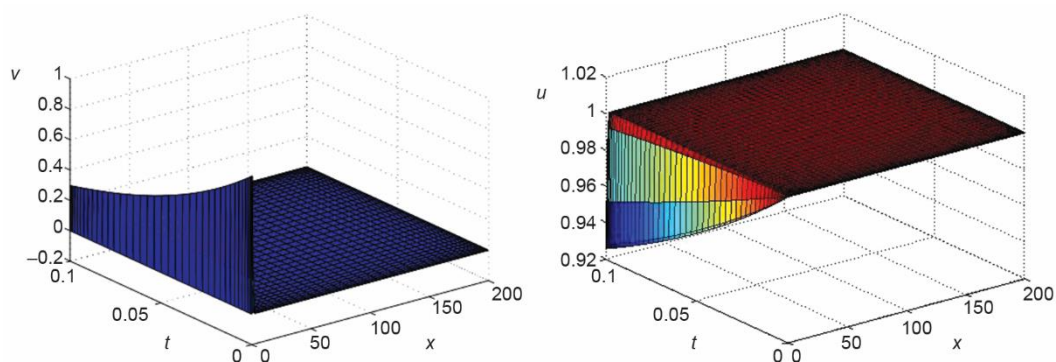


Figure 7. Numerical solutions obtained by the present method for *Experiment 4*

## Conclusion

In this paper, a class of reaction-diffusion systems have been solved by using the barycentric interpolation collocation method. The numerical experiments show that the algorithm is highly accurate. All computations are performed by a mathematical software.

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