

## ON LOCAL FRACTIONAL VOLTERRA INTEGRAL EQUATIONS IN FRACTAL HEAT TRANSFER

by

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*In the article, the fractal heat-transfer models are described by the local fractional integral equations. The local fractional linear and non-linear Volterra integral equations are employed to present the heat transfer problems in fractal media. The local fractional integral equations are derived from the Fourier law in fractal media.*

*Key words: fractal heat transfer, Fourier law, local fractional calculus, integral equations*

### Introduction

Integral equations have played important roles in engineering applications [1]. One of important applications is to deal with the heat flux in inverse heat conduction [2-6], and so on.

Local fractional calculus [7-11], has used to describe the non-differentiable problems in heat transfer. For example, the heat-conduction equation in fractal vector space was proposed in [12-14]. Many numerical and analytical technologies for fractal heat-transfer problems were also developed, such as Sumudu transform series expansion method [15], Laplace transform series expansion method [16], variational iteration algorithm [17, 18], Laplace transform variational iteration method [19], decomposition method [20, 21], Laplace transform decomposition method [22], and homotopy perturbation method [23, 24] via local fractional operator.

The main aim of this article is to propose the fractal heat-transfer models described by the local fractional integral equations.

### Fourier law of heat conduction in fractal media

The fractal temperature field reads as [12, 13, 18]:

$$\Lambda(x, y, z, \tau) = f(x, y, z, \tau) \text{ at } \tau > \tau_0 \text{ and in } \Omega_\alpha \quad (1)$$

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where  $f(x, y, z, \tau)$  is a non-differentiable function in fractal domain  $\Omega_\alpha$ .

For a given fractal temperature field  $\Lambda(x, y, z, \tau)$ , a local fractional temperature gradient of  $\Lambda(x, y, z, \tau)$  can be written [12-14, 18]:

$$\nabla^\alpha \Lambda(x, y, z, \tau) = \frac{\partial^\alpha \Lambda(x, y, z, \tau)}{\partial x^\alpha} e_1 + \frac{\partial^\alpha \Lambda(x, y, z, \tau)}{\partial y^\alpha} e_2 + \frac{\partial^\alpha \Lambda(x, y, z, \tau)}{\partial z^\alpha} e_3 \quad (2)$$

where  $\nabla^\alpha$  is the local fractional derivative operator in fractal time-space [12-14].

We consider the fractal heat flux per unit fractal area, denoted by  $\bar{q}(x, y, z, t)$ , is proportional to the fractal temperature gradient in fractal media. Fourier law of heat conduction in fractal media can be written [12-14]:

$$\bar{q}(x, y, z, t) = -\mu \nabla^\alpha \Lambda(x, y, z, t) \quad (3)$$

where  $\mu$  is the thermal conductivity of the fractal material, which is related to the fractal dimensions of materials. Meanwhile, we observe that  $\bar{q}(x, y, z, t)$  is the fractal Fourier flow and that  $\Lambda(x, y, z, t)$  is the non-differentiable temperature field.

Fourier law of low-dimensional heat conduction equation in fractal media is given by:

$$q(x, t) = -\mu \frac{\partial^\alpha \Lambda(x, t)}{\partial x^\alpha}, \quad \text{at } \tau > \tau_0 \quad \text{and in } \Gamma_\alpha \quad (4)$$

where  $\mu$  is the thermal conductivity of the fractal material and  $\partial^\alpha / \partial x^\alpha$  is the local fractional partial derivative operator with respect to the space  $x$  [7, 12, 13].

When  $t = t_0$ , eq. (4) can be written:

$$q(x) = -\mu \frac{d^\alpha \Lambda(x)}{dx^\alpha} \quad (5)$$

where  $\mu$  is the thermal conductivity of the fractal material and  $d^\alpha / dx^\alpha$  is the local fractional derivative operator [7-25]. Equation (5) can be applied to describe the fractal steady heat conduction problems.

### The Volterra integral equation in fractal heat transfer

In order to obtain the Volterra integral equations in fractal heat transfer, eq. (5) can be written:

$$\frac{d^\alpha \Lambda(x)}{dx^\alpha} = q[x, \Lambda(x)], \quad (0 \leq x \leq l) \quad (6)$$

where  $q(x) = -\mu q[x, \Lambda(x)]$ .

With the help of the result [25], we have the local fractional Volterra integral equation of second kind:

$$\Lambda(x) = \Lambda(0) + \frac{1}{\Gamma(1+\alpha)} \int_0^x q[\tau, \Lambda(\tau)] (d\tau)^\alpha \quad (7)$$

where  $\int_0^x (d\tau)^\alpha / \Gamma(1 + \alpha)$  is a local fractional integral operator [7, 12, 13, 25].

When the fractal flow is decomposed into the linear term of non-differentiable type:

$$q[x, \Lambda(x)] = \beta_0(x) + \beta_1(x)\Lambda(x) \quad (8)$$

Equation (7) is rewritten in the form:

$$\Lambda(x) = \Lambda(0) + \frac{1}{\Gamma(1 + \alpha)} \int_0^x [\beta_0(\tau) + \beta_1(\tau)\Lambda(\tau)](d\tau)^\alpha \quad (9)$$

where  $\beta_0(x)$  and  $\beta_1(x)$  are two parameters related to the fractal heat flow.

Taking:

$$\phi(x) = \Lambda(0) + \frac{1}{\Gamma(1 + \alpha)} \int_0^x \beta_0(\tau)(d\tau)^\alpha \quad (10)$$

we can rewrite eq. (9):

$$\Lambda(x) = \phi(x) + \frac{1}{\Gamma(1 + \alpha)} \int_0^x \beta_1(\tau)\Lambda(\tau)(d\tau)^\alpha \quad (11)$$

When  $\Lambda(0) = 0$  and  $\beta_0(x) = 0$ , we obtain the local fractional linear Volterra integral equation of first kind:

$$\Lambda(x) = \frac{1}{\Gamma(1 + \alpha)} \int_0^x \beta_1(\tau)\Lambda(\tau)(d\tau)^\alpha \quad (12)$$

Taking:

$$\beta_1(\tau) = \frac{(x - \tau)^\alpha}{\Gamma(1 + \alpha)} \quad (13)$$

from eqs. (11) and (12) we get the local fractional linear Volterra integral equations of convolution-type:

$$\Lambda(x) = \phi(x) + \frac{1}{\Gamma(1 + \alpha)} \int_0^x \frac{(x - \tau)^\alpha}{\Gamma(1 + \alpha)} \Lambda(\tau)(d\tau)^\alpha \quad (14)$$

and

$$\Lambda(x) = \frac{1}{\Gamma(1 + \alpha)} \int_0^x \frac{(x - \tau)^\alpha}{\Gamma(1 + \alpha)} \Lambda(\tau)(d\tau)^\alpha \quad (15)$$

When:

$$\beta_1(\tau) = 1 \quad (16)$$

the local fractional linear Volterra integral equations of convolution-type are:

$$\Lambda(x) = \phi(x) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \Lambda(\tau) (d\tau)^\alpha \quad (17)$$

and

$$\Lambda(x) = \frac{1}{\Gamma(1+\alpha)} \int_0^x \Lambda(\tau) (d\tau)^\alpha \quad (18)$$

When the fractal flow is decomposed into the non-linear term of non-differentiable type:

$$q[x, \Lambda(x)] = \chi(x) \Lambda^2(x) \quad (19)$$

we obtain the local fractional non-linear Volterra integral equation of second kind:

$$\Lambda(x) = \Lambda(0) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \chi_3(x) \Lambda^2(x) (d\tau)^\alpha \quad (20)$$

where  $\chi(x)$  is the parameter related to the fractal heat flow.

In view of eq. (19), the local fractional non-linear Volterra integral equation of first kind is written as:

$$\Lambda(x) = \frac{1}{\Gamma(1+\alpha)} \int_0^x \chi(x) \Lambda^2(x) (d\tau)^\alpha \quad (21)$$

Taking:

$$\chi(x) = \frac{(x-\tau)^\alpha}{\Gamma(1+\alpha)} \quad (22)$$

the local fractional non-linear Volterra integral equations of convolution-type are:

$$\Lambda(x) = \Lambda(0) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-\tau)^\alpha}{\Gamma(1+\alpha)} \Lambda^2(x) (d\tau)^\alpha \quad (23)$$

and

$$\Lambda(x) = \frac{1}{\Gamma(1+\alpha)} \int_0^x \frac{(x-\tau)^\alpha}{\Gamma(1+\alpha)} \Lambda^2(x) (d\tau)^\alpha \quad (24)$$

When

$$\chi(x) = 1 \quad (25)$$

we obtain the local fractional non-linear Volterra integral equations:

$$\Lambda(x) = \Lambda(0) + \frac{1}{\Gamma(1+\alpha)} \int_0^x \Lambda^2(x) (d\tau)^\alpha \quad (26)$$

and



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