NUMERICAL SIMULATION OF THE FORMATION AND THE DRIPPING OF DROPLET IN THE ELECTROSLAG REMELTING PROCESS

by

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Based on the magnetohydrodynamic module of the commercial computational fluid dynamics software FLUENT, a 3-D mathematical model was developed to describe the multi-field coupling phenomenon in the electroslag remelting process. The model predicted value is in good agreement with the experimental measurements. The simulation results show that, during the melt dripping process, the resistance of the electroslag remelting system is decreasing while the current density and velocity, as well as the temperature are increasing gradually. The maximum value of current density, electromagnetic force, and velocity appears when the melt droplet detaches and the quality of large droplets fall into the slag-metal interface. The maximum value of the current density is increased by almost one order of magnitude. The maximum value of the electromagnetic force and velocity increased by 2.5 and 4.7 times, respectively. The maximum value of Joule heating and temperature appeared when the droplets fall into the slag-metal interface. The maximum value of the Joule heating and temperature increased for about 174.7% and 26.8%, respectively, when compared with the moment of the melt droplet formation.

Key words: electroslag remelting, magnetohydrodynamics, multi-field coupling, formation and dripping of the melt droplet

Introduction

Electroslag remelting (ESR) is an advanced metallurgical technology for production of high quality special steel and alloy steel, which couples many phenomena such as, MHD and heat and mass transfer. The fine grained, high-purity ingot products are widely used in the production of turbine disks for aeronautical bearings, submarine shells, aircraft engines, gun steels, titanium alloys for astronautics, crank axles for ships, and high-speed tool steels.

In the ESR process, the alternating or direct current is passed from a conventionally cast electrode through the molten slag and the liquid metal pool to the ingot, as shown in fig. 1. Due to the high electrical resistivity of the slag, Joule heating is generated in the slag layer, which melts the electrode tip. The molten metal passes through the slag layer in the form of droplets, accumulates in the liquid metal pool and then solidifies directionally. The metallurgical properties of the ingot are directly related to the solidification process, which strongly depends on hydrodynamic and thermal condition in the liquid pool and the mushy zone, as well as the

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development of the heat and mass transfer at the slag/droplet interface. On the other hand, fluid flow, heat and mass transfer in the slag phase influence the electrode melting process and, thus, the energy efficiency of the operation. Therefore, a thorough understanding of the transport phenomena in the ESR system is needed to improve the control of the ESR process.

Due to the opaque nature of the slag and the metal phase, as well as the high temperature, experimental observation of the process is very difficult. More and more researchers have devoted significant efforts to develop the mathematical models for better understanding of the electromagnetic phenomenon, fluid flow, heat and mass transfer during the ESR process. Mitchell and Joshi [1] developed the mathematical models to predict the voltage distribution within the slag. Lately, Li et al. [2] developed a 3-D finite element model for the ESR system and obtained good agreement between the computed results and the measurement. Based on the finite difference technique, the general model have been developed for analyzing the electromagnetic fields, fluid flow and heat transfer in ESR [3]. The effects of electromagnetic force, natural convection, momentum, and heat transfer by the moving droplets were thoroughly investigated. However, the interface between the slag-pool and pool-ingot were assumed to be flat. In addition, the influence of the droplets was taken into account in the form of a momentum and energy source applied to slag and pool regions [4-10]. These simplifying assumption restrict the further application of the aforementioned models. Kharicha et al. [11] have showed that the slag/steel interface is never flat and may even become unstable. Recently, models fully coupled the MHD and the multiphase flow fields have been used successfully to solve 2-D problems related to the ESR process [12].

Thus, it is fundamentally important to study how the droplets forms and behaves in the ESR system and their effects on electromagnetic fields, fluid flow, as well as heat and mass transfer. The purpose of the current research is to report on numerical investigation, which have been undertaken to study the unsteady distribution of current, electromagnetic force, velocity, and temperature filed based on the 3-D MHD model coupled with a volume of fluid (VOF) model for the phases distribution.

The numerical model

In the development of a mathematical statement of this problem, the following conservation equations of fluid flow and electromagnetic force field, as well as the heat transfer have to be considered. The coupled nature of these equations is readily apparent, in which the temperature field affects the fluid flow equations through the buoyancy term, and the knowledge of the velocity field and of the turbulence parameters is required in the solution of the convective heat transfer equation. The geometrical parameters, Physical properties of the metal and slag, as well as the process parameters are list in tab. 1-3.

**Governing equations**

This model is based on the numerical solution of the Maxwell, Navier-Stokes, and convective heat balance equations. It is assumed that the eff-

<table>
<thead>
<tr>
<th>Geometry</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Electrode diameter [m]</td>
<td>0.057</td>
</tr>
<tr>
<td>Immersion depth of electrode [m]</td>
<td>0.02</td>
</tr>
<tr>
<td>Slag height [m]</td>
<td>0.07</td>
</tr>
<tr>
<td>Mold diameter [m]</td>
<td>0.10</td>
</tr>
</tbody>
</table>
The effects of the electrode movement are neglected, which is valid for the slag and the upper part of the ingot once the ingot height reaches the order of magnitude of the ingot diameter [13].

The equation of continuity for incompressible fluid is:
\[ \nabla \vec{u} = 0 \]  
(1)

and the equation of motion is given:
\[ \frac{\partial}{\partial t}(\rho \vec{u}) + \nabla \left( \rho \vec{u} \vec{u} \right) = -\nabla p + \mu_{eff} \nabla^2 \vec{u} + \vec{F}_b + \vec{F}_e \]  
(2)

where \( \vec{u} \) is the velocity of the flow, \( \rho \) – the density of the fluid, and \( p \) – the pressure. The turbulent viscosity, \( \mu_{eff} \), is obtained with the aid of the RNG \( k-\varepsilon \) model. \( \vec{F}_b \) represents the effects of buoyancy force determined by the Boussinesq approximation and \( \vec{F}_e \) is the Lorenz force.

The energy equation of the enthalpy-porosity formulation is:
\[ \frac{\partial}{\partial t}(\rho H) + \nabla \rho \vec{u} H = \nabla \left( \mu_{eff} \nabla T \right) + Q_{source} \]  
(3)

where \( H \) is the enthalpy of the metal and the slag which is the sum of the sensible enthalpy and the latent heat released in the mushy zone:
\[ H = h + f_l L \]  
(4)

\[ h = h_{lat} + \int_{T_{lat}}^T c_p dT \]  
(5)

where \( f_l \) is the liquid fraction and defined:
\[ f_l = \frac{T - T_s}{T_l - T_s} \]  
(6)

where \( T_s \) and \( T_l \) are solidus and liquidus temperature, respectively.

The mathematical formulation of the electric and magnetic field is composed of the Maxwell’s equations, Lorenz law, and Joule law.
\[ \nabla \times \vec{H} = \vec{J} \]  
(7)

\[ \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \]  
(8)
where \( \mathbf{H} \) is the magnetic field, \( \mathbf{J} \) – the current density, \( \mathbf{E} \) – the electric field, \( \mathbf{B} \) – the magnetic flux density, \( \mu_0 \) – the magnetic permeability, and \( \sigma \) – the electric conductivity.

The metal/slag interface is tracked with the VOF method which follows conservation equation:

\[
\frac{\partial \alpha}{\partial t} + \nabla \cdot (\alpha \mathbf{u}) = 0
\]

The mushy zone is treated as a pseudo porous medium through the enthalpy-porosity formulation in which the porosity decreases from 1 to 0 as the metal solidifies. The liquid fraction in each cell is set to the porosity in that cell.

**Boundary conditions**

At the inlet the melting rate obtained from experiment is imposed and outflow is applied on the outlet. The no-slip condition is applied on the wall where the slag and metal are in contact with the mold. Along the exposed slag surface, a zero shear condition is assumed for the radial velocity. For the electromagnetic field, zero potential is applied on inlet and a corresponding potential gradient is applied on outlet. The mold lateral wall is assumed to be an insulating wall. The top surface of slag as well as the mold wall is supposed to exchange heat with the surrounding air by natural convection and thermal radiation.

**Solution method**

The electromagnetic field is obtained by the electrical potential method \([14]\) which simultaneously solve the electric potential, \( \varphi \), as well as magnetic potential vector, \( \mathbf{A} \), equations. A control volume-based technique was used to convert the non-linear governing equations to algebraic equations that can be solved numerically. The SIMPLE algorithm was applied for the pressure-velocity coupling. A segregated and implicit solver was used to solve the governing equations and the second-order accuracy central differencing scheme was adopted to improve the accuracy of the solution. The commercial finite volume CFD code, FLUENT was used to compute the two-phase flow, heat transfer, and solidification and the MHD model coupled user defined functions was employed to describe the electromagnetic phenomena. The electric current distribution is dynamically adapted to the transient phase distribution and then the electromagnetic forces and the Joule heating are recalculated at each time step, which are incorporated into the momentum and energy equation as a source term, respectively. The mesh is refined to the point where further refinements will not induce noticeable change (not more than 5\%) in results. The convergence criterion for scaled residuals was set to be less than \( 10^{-6} \). The computational model and mesh are shown in fig. 2.
Results and discussion

Figure 3 shows a comparison between the experimentally measured axial temperature profiles at radial position \( r = 2.5 \text{ cm} \) and those predicted from the present numerical model. The discrete data points denote measurements and the line denotes predictions for the condition. The computed values are in good agreement with measurement and the deviation between the experiment and computation is less than 5%, which proves the accuracy of the present mathematical model.

In figs. 4(a)-(d), the representative snapshots of the droplet dripping during the ESR process are shown at four different instants. There, the first figure corresponds to an instant immediately following the initial formation of the metal droplet due to the temperature rising and melting of the electrode tip. In fig. 4(b), the metal droplet is enlarged and necking phenomenon occurs. Due to the joint effects of the gravity, Lorentz force and surface tension, the metal droplet begins to detach from the electrode tip, as shown in fig. 4(c). In fig. 4(d), most of the metal has entered in the liquid metal pool, accompanied with the perturbations of the slag/metal interface. The moment corresponding to figs. 4(a)-(d) will be referred as \( T_a, T_b, T_c, \) and \( T_d \) in the following content, respectively.

The presence of two phases in the present work makes the local value of the electrical conductivity vary in time and space. The distribution of the electric current coming from the top electrode is controlled by the distribution of the phase fraction, that is, the movement of the metal droplet, as shown in figs. 5(a)-(d), on which the volume fraction of the metal is superimposed. Because of the high conductivity of the metal droplet, most electrical current flows following the metal droplet, rather than through the slag. With the electrode melting, the metal droplet becomes larger and larger, the total resistance of ESR system decreases correspondingly during the process of metal droplet formation and dripping. When the metal droplets pass through the slag layer and entering into the metal pool, the volume fraction of the molten metal is small and the current density decrease.
The distributions of the Lorentz force at different instant are shown in figs. 6(a)-(d) with the volume fraction of the metal superimposed. The Lorentz force is produced by an interaction of the oscillating current and magnetic fields. In most of the metal region, the Lorentz force is radially inwards and the presence of the metal droplet have strong effects on the distribution of the Lorentz force. The direction of the Lorentz force imposed on the metal droplet is opposed to that of its movement, and the Lorentz force continuously increases as the metal droplet grows.

One of the roles of the slag in ESR is to produce the necessary Joule heating to melt electrode. The Joule heating per unit volume is determined from the current density and inversely proportional to the electrical conductivity of slag, hence ESR slag with lower electrical conductivity would provide higher Joule heating to melt the electrode. Figures 7(a)-(d) shows the contours of the Joule heating produced by the current as different instant. The Joule heating is concentrated especially at the location of the metal droplet. Further, near the corner of the electrode in the slag and just under the slag/metal interface, the Joule heating are also enhanced.

The temperature distribution in ESR system plays an important role to the organization and the quality of the ingot. Unlike the distribution of Joule heating, the temperature distribution in the slag layer and metal pool is the results of interaction between the productions of Joule heating and the metal droplet movements. The Joule heating carried by the metal droplet can be transferred to the region around them and the slag is hottest on the two sides beneath the electrode. It can also be seen that very large temperature gradients exist at the mold wall, as shown in figs. 8(a)-(d). At the moment of $T_a$, $T_b$, and $T_c$ during the dripping of the metal droplet,
Figure 7. Joule heating field distribution of ESR system at different time instant

Figure 8. Temperature field distributions of ESR system at different time instant
the temperature just below the electrode tip is slightly lower than the highest temperature region. While, at the moment of $T_d$, the temperature below the electrode tip is the highest.

In the slag phase, two symmetrical re-circulating cells are observed and these two vortices, is induced by electromagnetic forces, as shown in figs. 9(a)-(d). Beneath the electrode, some instabilities of a Rayleigh-Benard type arise because of Joule heating and cooling by contact with the electrode tip.

Conclusions

A 3-D mathematical model coupled with MHD model has been developed to investigate the two-phase flow behavior during the ESR process, in which the electromagnetic forces and the Joule heating are recalculated at each time step, which are incorporated into the momentum and energy equation as a source term, respectively. This model has been used to investigate the droplet formation and dripping as well as the corresponding electromagnetic field, velocity field and temperature field. The model predicted value is in good agreement with the experimental measurements in literature. The simulation results show that, during the melt dripping process, the resistance of ESR system is decreasing while the current density and velocity, as well as the temperature are increasing gradually. The maximum value of current density, elec-
tromagnetic force and velocity appears when the melt droplet detaches and the large droplets fall into the slag-metal interface. The maximum value of the current density is increased by almost one order of magnitude. The maximum value of the electromagnetic force and velocity increased by 2.5 and 4.7 times, respectively. The maximum value of Joule heating and temperature appears when the droplets fall into the slag-metal interface. The maximum value of the Joule heating and temperature increased for about 174.7% and 26.8% when compared with the moment of the melt droplet formation.

Acknowledgement

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Nomenclature

\begin{align*}
\mathbf{A} & \quad \text{magnetic potential vector, [Vsm}^{-1}\text{]} \\
\mathbf{B} & \quad \text{magnetic flux density, [T]} \\
c_r & \quad \text{specific heat at constant pressure, [Jkg}^{-1}\text{K}^{-1}\text{]} \\
\mathbf{E} & \quad \text{electric field intensity, [Vm}^{-1}\text{]} \\
\mathbf{E}_b & \quad \text{buoyancy force, [Nm}^{-3}\text{]} \\
\mathbf{E}_l & \quad \text{Lorentz force, [Nm}^{-3}\text{]} \\
f_l & \quad \text{liquid fraction} \\
H & \quad \text{enthalpy of the metal and the slag, [J]} \\
\mathbf{H} & \quad \text{magnetic field intensity, [Am}^{-2}\text{]} \\
h & \quad \text{sensible enthalpy, [J]} \\
J & \quad \text{current density, [Am}^{-2}\text{]} \\
k_{eff} & \quad \text{effective thermal conductivity, [Wm}^{-1}\text{K}^{-1}\text{]} \\
L & \quad \text{latent heat of fusion, [Jkg}^{-1}\text{]} \\
p & \quad \text{pressure, [Pa]} \\
Q_{Joule} & \quad \text{Joule heating, [Wm}^{-3}\text{]} \\
r & \quad \text{radius of the electrode, [m]} \\
T & \quad \text{temperature, [K]} \\
T_l & \quad \text{liquidus temperature, [K]} \\
T_s & \quad \text{solidus temperature, [K]} \\
t & \quad \text{time, [s]} \\
\mathbf{u} & \quad \text{velocity, [ms}^{-1}\text{]} \\
x, y, z & \quad \text{Cartesian co-ordinates} \\
\alpha & \quad \text{volume fraction of metal} \\
\mu_{eff} & \quad \text{turbulent viscosity, [Pa}^{-1}\text{s]} \\
\mu_0 & \quad \text{vacuum permeability, [Hm}^{-1}\text{]} \\
\rho & \quad \text{density of fluid, [kgm}^{-3}\text{]} \\
\sigma & \quad \text{electric conductivity, [\Omega}^{-1}\text{m}^{-1}\text{]} \\
\phi & \quad \text{electric potential, [Vm}^{-1}\text{]} \\
\end{align*}

References